

ATOC 5051 INTRODUCTION TO PHYSICAL OCEANOGRAPHY

Lecture 11

Learning objectives: understand internal waves, explain storm surge, and identify the effect of the Earth's rotation

1. Internal waves; concepts -barotropic and baroclinic modes;
2. Storm surges
3. Effects of rotation, Rossby radius of deformation

Previous classes: we discussed surface gravity waves. What assumptions did we make to isolate “*surface*” gravity waves?

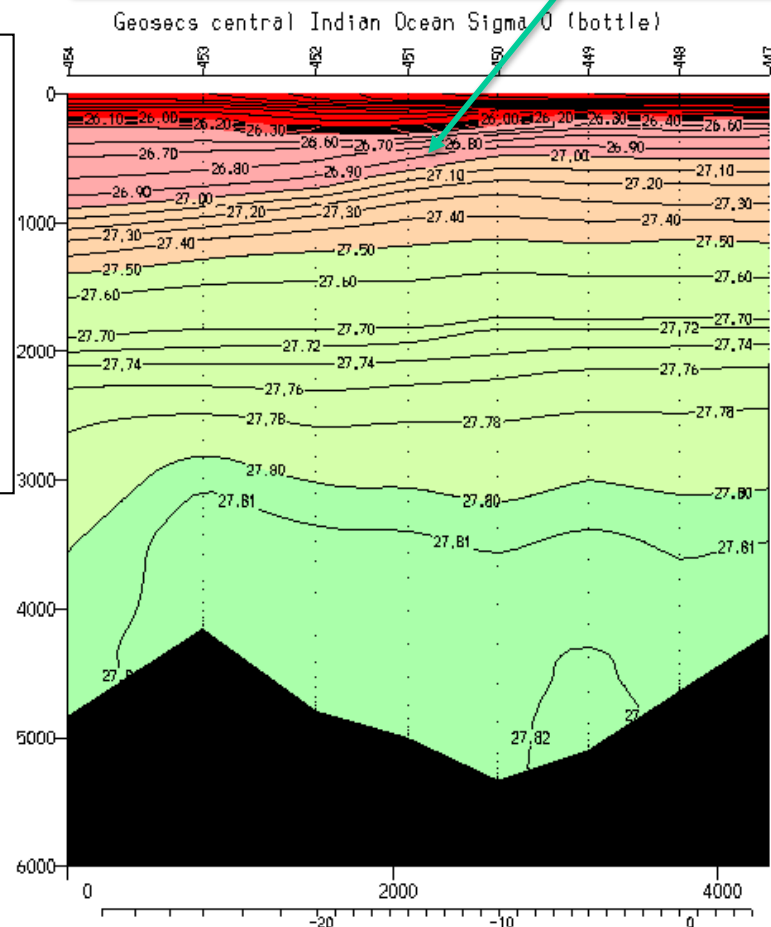
Constant density – homogeneous ocean

Observed ocean density:

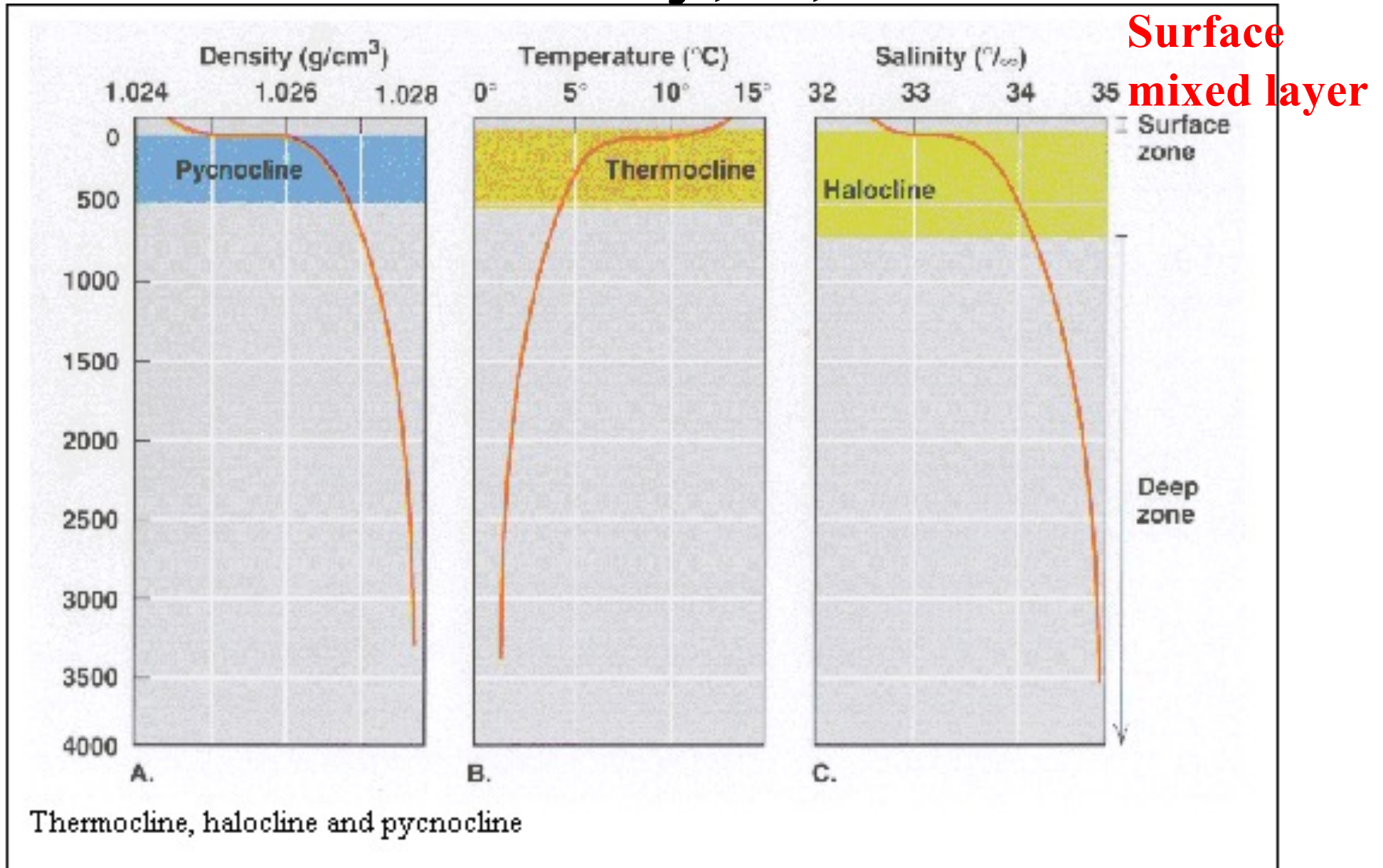
$$\text{Sigma}(0) = (\text{density} - 1000) \text{ kg/m}^3$$

Potential density relative to surface

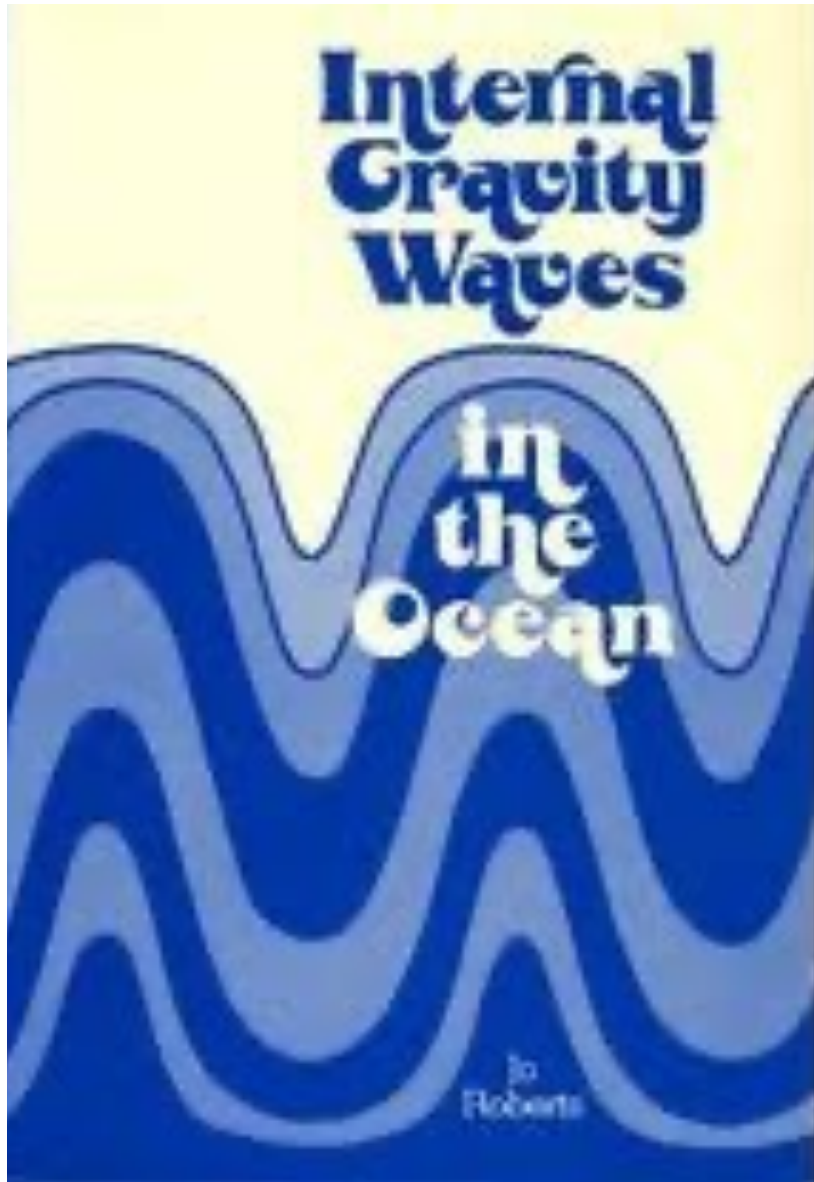
Sharp vertical density gradient



Concepts: Vertical profiles of density, T, S



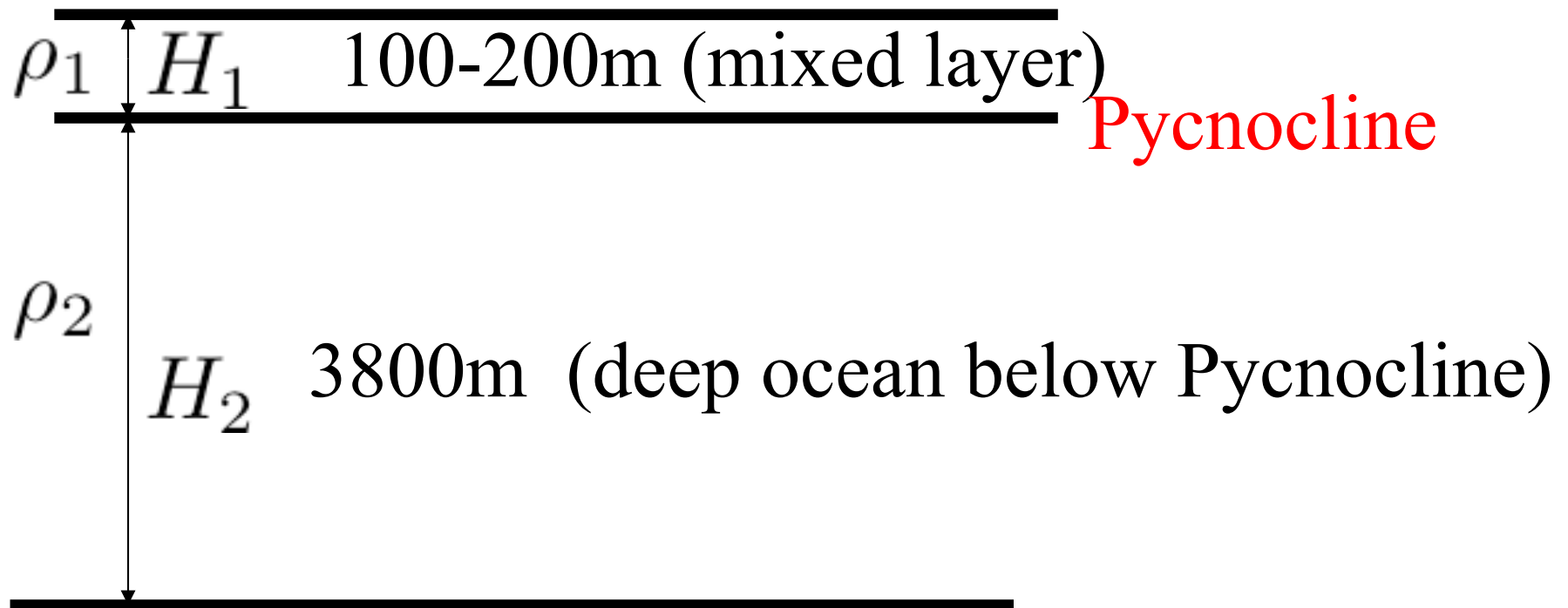
1. Internal gravity waves in density stratified ocean



**Largest amplitudes:
in pycnocline**

Internal gravity waves in density stratified ocean

a) A 2-layer model



Barotropic and baroclinic modes

The 2-layer system has 2 vertical normal modes.

Total Solution is the superposition of the two modes.

Barotropic mode: independent of z , which represents vertically-averaged motion.

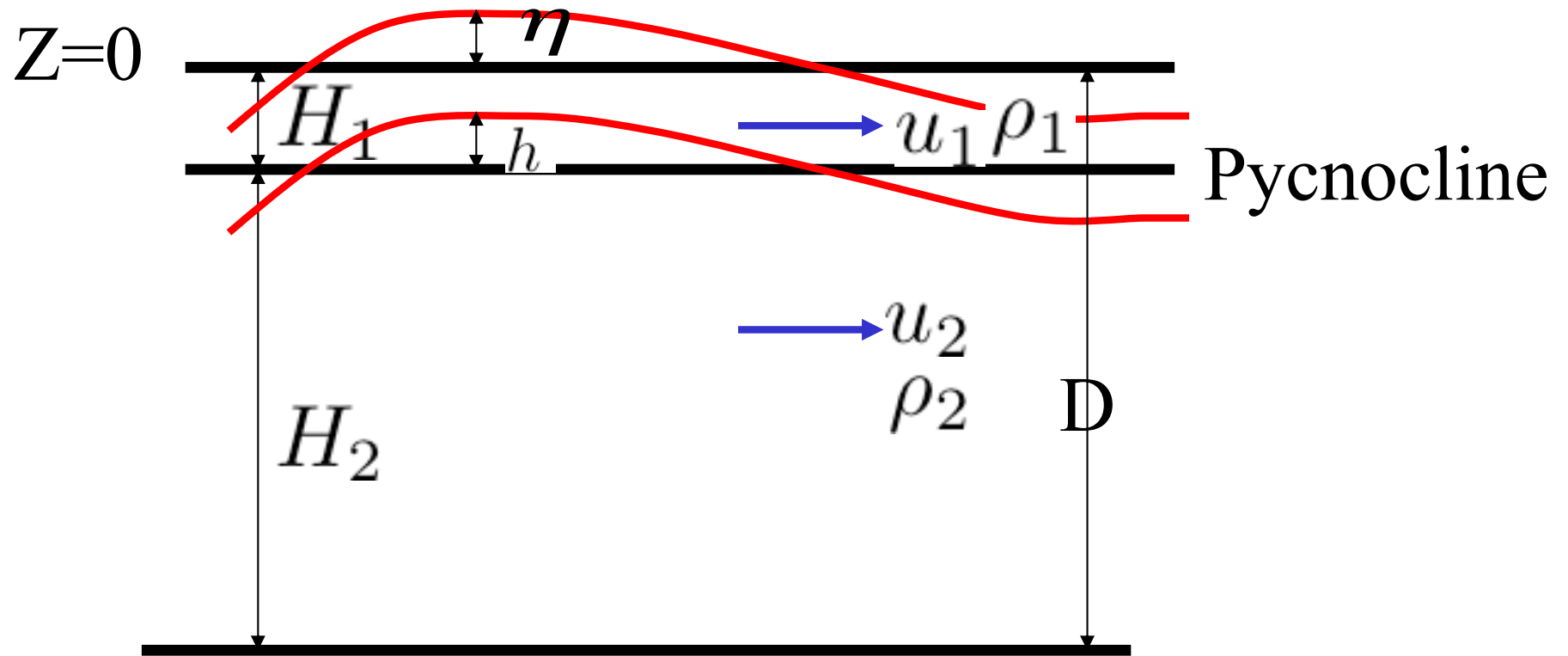
Restoring force: gravity g

Baroclinic mode: vertical shear flow, and vertically-integrated transport is zero. **Restoring force is reduced**

gravity,
$$g' = \frac{g(\rho_2 - \rho_1)}{\rho_2} \sim 0.03m/s^2$$

Why? At ocean surface – air density is \ll water density & thus ignored, and restoring force is gravity (g); in the pycnocline, the restoring force is reduced gravity g' : due to the small difference between the water density above and below

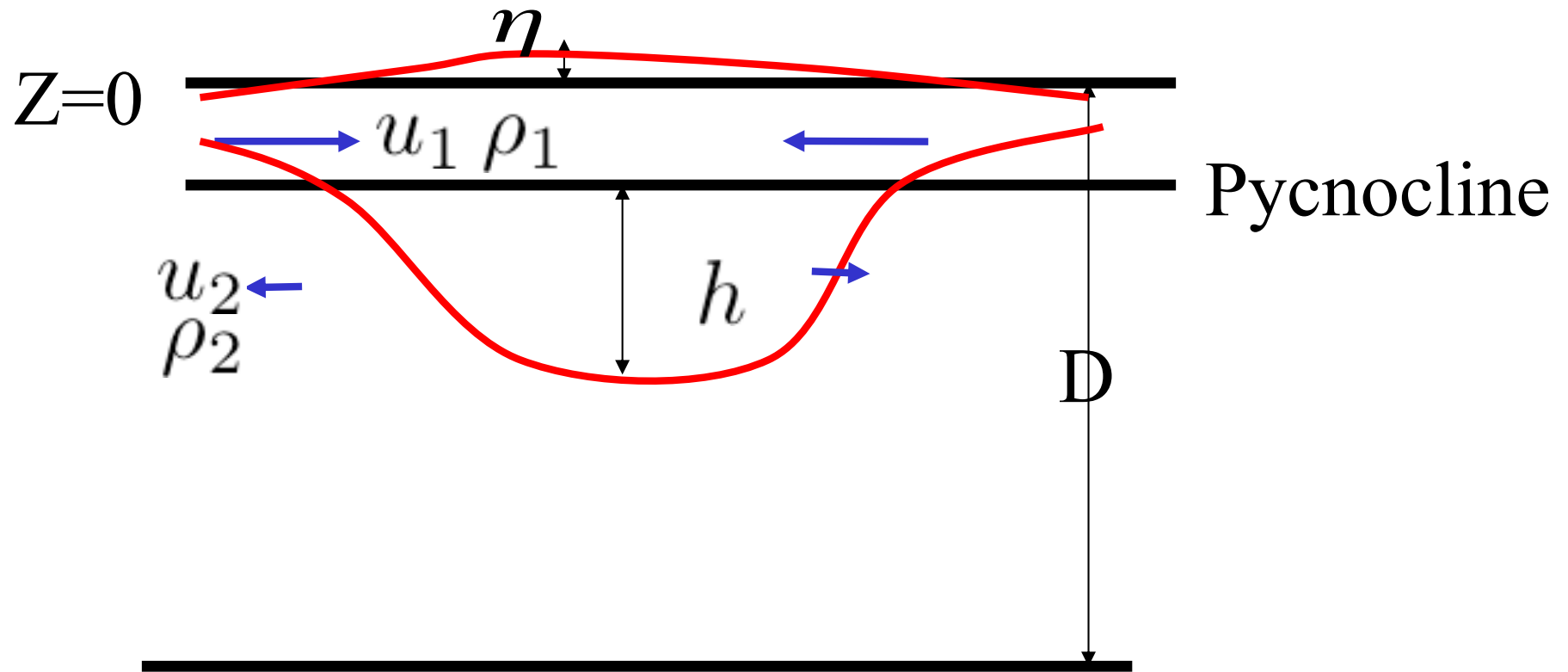
Barotropic mode



$$C_0 = \sqrt{g(H_1 + H_2)} = \sqrt{gD} \sim 200 \text{ m/s}$$

$$\frac{h}{\eta} = \frac{H_2}{H_1 + H_2} \leq 1 \sim 1, \quad \frac{u_2}{u_1} \sim 1$$

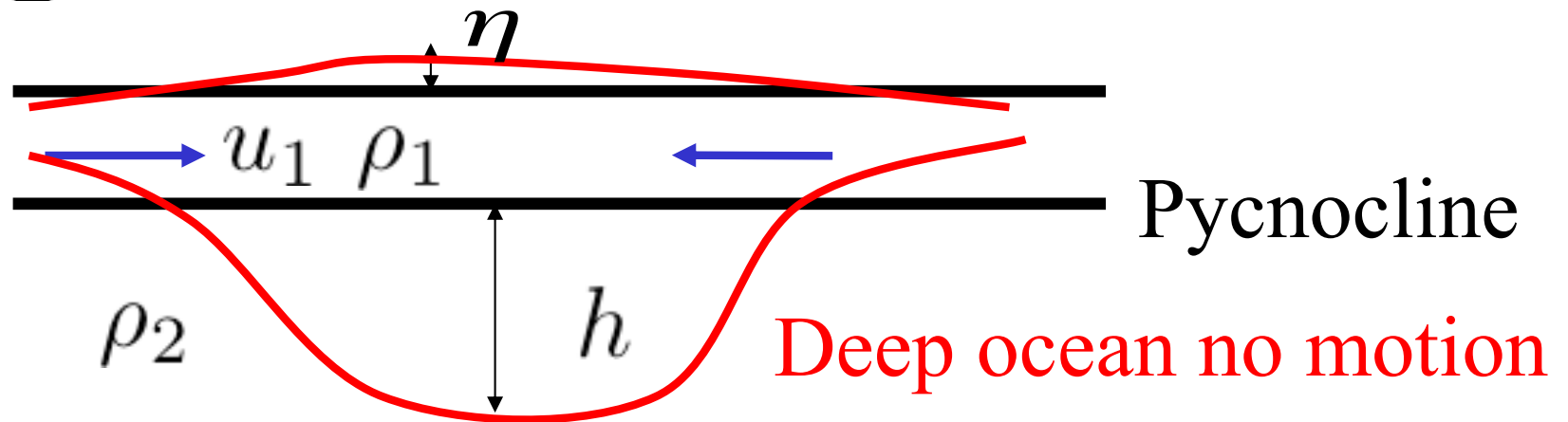
Baroclinic mode



$$C_1 = \sqrt{g' \frac{H_1 H_2}{H_1 + H_2}} \sim 2 - 3 \text{ m/s}$$

$$\frac{h}{\eta} = -\frac{g}{g'} \sim -300 \text{ (internal gravity waves)}$$

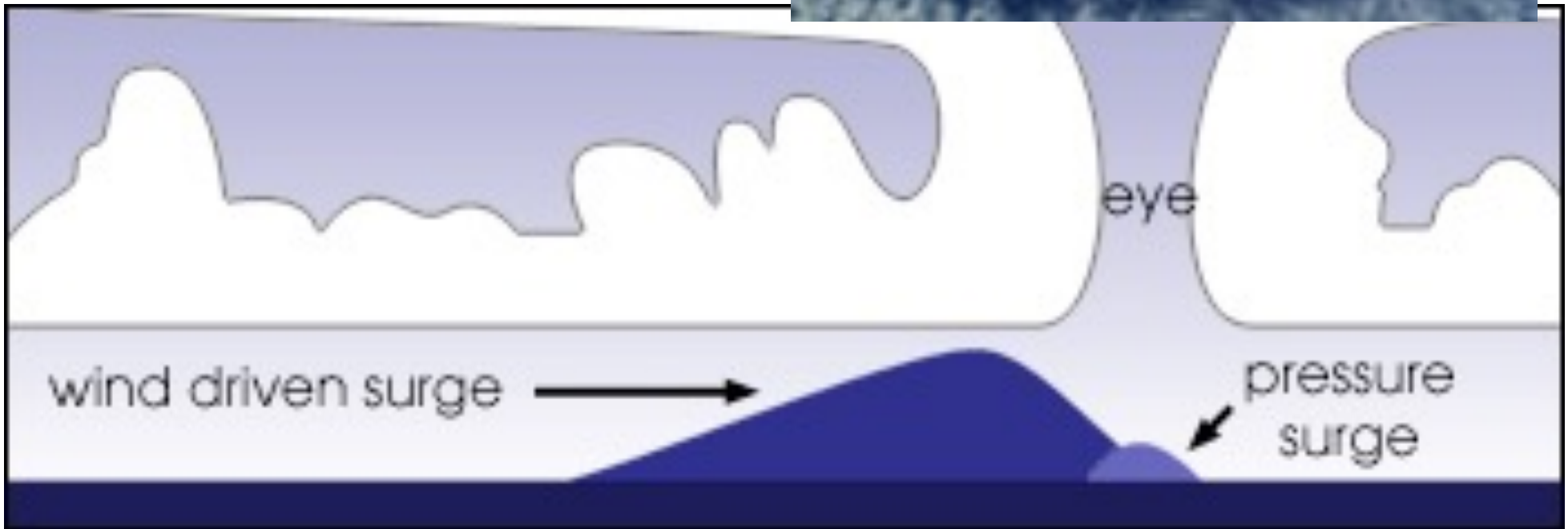
A $1\frac{1}{2}$ - layer model



Deep Ocean: infinitely deep and $\nabla p = 0$

The system has only **one baroclinic mode**;
No barotropic mode since we assumed $\nabla p = 0$
below the pycnocline. *We can also view it as
the deep ocean is infinitely deep.*

2. Storm surge



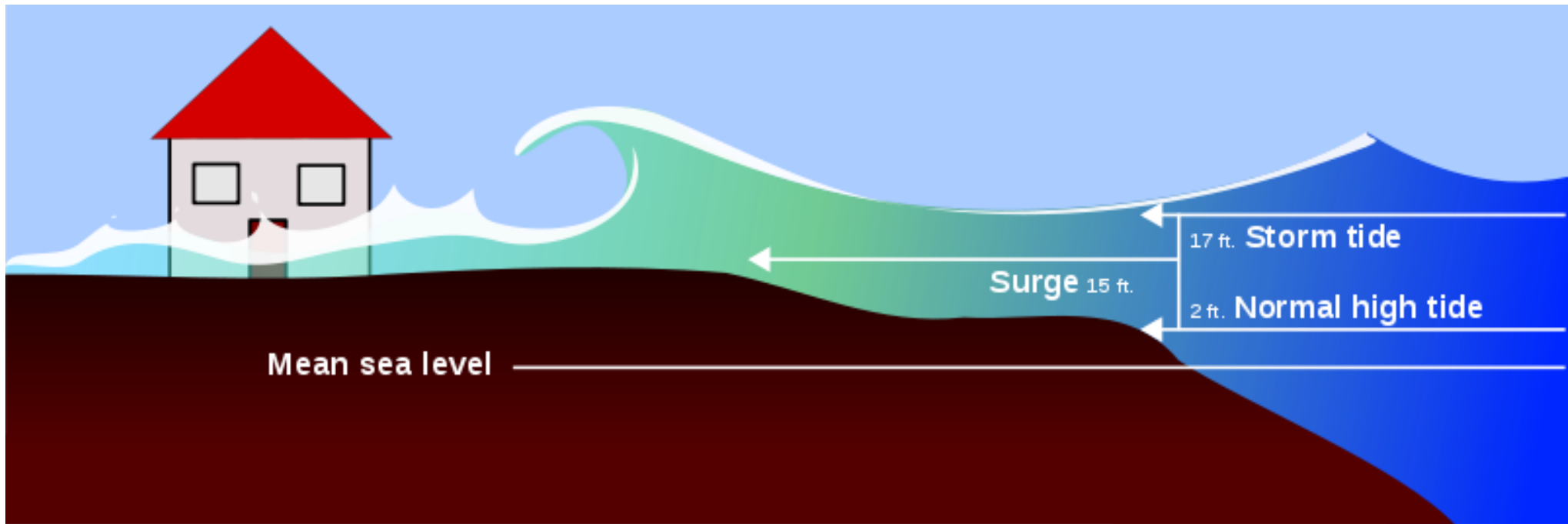
Coastal shallow water: amplify. **Why?**

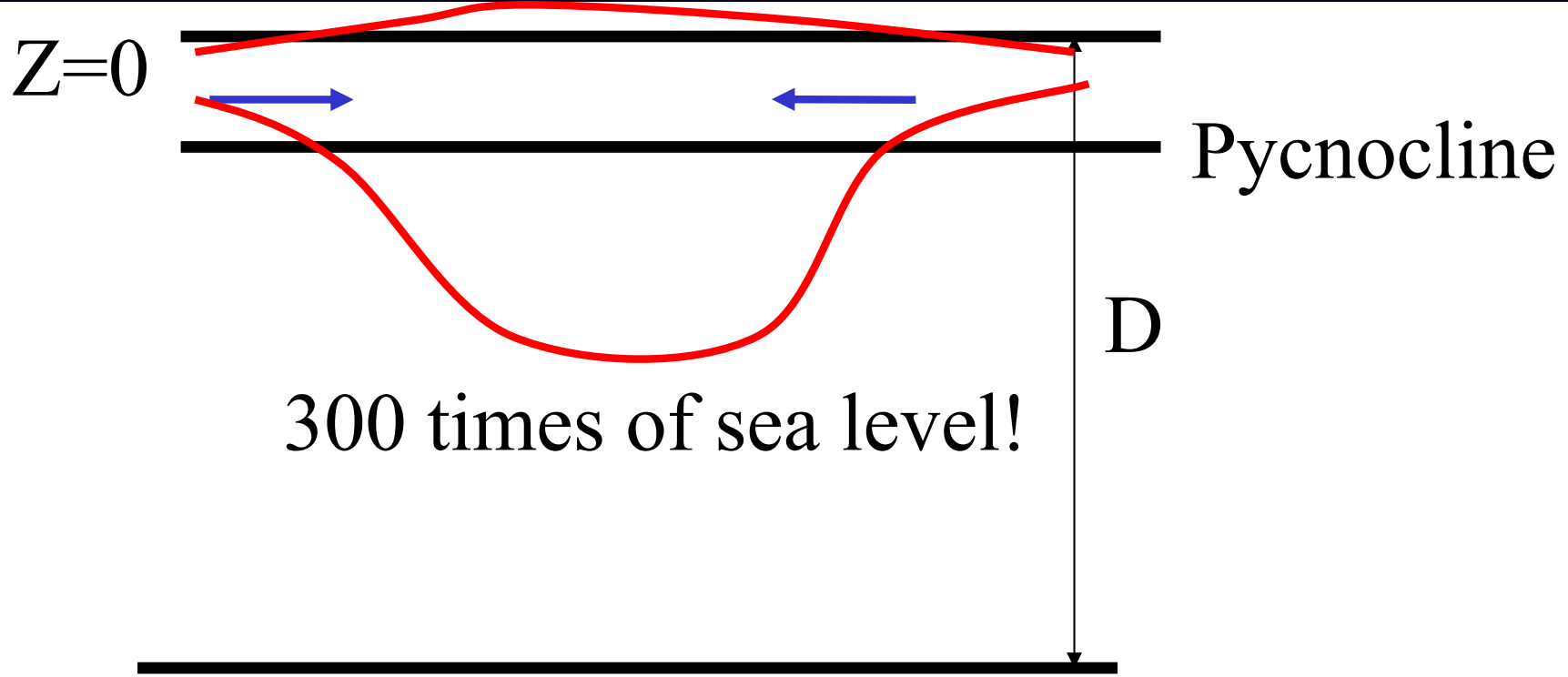
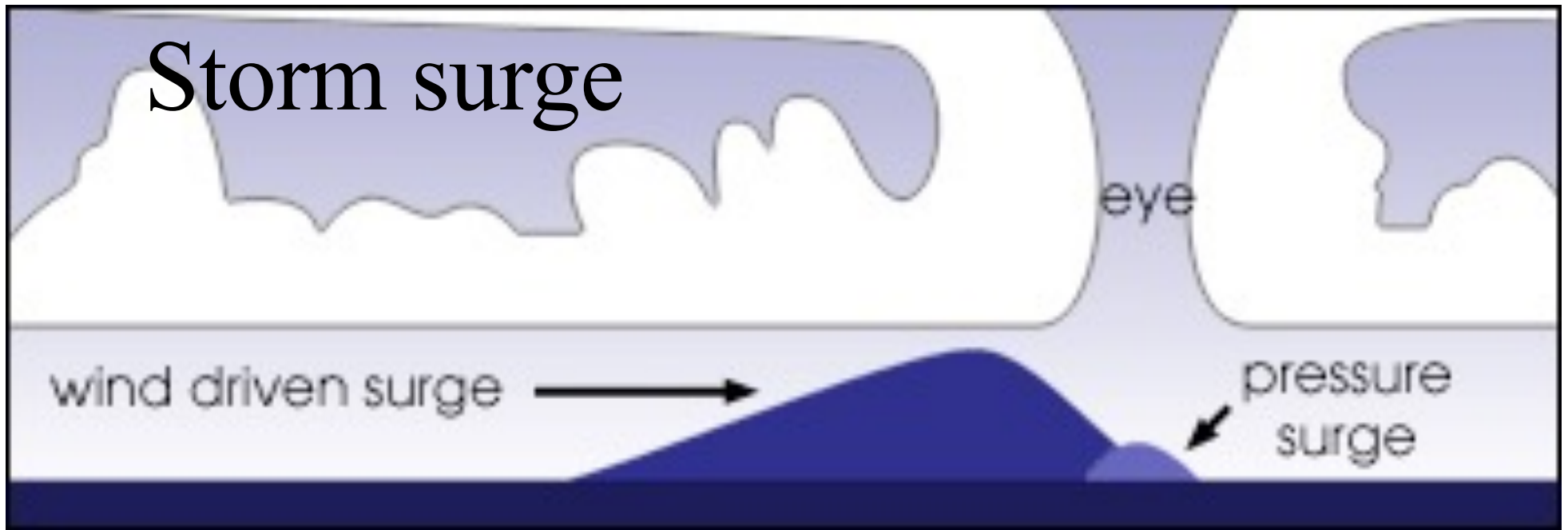
Basic dynamics :

- (i) winds associated with storms or hurricanes pile up the water – **wind surge**;
- (ii) low sea level pressure at the storm center (minimal comparing with wind) – **pressure surge**;
- (iii) Shallow, gently sloping coastal region: intensify;
- (iv) Overlapping with high tide – more devastating.

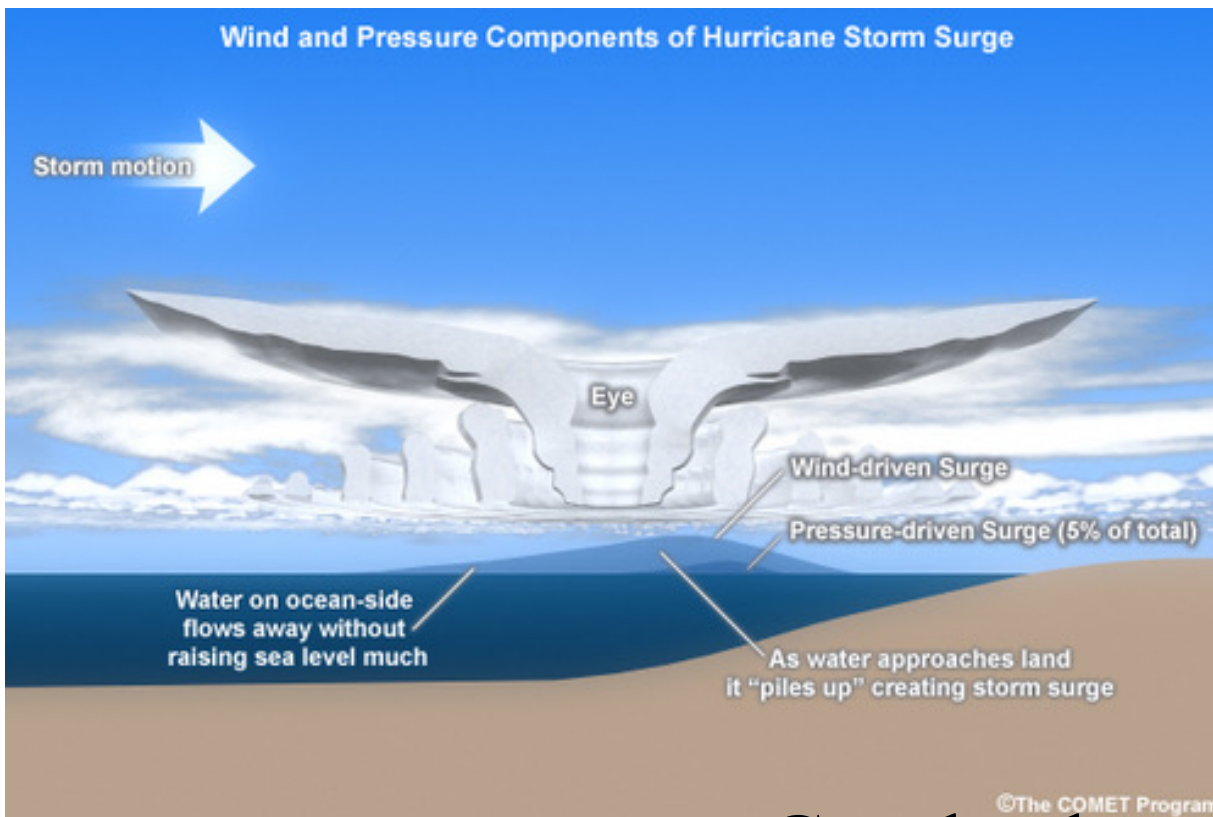
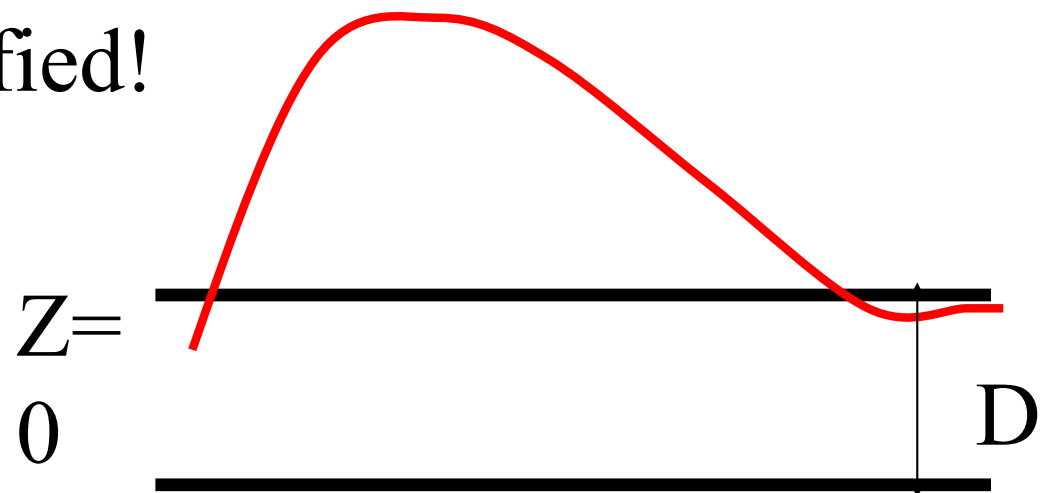
Very complicated, depend on many factors: Storm surge is a very complex phenomenon because it is sensitive to the slightest changes in storm intensity, forward speed, size (radius of maximum winds-RMW), angle of approach to the coast, central pressure (minimal contribution in comparison to the wind), and the shape and characteristics of coastal features such as bays and estuaries.

Surge + high tide





Approach the coast, D is shallow: sea level has to go up, surge amplified!



Gentle slope: amplify

Surge examples:

- The highest storm surge in record: 1899 Cyclone Mahina: 13 meters (43 feet) storm surge at Bathurst Bay, Australia (high tide);
- In the U.S., the greatest storm surge was generated by Hurricane Katrina: 9 meters (30feet) high storm surge in Bay St. Louis, Mississippi, and surrounding counties. (Low elevation above sea level, larger impact)

Hurricane Katrina
Near peak strength:
Aug 28, 2005

Formed: Aug 23;
Dissipated: Aug 31

Highest: 175mph
Lowest pressure:
902mbar



- Damages: \$81.2 billion (costliest Atlantic hurricane in history), the 6th strongest hurricane;
- Fatalities: greater than 1836 total;
- Areas affected: Bahamas, South Florida, Cuba, Louisiana (especially greater New Orleans), Mississippi, Alabama, Florida Panhandle, most of the eastern North America.



Aftermath of Katrina



Storm Surge video: NOAA National Weather Service:

https://www.youtube.com/watch?v=2GgUn2QTJtE&feature=emb_rel_end

Sea, Lake, and Overland Surges from Hurricanes (SLOSH)

3. Effects of rotation ($f \neq 0$) and Rossby radius of deformation

With $f=0$, what transient waves are available in the system?

What is the equilibrium state of the ocean after the waves propagation?

Critical thinking: what effects do you think f will have on the transient waves and equilibrium state?

Effects of rotation ($f \neq 0$) and Rossby radius of deformation

With a uniform rotation (f is assumed to be a constant), the equations of motion for the unforced, inviscid ocean are:

$$\frac{dU}{dt} - fV = -\frac{1}{\rho_0} \frac{\partial P}{\partial x},$$

$$\frac{dV}{dt} + fU = -\frac{1}{\rho_0} \frac{\partial P}{\partial y},$$

$$\frac{dW}{dt} = -\frac{1}{\rho_0} \frac{\partial P}{\partial z} - g.$$

$$\frac{\partial U}{\partial x} + \frac{\partial V}{\partial y} + \frac{\partial W}{\partial z} = 0.$$

Assumptions:

(i) f constant;

(ii) $\rho = \rho_0 = \text{constant}$;

(iii) $\frac{\partial P}{\partial z} = -\rho_0 g$.
Total P

(iv) $W = 0$, at $z=H$;
bottom

(v) Background state:

$$U_0 = V_0 = W_0 = 0$$

$$(\text{Ro} \ll 1, \text{E} \ll 1)$$

For small perturbations u, v, w, p about the resting state, we have:

$$U = U_0 + u = u,$$

$$V = V_0 + v = v,$$

$$W = W_0 + w = w,$$

$$P = P_0(z) + p, \longrightarrow \frac{\partial P}{\partial z} = -\rho_0 g.$$

$$\rho = \rho_0.$$

The linearized first order equations for perturbation:

$$\frac{\partial u}{\partial t} - fv = -\frac{1}{\rho_0} \frac{\partial p}{\partial x},$$

$$\frac{\partial v}{\partial t} + fu = -\frac{1}{\rho_0} \frac{\partial p}{\partial y},$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0.$$

Apply boundary conditions

$$Z=0, \quad p = \rho_0 g \eta, \quad w = \frac{\partial \eta}{\partial t}$$

$$Z=H, \quad w = 0;$$

and vertically integrate the perturbation equations:

$$\frac{\partial u}{\partial t} - f v = -g \frac{\partial \eta}{\partial x},$$

$$\frac{\partial v}{\partial t} + f u = -g \frac{\partial \eta}{\partial y},$$

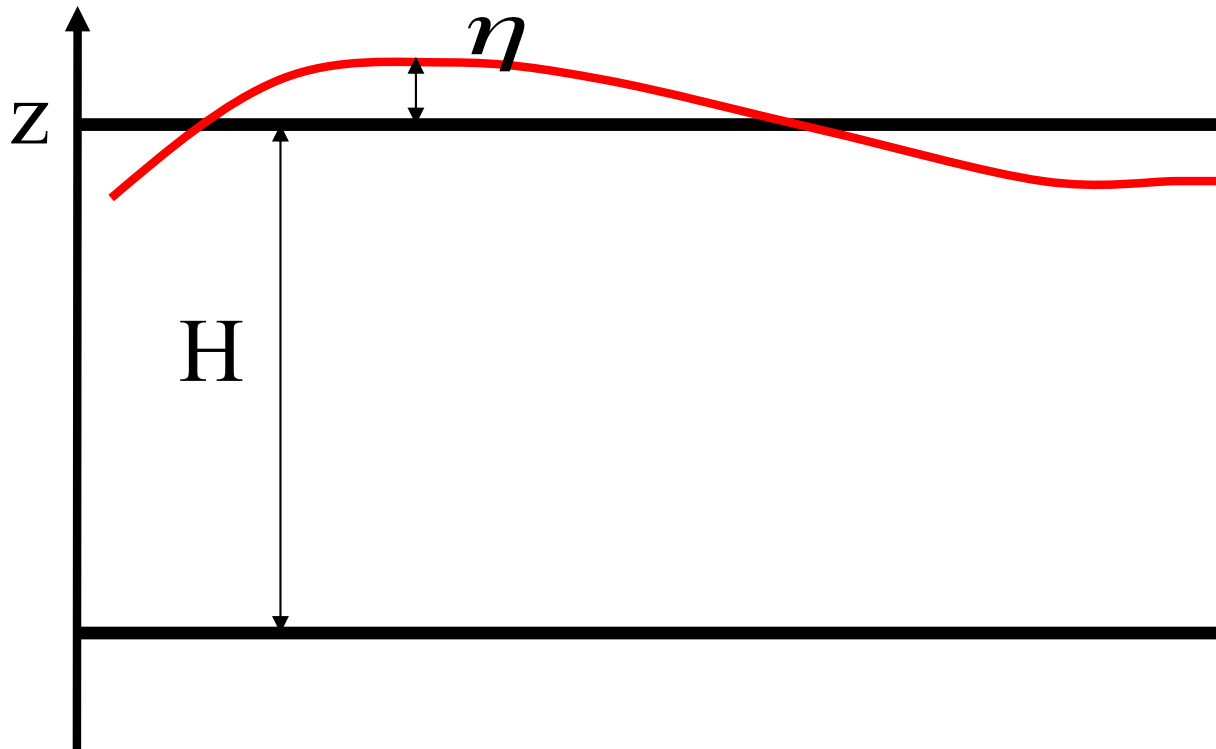
$$\frac{\partial \eta}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0.$$

Following the same procedure as in the non-rotating case, write an equation in η ,

$$\frac{\partial^2 \eta}{\partial t^2} - c^2 \nabla^2 \eta + fH\zeta = 0,$$

Where $c^2 = gH, \zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$

Relative
vorticity



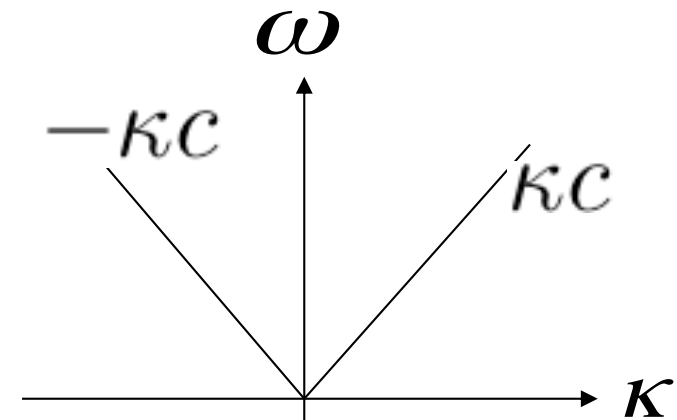
Note: $f + \zeta$ is referred to as absolute vorticity;

Planetary
vorticity

Relative
vorticity

a) Non-rotating case ($f=0$):

$$\frac{\partial^2 \eta}{\partial t^2} - c^2 \nabla^2 \eta = 0$$



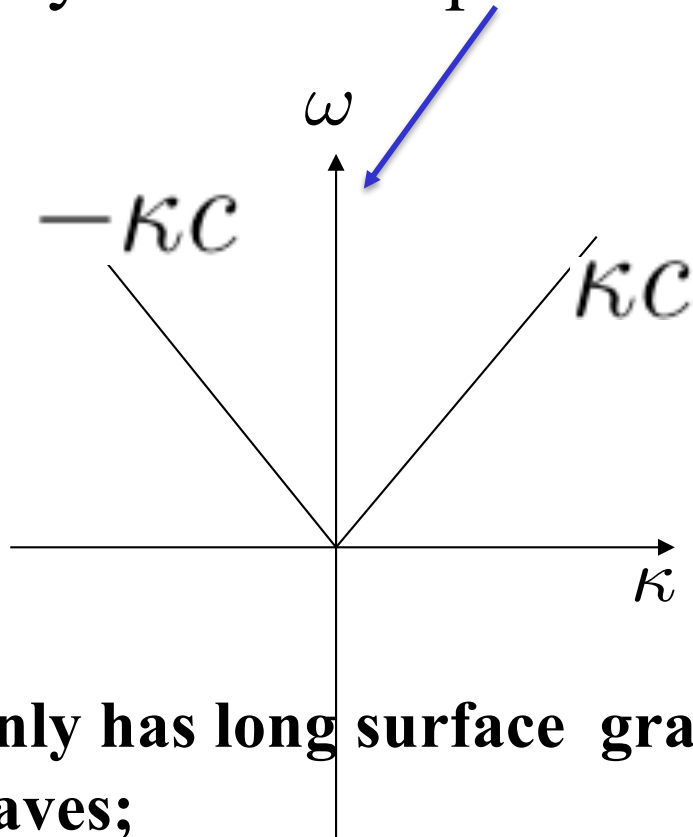
Dispersion curve

Assume $\eta = \eta_0 \cos(kx - \omega t)$ (1-dimensional exp)

$$\omega^2 = c^2 k^2 \longrightarrow$$
$$c^2 = gH$$

Recall this is the dispersion relation for long-surface gravity waves when $f=0$!

Today's class: dispersion relation

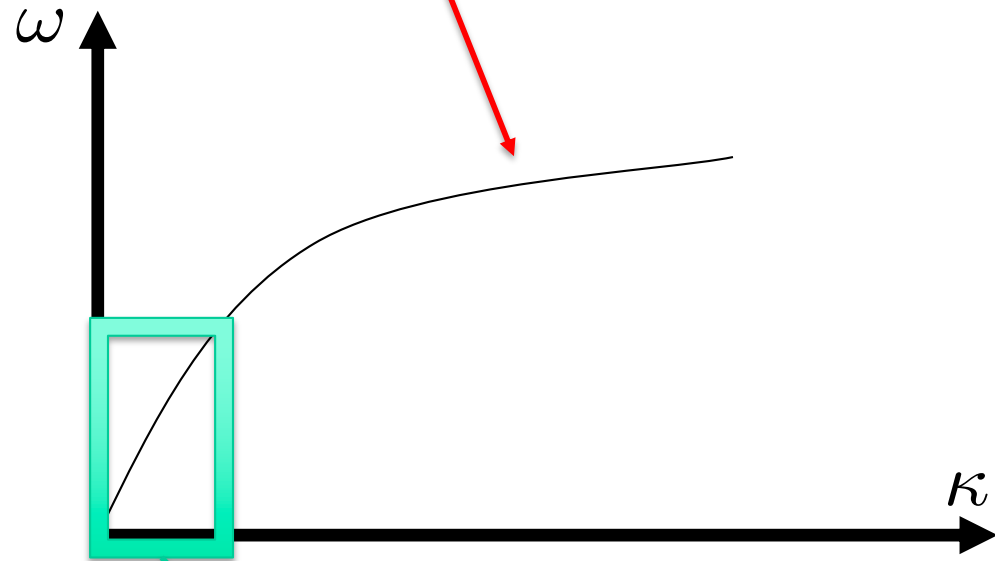


Only has long surface gravity waves;

short waves: **distorted by hydrostatic approximation**

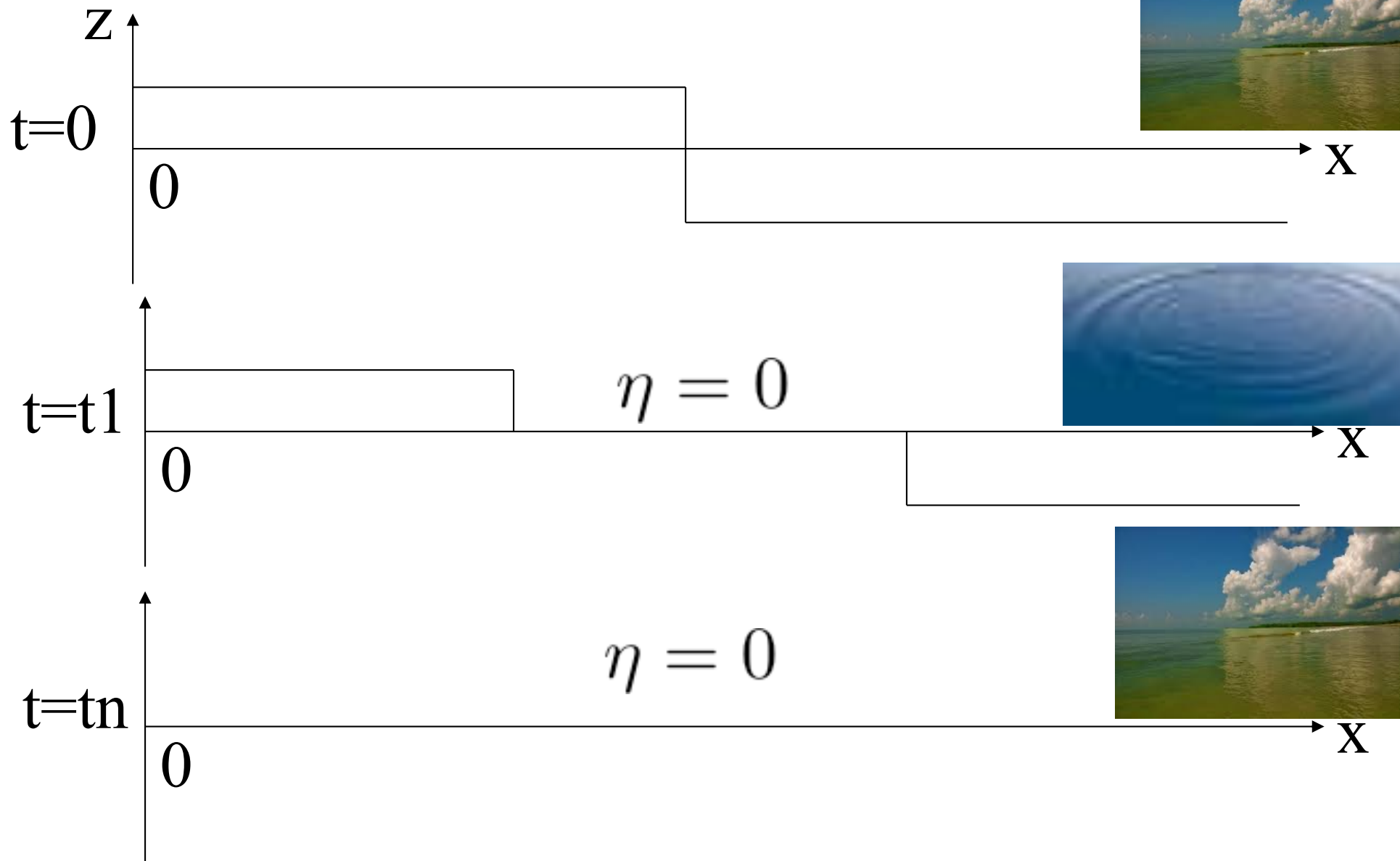
→ $\frac{\partial P}{\partial z} = -\rho_0 g.$

Previous: dispersion relation



$L = \frac{2\pi}{\kappa}$: long surface gravity wave with $f=0$

$$\omega^2 = c^2 k^2 \quad \omega = \pm kc \quad f=0$$



b) Rotating case ($f \neq 0$)

Assume wave form of solution

$$\eta = \eta_0 e^{i(kx + ly - \omega t)},$$

$$u = u_0 e^{i(kx + ly - \omega t)},$$

$$v = v_0 e^{i(kx + ly - \omega t)},$$

Substitute into the vertically-integrated
perturbation equation for η, u, v

$$\frac{\partial u}{\partial t} - fv = -g \frac{\partial \eta}{\partial x},$$

$$\frac{\partial v}{\partial t} + fu = -g \frac{\partial \eta}{\partial y},$$

$$\frac{\partial \eta}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0.$$

let coefficient matrix = 0,

$$\longrightarrow \omega^2 = \kappa^2 c^2 + f^2, \quad \text{where } \kappa^2 = k^2 + l^2$$

$$\begin{vmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{vmatrix} =$$

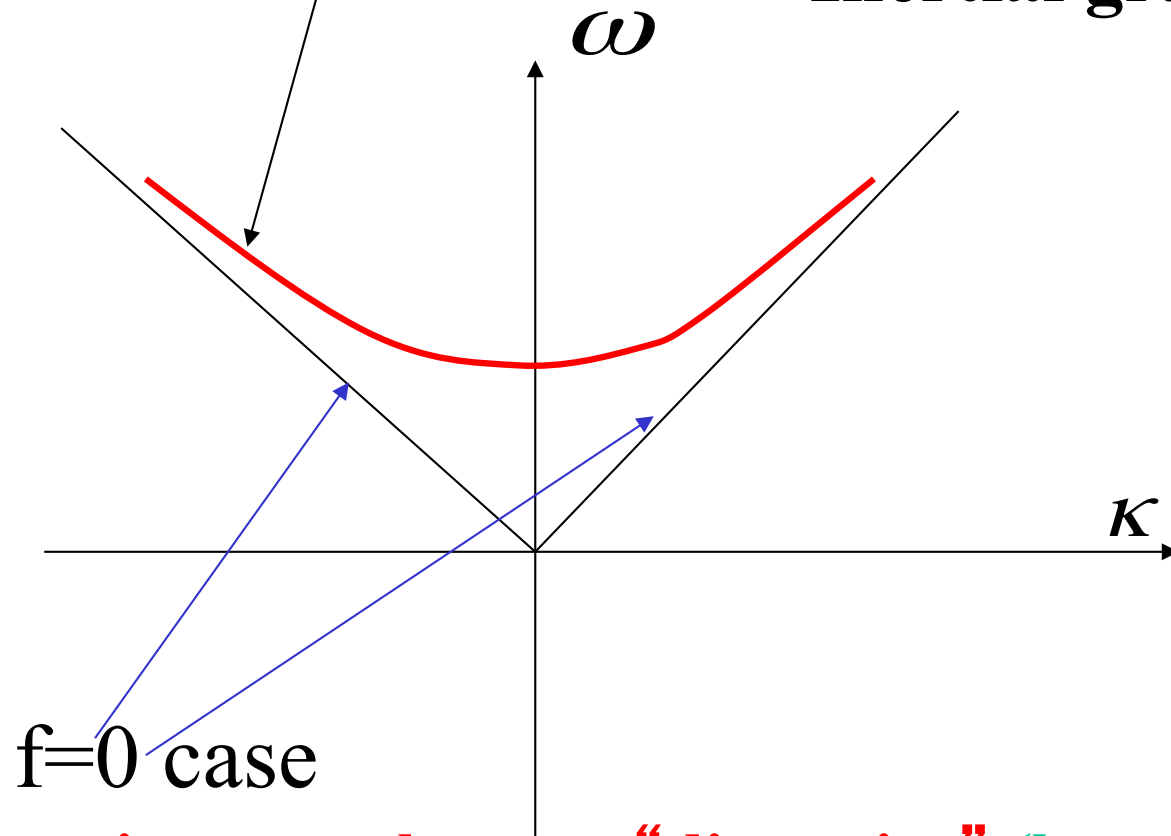
$$a_{11}a_{22}a_{33} + a_{12}a_{23}a_{31} + a_{13}a_{21}a_{32} \\ - a_{11}a_{23}a_{32} - a_{12}a_{21}a_{33} - a_{13}a_{22}a_{31}$$

$$\rightarrow \omega^2 = \kappa^2 c^2 + f^2,$$

What does f do to the system?

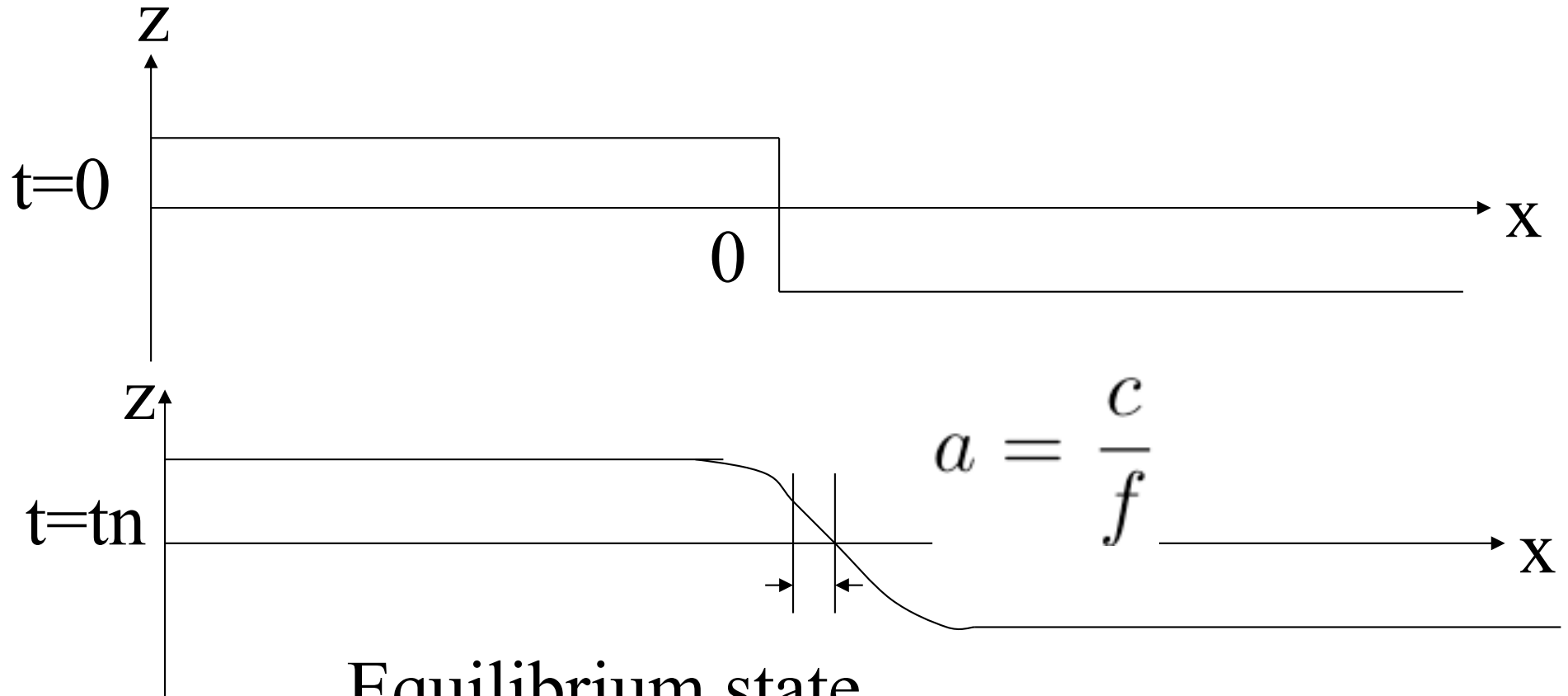
$$\omega^2 = \kappa^2 c^2 + f^2,$$

Inertial gravity waves



- (i) Long gravity waves become “dispersive” (board demo)
- (ii) Long gravity waves do not have “zero” frequency anymore. Their lowest frequency is “ f ”, which has a period of a few days in mid latitude.

Adjustment with f (1-dimensional in x)



Equilibrium state

$$\cancel{\frac{\partial^2 \eta}{\partial t^2}} - c^2 \frac{\partial^2 \eta}{\partial x^2} + fH \frac{\partial v}{\partial x} = 0 \longrightarrow$$

Geostrophic balance $f v = g \frac{\partial \eta}{\partial x}$

Solutions:

$$\eta = \eta_0 \left[-1 + \exp\left(-\frac{x}{a}\right) \right], \text{ for } x > 0,$$

$$\eta = \eta_0 \left[1 - \exp\left(\frac{x}{a}\right) \right], \text{ for } x < 0,$$

where $a = \frac{c}{f}$ is Rossby radius of deformation.

