• The time average Poynting vector (average power density) does not depend on time, and can be written

as

$$\mathbf{W}_{av}(x, y, z) = [\mathscr{W}(x, y, z; t)]_{av} = \frac{1}{2} \operatorname{Re}[\mathbf{E} \times \mathbf{H}^*] \qquad (W/m^2)$$

Where, average power density is obtained by integrating the instantaneous Poynting vector over one period and dividing by the period.

E and H fields represent peak values, complex fields

Instantaneous Poynting vector represents the energy transfer per unit area per unit time of an electromagnetic field.

$$\mathcal{W} = \mathcal{E} \times \mathcal{H} = \frac{1}{2} \operatorname{Re}[\mathbf{E} \times \mathbf{H}^*] + \frac{1}{2} \operatorname{Re}[\mathbf{E} \times \mathbf{H}e^{j2\omega t}]$$
Avg=0

Average power radiated by an antenna (radiated power)

$$P_{rad} = \iint W_{rad} \, ds$$

$$P_{rad} = P_{ave} = \int_{0}^{2\pi} \int_{0}^{\pi} \left(\frac{1}{2} \operatorname{Re}\left[\underline{E} \times \underline{H}^{*}\right] \right) \cdot \hat{a}_{r} r^{2} \sin\theta d\theta d\phi$$
(2-9)

For Isotropic radiator which radiates equally in all directions(not exist but used as reference to compared with other antenna), the power density equal to

$$\mathbf{W}_0 = \hat{\mathbf{a}}_r W_0 = \hat{\mathbf{a}}_r \left(\frac{P_{\text{rad}}}{4\pi r^2}\right) \quad (W/\text{m}^2)$$

Example 1

A hypothetical isotropic antenna is radiating in free pace. At a distance of 100 m from the antenna, the total electric field (E_{θ}) is measured to be 5 V/m. Find the (a) Power density (W_{rad}) (b) Power radiated (P_{rad})

(a)
$$W_{rad} = \frac{1}{2} [E \times H^*] = \frac{E^2}{2\eta} \hat{a}_r = \frac{5^2 \hat{a}_r}{2(120\pi)} = 0.03315 \hat{a}_r \text{ watts/m}^2$$

(b) Since isotropic:

Prad = $w_{rad}^* 4\pi r^2$

 $=0.03315*4\pi*100^{2}=4165.75$ watts

Note E depend on r so prad always const and wrad decrease as r increase

Example 2

$$\underline{\text{Example 2.2}}$$

$$\underline{W}_{ave} = \underline{W}_{rad} = \hat{a}_r A_0 \frac{\sin \theta}{r^2}$$

$$P_{ave} = P_{rad} = \bigoplus_s \underline{W}_{ave} \cdot d\underline{s}$$

$$= \int_0^{2\pi} \int_0^{\pi} \left[\hat{a}_r A_0 \frac{\sin \theta}{r^2} \right] \cdot \left[\hat{a}_r r^2 \sin \theta d\theta d\phi \right]$$

$$P_{ave} = P_{rad} = A_0 \int_0^{2\pi} \int_0^{\pi} \sin^2 \theta d\theta d\phi = 2\pi A_0 \int_0^{\pi} \sin^2 \theta d\theta$$

$$P_{rad} = A_0 \pi^2$$

Radiation intensity: defined as

The power radiated from antenna per unit solid angle $U(\Theta, \phi)$

$$U(\theta,\phi) = \frac{r^2}{2\eta} |\mathbf{E}(r,\theta,\phi)|^2 \simeq \frac{r^2}{2\eta} \left[|E_\theta(r,\theta,\phi)|^2 + |E_\phi(r,\theta,\phi)|^2 \right]$$
(2-12a)

where

 $\mathbf{E}(r, \theta, \phi) = \text{far-zone electric-field intensity of the antenna}$ $E_{\theta}, E_{\phi} = \text{far-zone electric-field components of the antenna}$ $\eta = \text{intrinsic impedance of the medium}$ U = radiation intensity (W/unit solid angle)

The total power is obtained by integrating the radiation intensity, as given by (2-12), over the entire solid angle of 4π . Thus

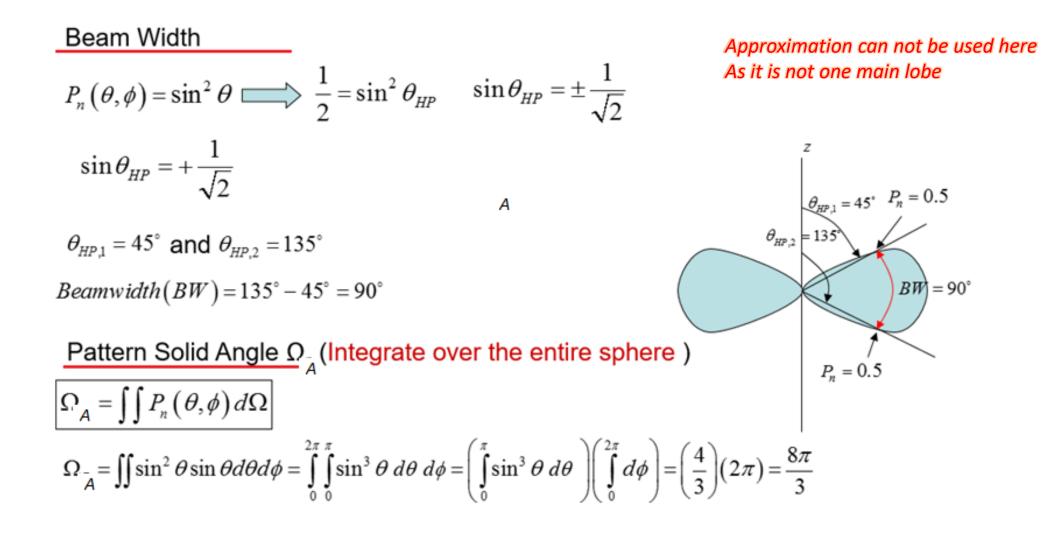
$$P_{\rm rad} = \oint_{\Omega} U \, d\Omega = \int_0^{2\pi} \int_0^{\pi} U \sin\theta \, d\theta \, d\phi \tag{2-13}$$

radiation intensity of an isotropic source
$$U_0 = \frac{P_{\text{rad}}}{4\pi}$$

| W rad = | ¥ watt∕m² |
|--------------------------|------------------------------------|
| $\mathbf{U} \Rightarrow$ | watt/solid angle = watt/ m^2/r^2 |
| U = | r ² W _{rad} |

Example 3

Compute Beam width, Beam solid angle and the power radiated through HPBW for normalized field pattern Fn=sin0



Half-power Pattern Solid Angle $\Omega_{A,HP}$ (*integrate over half power Beam width*)

$$\begin{split} & \boxed{\Omega_{A,HP} = \iint P_n(\theta, \phi) d\Omega} \\ & \Omega_{A,HP} = \iint \sin^2 \theta \sin \theta d\theta d\phi = \int_0^{2\pi} \int_{45^\circ}^{135^\circ} \sin^3 \theta \ d\theta \ d\phi = \left(\int_{45^\circ}^{135^\circ} \sin^3 \theta \ d\theta \ \right) \left(\int_0^{2\pi} d\phi \right) = \left(\frac{5}{3\sqrt{2}}\right) (2\pi) = \frac{5\pi\sqrt{2}}{3} \\ & \int_{45^\circ}^{135^\circ} \sin^3 \theta \ d\theta = \left[-\cos\theta + \frac{\cos^3 \theta}{3}\right]_{45^\circ}^{135^\circ} = \left[\left(-\cos(135^\circ) + \frac{\cos^3(135^\circ)}{3}\right) - \left(-\cos(45^\circ) + \frac{\cos^3(45^\circ)}{3}\right)\right] \\ & = \left[\left(\frac{1}{\sqrt{2}} - \frac{1}{6\sqrt{2}}\right) - \left(-\frac{1}{\sqrt{2}} + \frac{1}{6\sqrt{2}}\right)\right] = \frac{2}{\sqrt{2}} - \frac{2}{6\sqrt{2}} = \frac{10}{6\sqrt{2}} = \frac{5}{3\sqrt{2}} \\ & \text{Power radiated through the beam width} \\ & P_{BW} = \frac{\Omega_{A,HP}}{\Omega_A} = \frac{\frac{5\pi\sqrt{2}}{3}}{\frac{8\pi}{3}} = \frac{5\sqrt{2}}{8} \cong 0.88 \text{ (or)} \\ & 88\% \text{ of Prad} \end{split}$$

Directivity

is the ratio of radiation intensity in a given direction to isotropic radiation intensity

$$D = \frac{U}{U_0} = \frac{4\pi U}{P_{\rm rad}}$$

If the direction is not specified \rightarrow (maximum directivity)

$$D_{\text{max}} = D_0 = \frac{U|_{\text{max}}}{U_0} = \frac{U_{\text{max}}}{U_0} = \frac{4\pi U_{\text{max}}}{P_{\text{rad}}}$$

Example 4

Example 2.5

find the maximum directivity of the antenna whose radiation intensity is $U = A_0 \sin \theta$ Write an expression for the directivity as a function of the directional angles θ and ϕ .

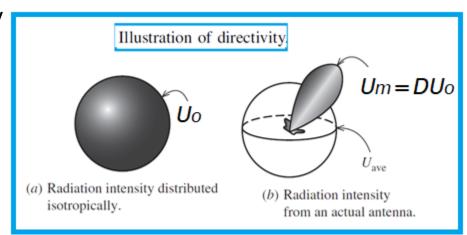
$$P_{\rm rad} = \int_0^{2\pi} \int_0^{\pi} U \sin \theta \, d\theta \, d\phi = A_0 \int_0^{2\pi} \int_0^{\pi} \sin^2 \theta \, d\theta \, d\phi = \pi^2 A_0$$

Using (2-16a), we find that the maximum directivity is equal to

$$D_0 = \frac{4\pi U_{\text{max}}}{P_{\text{rad}}} = \frac{4}{\pi} = 1.2$$

Since the radiation intensity is only a function of θ , the directivity as a function of the directional angles is represented by

$$D = D_0 \sin \theta = 1.27 \sin \theta$$



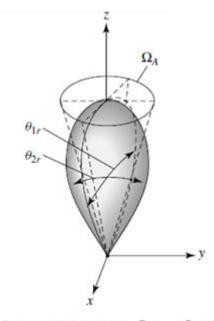
Approximated Directivity (by Kraus)

For antennas with one narrow major lobe and very negligible minor lobes, the beam solid angle is approximately equal to the product of the half-power beam widths in two perpendicular planes

For radiation intensity $U(\theta,\phi) = B_0 \cos^n(\theta)$ at $0 < \theta < \pi/2$, $0 < \phi < 2\pi$ and 0 elsewhere

$$D_0 = \frac{4\pi}{\Omega_A} \simeq \frac{4\pi}{\theta_{1r} \cdot \theta_{2r}} \quad \text{(Kraus)}$$

 θ_{1r} = half-power beamwidth in one plane (rad) θ_{2r} = half-power beamwidth in a plane at a right angle to the other (rad)



Beam solid angles $\Omega_A = \theta_{1r} \cdot \theta_{2r}$

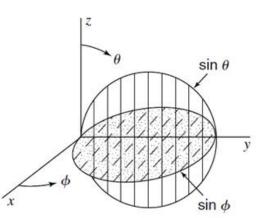
Example 5

The normalized field pattern of an antenna is given by $E_n = \sin \theta \sin \phi$, where $\theta = \text{zenith}$ angle (measured from *z* axis) and $\phi = \text{azimuth}$ angle (measured from *x* axis). E_n has a value only for $0 \le \theta \le \pi$ and $0 \le \phi \le \pi$ and is zero elsewhere (pattern is unidirectional with maximum in +*y* direction). Find (*a*) the exact directivity, (*b*) the approximate directivity and (*c*) the decibel difference.

Solution

(a)
$$G = \frac{4\pi}{\iint_{4\pi} P_n(\theta, \phi) d\Omega} = \frac{4\pi}{\iint_{0}^{\pi} \iint_{0}^{\pi} (\sin \theta \sin \phi)^2 \sin \theta d\theta d\phi}$$

 $G = \frac{4\pi}{\iint_{0}^{\pi} \iint_{0}^{\pi} \sin^3 \theta \sin^2 \phi d\theta d\phi} = \frac{4\pi}{2\pi/3} = 6$
(b) $D = \frac{41,253}{\theta_{HP}^o \phi_{HP}^o} = \frac{41,253}{(90^o)(90^o)} = 5.1$
(c) $\Delta D(dB) = 10 \log \frac{6}{5.1} = 0.7 dB$



Unidirectional $\sin \theta$ and $\sin \phi$ field patterns.

Example 6Directivity of a Sector Omnidirectional Pattern

An ideal omnidirectional antenna would have constant radiation in the horizontal plane (θ =90) and would fall rapidly to zero outside that plane. if power pattern in the vertical plane is constant out to ±30° **from horizontal**. find Directivity.

The power pattern expression is then written as:

$$F(\theta) = \begin{cases} 1 & \frac{1}{3}\pi < \theta < \frac{2}{3}\pi \\ 0 & \text{elsewhere} \end{cases}$$

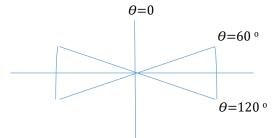
Solution:

The solid angle of the pattern is

$$\Omega_A = \int \int |F(\theta, \phi)| \, \mathrm{d}\Omega = \int_0^{2\pi} \int_{\pi/3}^{2\pi/3} \sin\theta \, d\theta \, d\phi$$
$$= (2\pi) [-\cos\theta]_{\pi/3}^{2\pi/3} = (2\pi)(0.5 + 0.5) = 2\pi$$

The directivity is

$$D = \frac{4\pi}{\Omega_A} = \frac{4\pi}{2\pi} = 2.$$



Example 7

Example 2.8

Design an antenna with omnidirectional amplitude pattern with a half-power beamwidth of 90°. Express its radiation intensity by $U = \sin^n \theta$. Determine the value of *n*, max directivity

Solution: Since the half-power beamwidth is 90°, the angle at which the half-power point occurs is $\theta = 45^{\circ}$. Thus

$$U(\theta = 45^{\circ}) = 0.5 = \sin^{n}(45^{\circ}) = (0.707)^{n}$$

or

n = 2

Therefore, the radiation intensity of the omnidirectional antenna is represented by $U = \sin^2 \theta$. An infinitesimal dipole (see Chapter 4) or a small circular loop (see Chapter 5) are two antennas which possess such a pattern.

Using the definition of (2-16a), the exact directivity is

$$U_{\text{max}} = 1$$

$$P_{\text{rad}} = \int_{0}^{2\pi} \int_{0}^{\pi} \sin^{2}\theta \, \sin\theta \, d\theta \, d\phi = \frac{8\pi}{3}$$

$$D_{0} = \frac{4\pi}{8\pi/3} = \frac{3}{2} = 1.761 \text{ dB}$$

$$\int_{0}^{\pi} \sin^{3}\theta \, d\theta = \frac{4}{3}$$

Example 8

The antenna is a lossless end-fire array of 10 isotropic point sources spaced $\lambda/4$ and operating with increased directivity. The normalized field pattern is

$$E_n = \sin\left(\frac{\pi}{2n}\right) \frac{\sin(n\psi/2)}{\sin(\psi/2)}$$

where

$$\psi = d_r (\cos \phi - 1) - \frac{\pi}{n}$$
 $d_r = \pi/2$ $n = 10$

Since the antenna is lossless, gain = directivity.

(a) Calculate the gain G.

(b) Calculate the approximate gain

(*c*) What is the difference?

Solution

$$E_{n} = \sin(\frac{\pi}{2}) \frac{\sin(25\pi (c_{0}\beta-1) - 0.5\pi)}{\sin(0.25\pi (c_{0}\beta-1) - 0.05\pi)}$$
To Find HPBW Use Table
 $\theta' = 0 + \frac{1}{10} \frac{4}{20} = 30 + 40 = 50 = 60 = 70 = 50 = 30$
 $E_{n} = \frac{1}{100} \frac{0.42}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{120}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{120}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{160}{100} \frac{120}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{170}{100} \frac{160}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{160}{100} \frac{120}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{120}{100} \frac{130}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{160}{100} \frac{120}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{160}{100} \frac{120}{100} \frac{120}{100} \frac{140}{100} \frac{150}{100} \frac{160}{100} \frac{160}{100$