



Phase Diagram of QCD



Kenji Fukushima

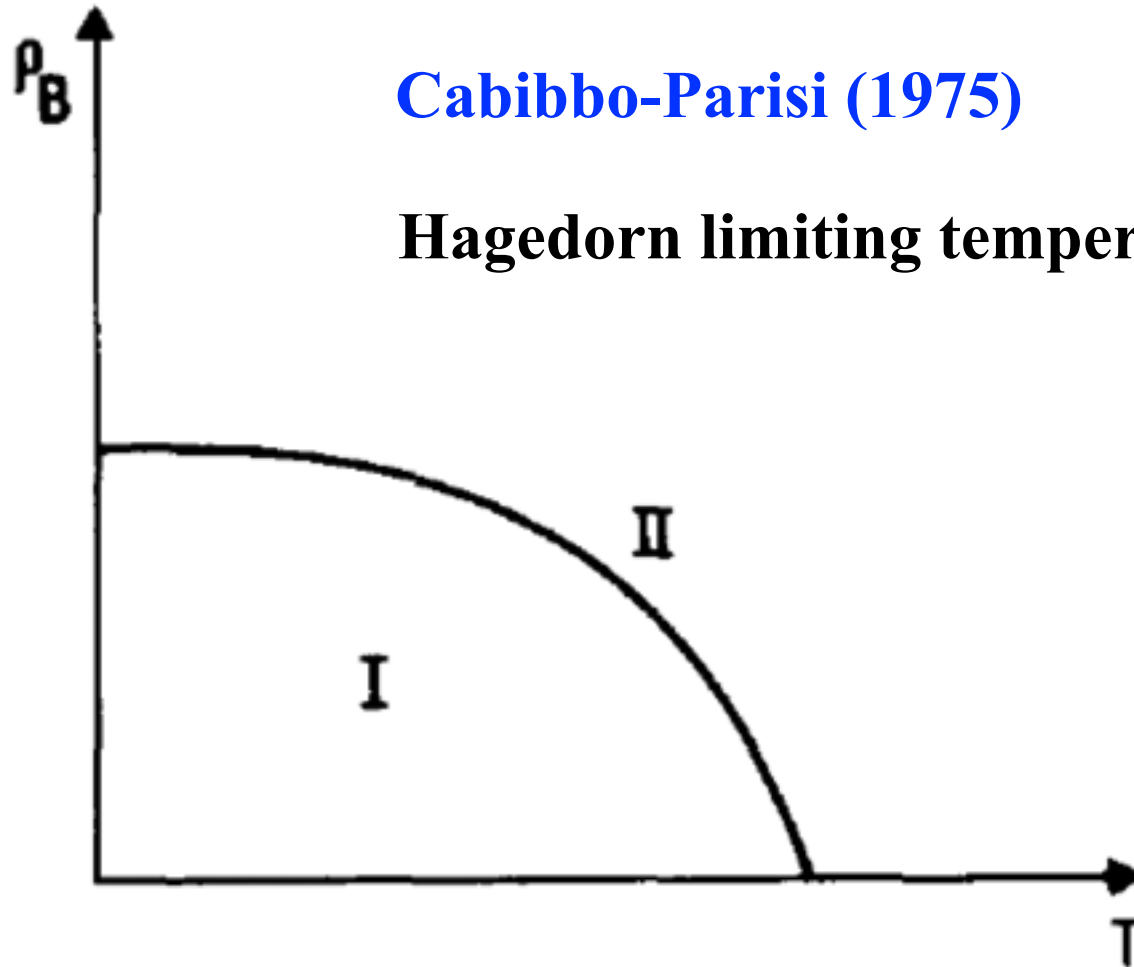
The University of Tokyo

— QCD Phase Structure III in CCNU Wuhan —

A bit about the History

as an introduction to diquarks

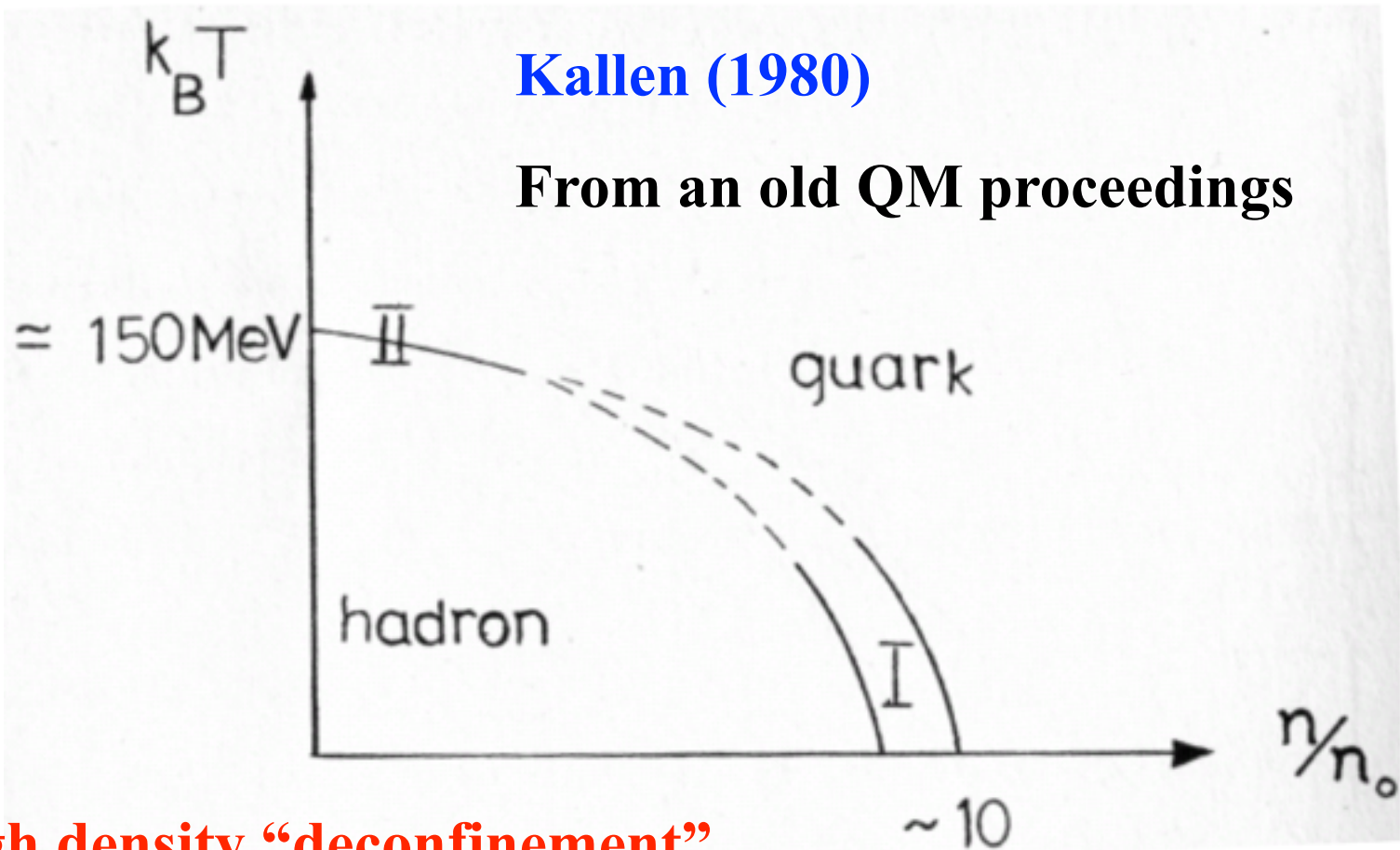
Birth of the QCD Phase Diagram



Cabibbo-Parisi (1975)

Hagedorn limiting temperature

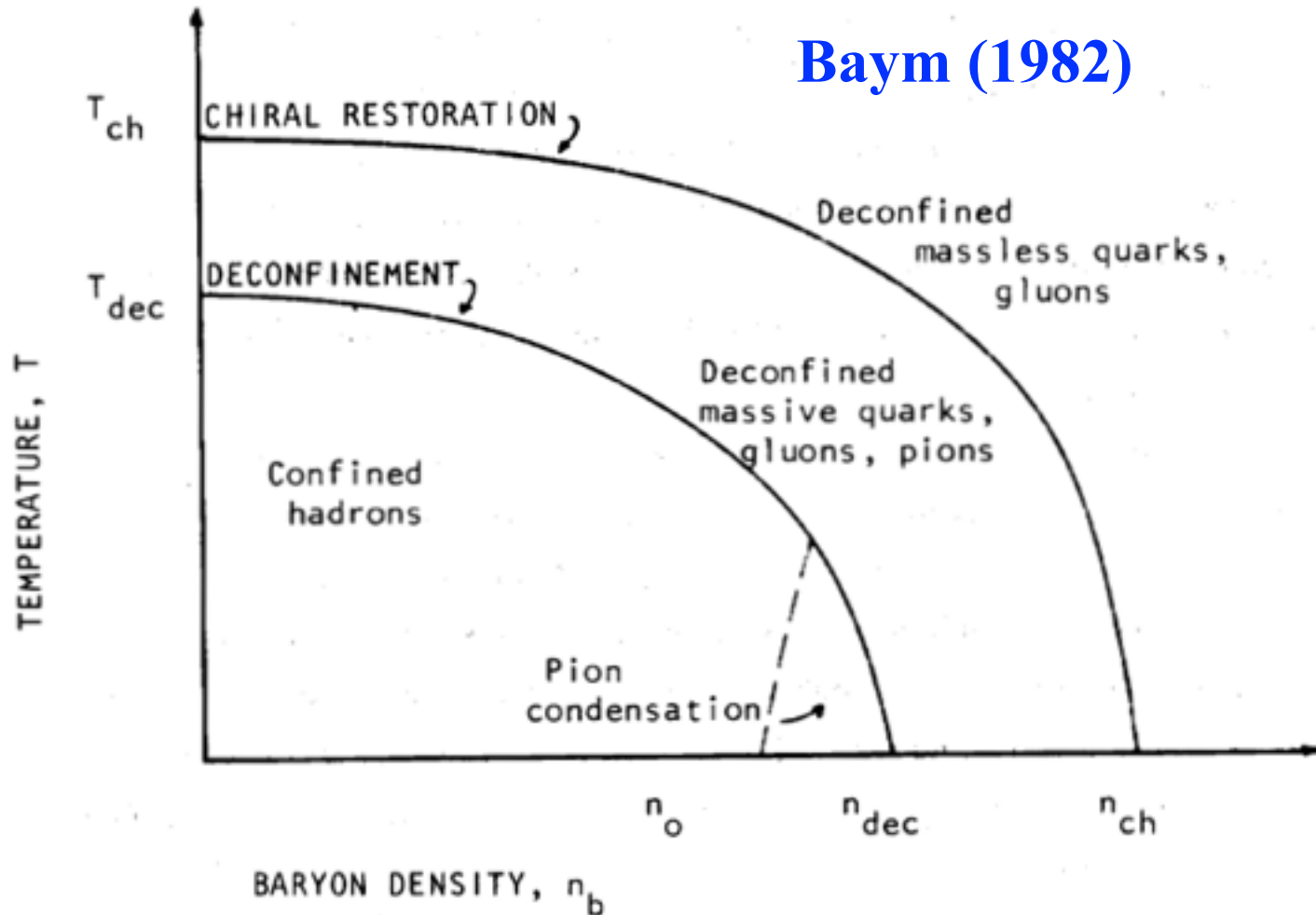
Forgotten Candidate



High density “deconfinement”

Forgotten Candidate

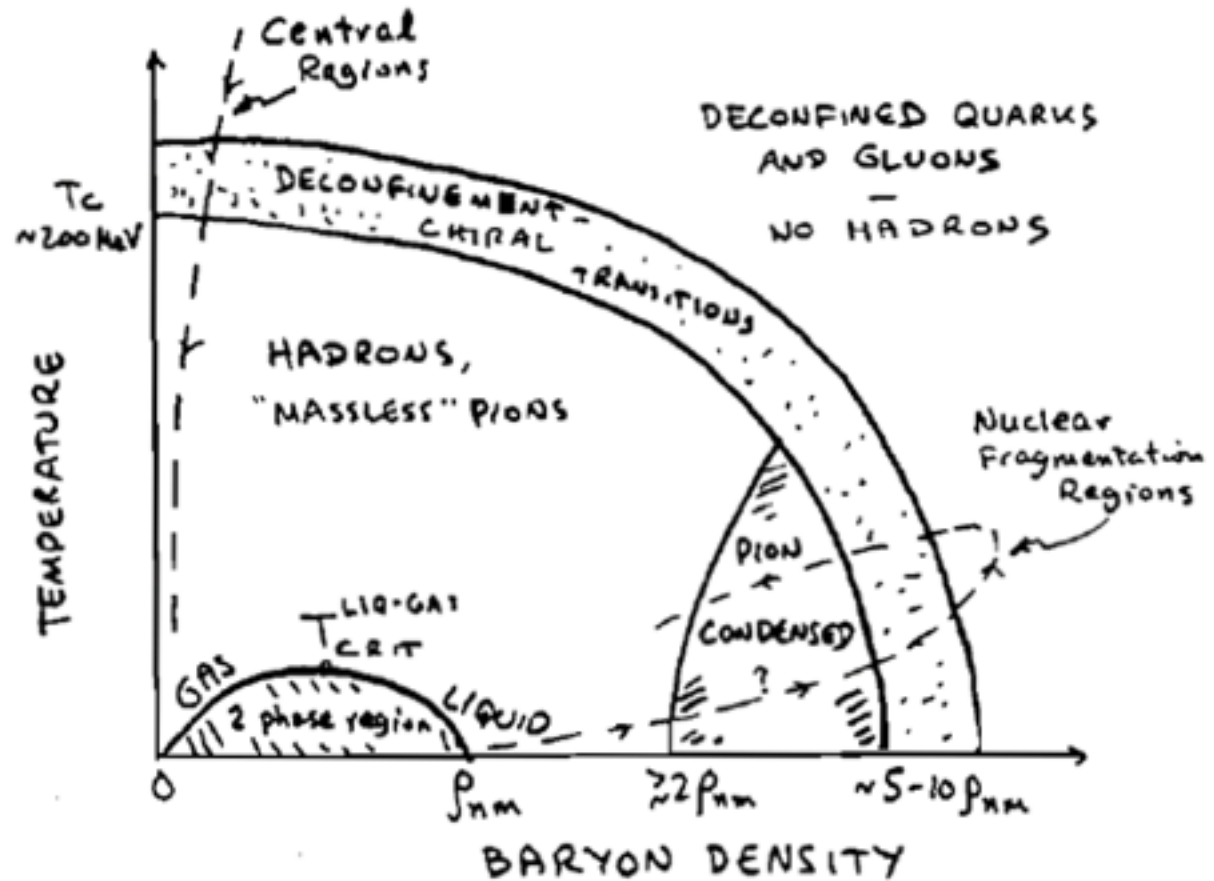
Baym (1982)



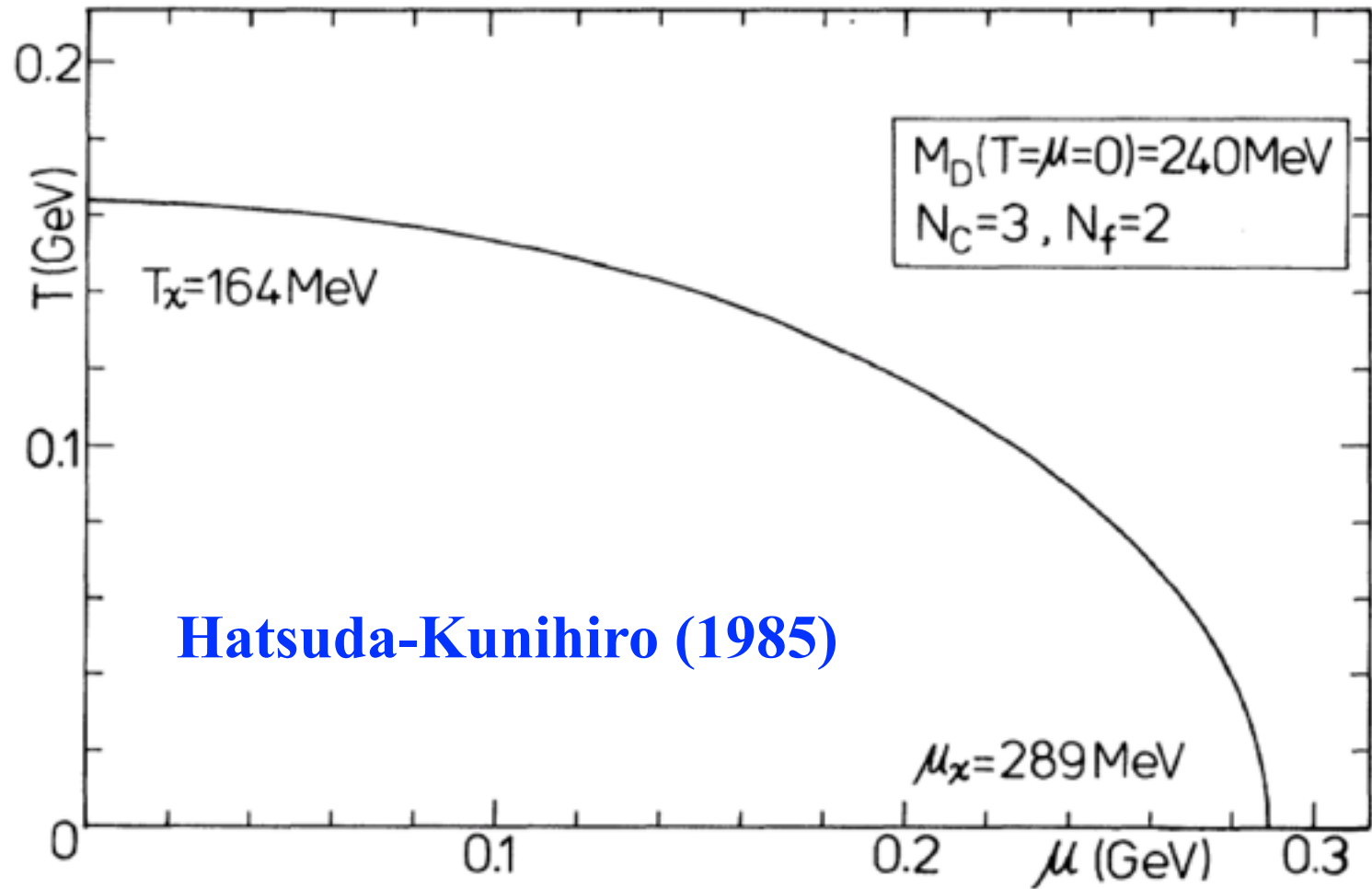
Prototype of Phase Diagram

Baym (1983)

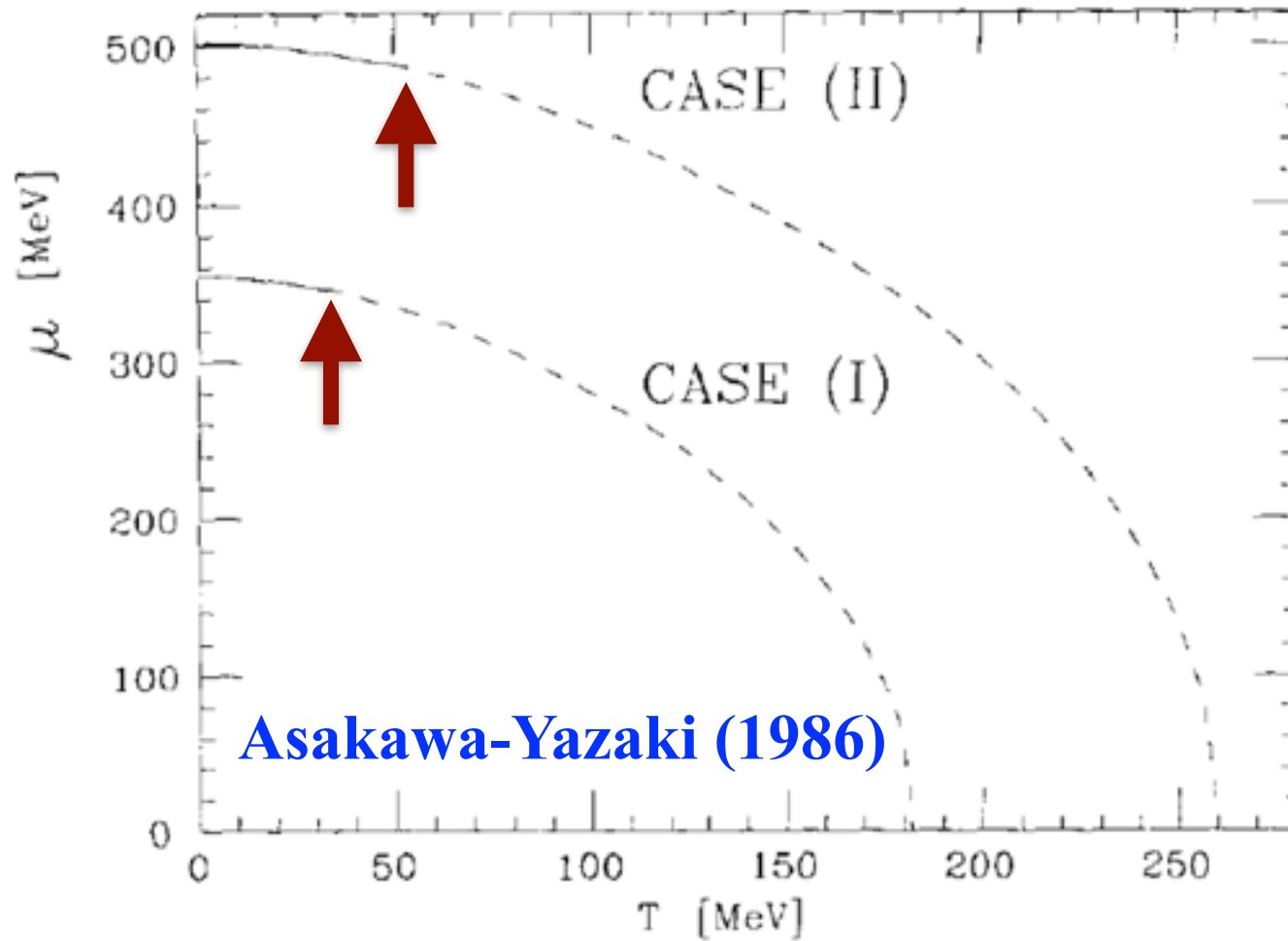
PHASE DIAGRAM OF NUCLEAR MATTER



Chiral Phase Diagram

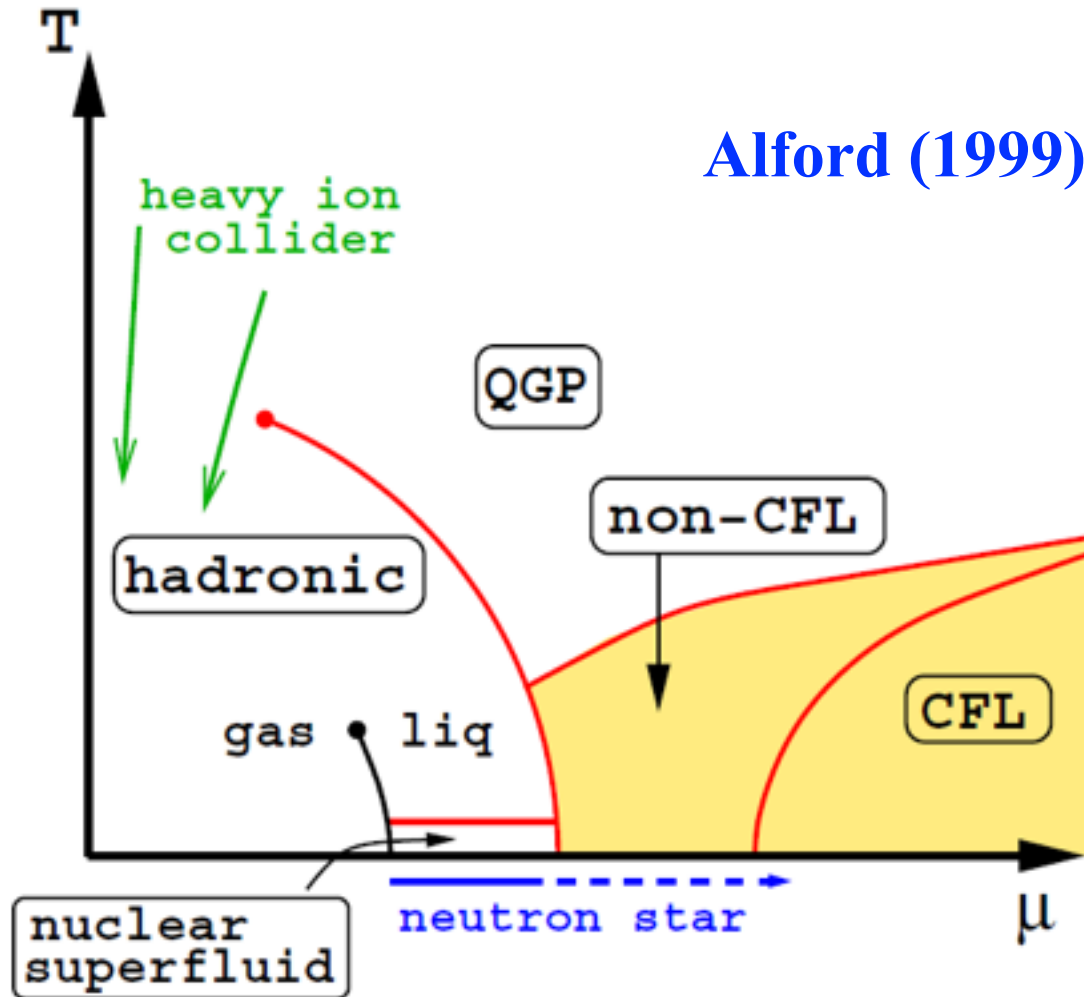


QCD Critical Point

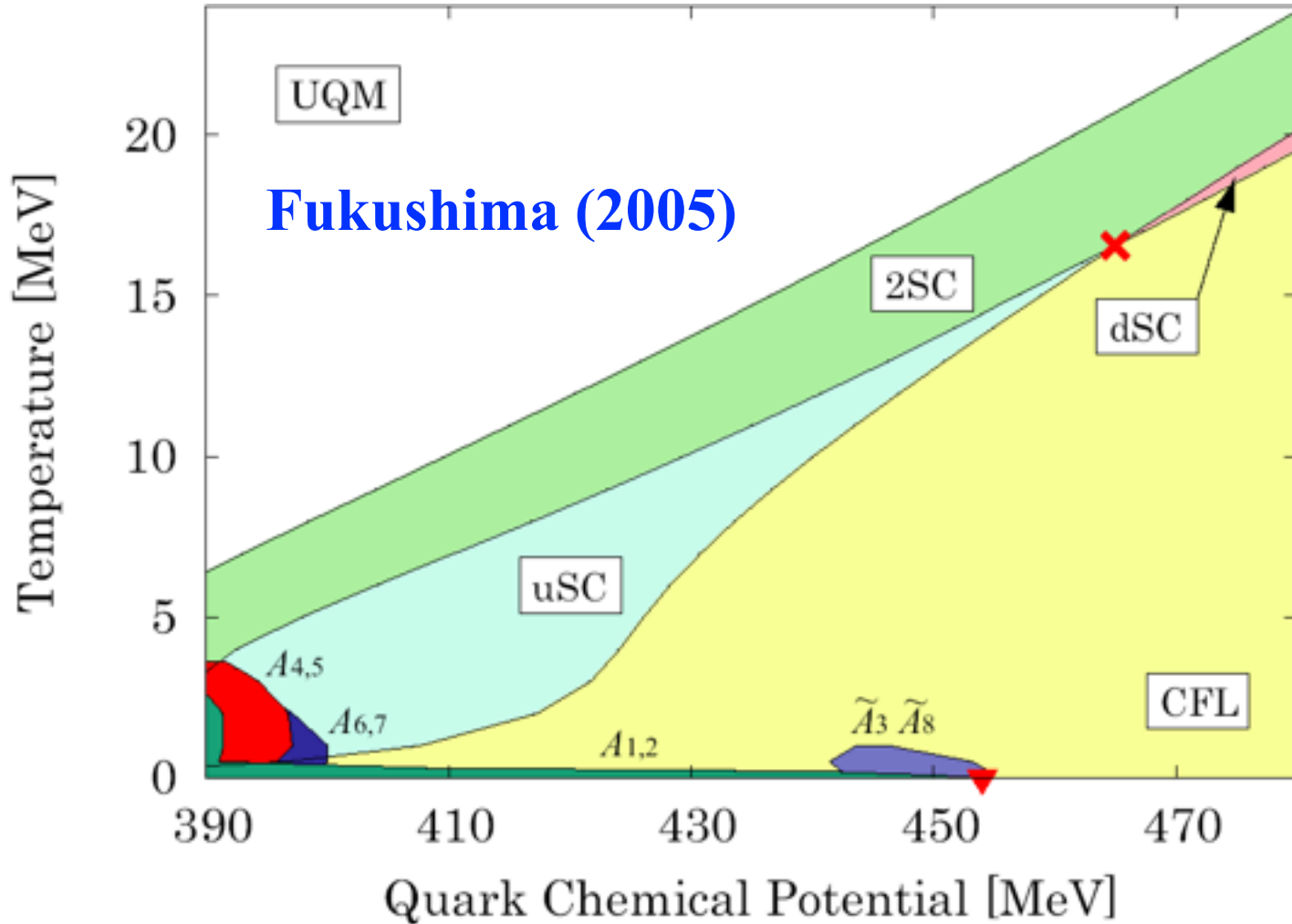


Asakawa-Yazaki (1986)

Color Superconductivity



Chromomagnetic Instabilities



Some of Highlights



“Charge Neutrality Effects on 2-flavor Color Superconductivity”

Mei Huang, **Pengfei Zhuang**, Weiqin Chao: hep-ph/0207008

“Breached Pairing Superfluidity at Finite Temperature and Density”

Jinfeng Liao, **Pengfei Zhuang**: cond-mat/0307516

“Pion Condensation in Baryonic Matter: from Sarma Phase to Larkin-Ovchinnikov-Fudde (Fulde)-Ferrell Phase”

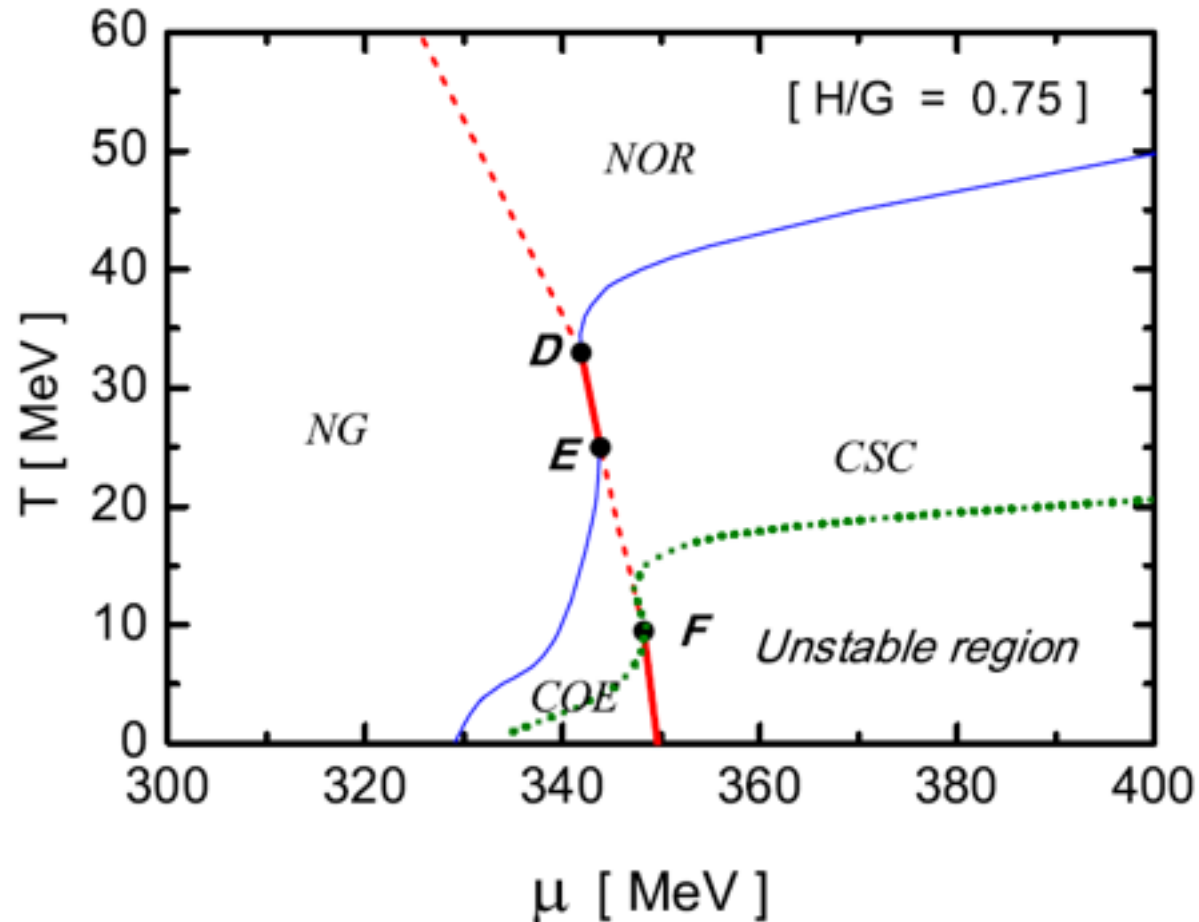
Lianyi He, Meng Jin, **Pengfei Zhuang**: hep-ph/0604224

“Neutral Color Superconductivity Including Inhomogeneous Phases at Finite Temperature”

Lianyi He, Meng Jin, **Pengfei Zhuang**: hep-ph/0610121

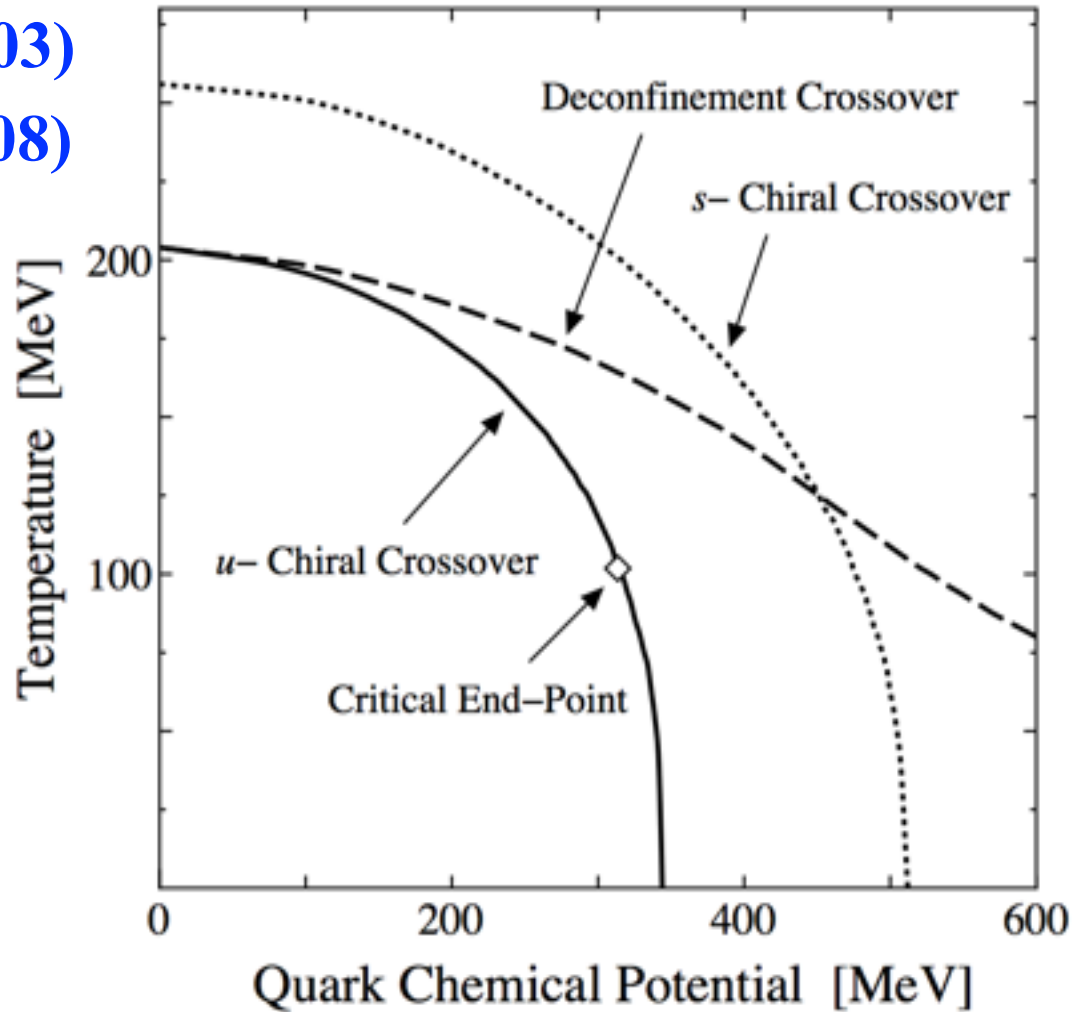
Many QCD Critical Points?

Zhang-Fukushima-Kunihiro (2008)

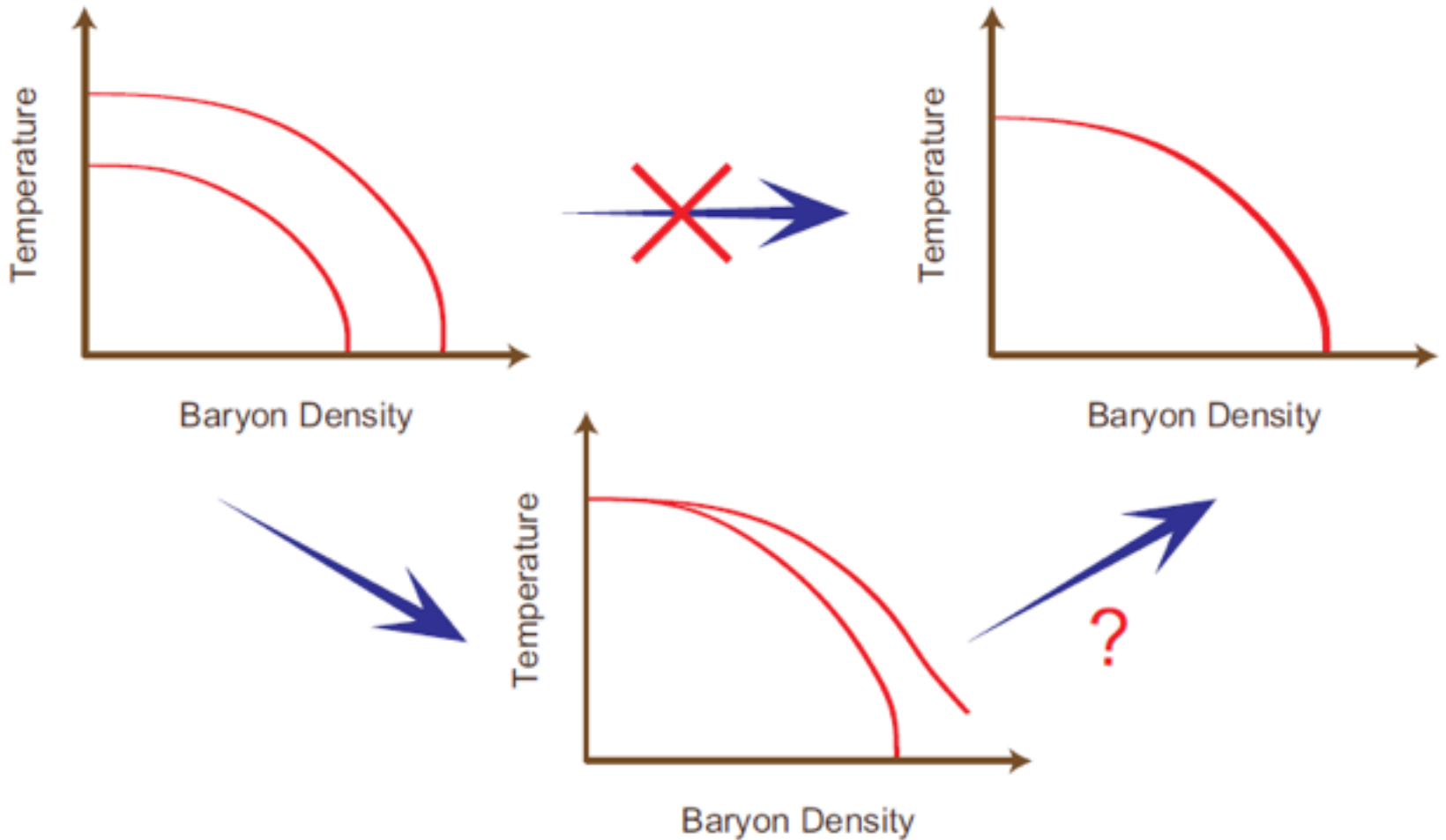


Polyakov Loop

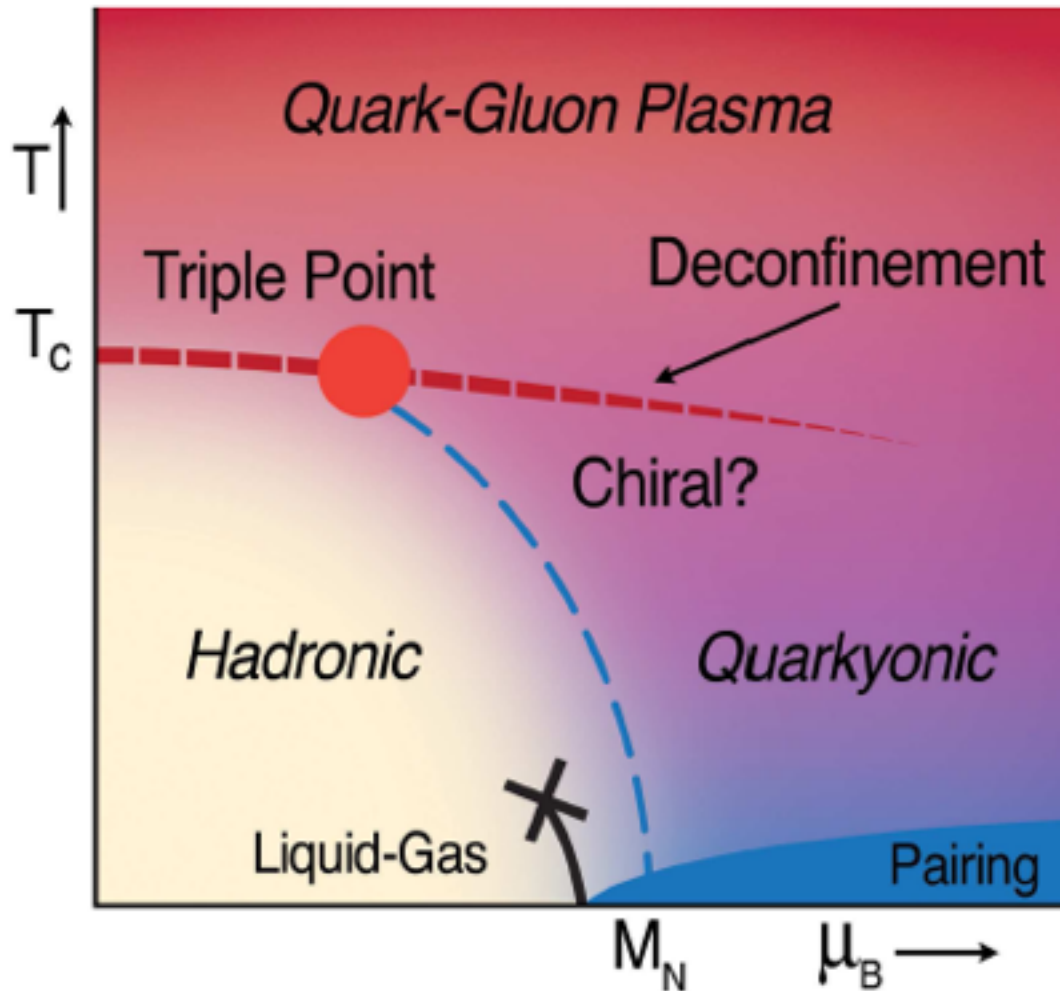
Fukushima (2003)
(2008)



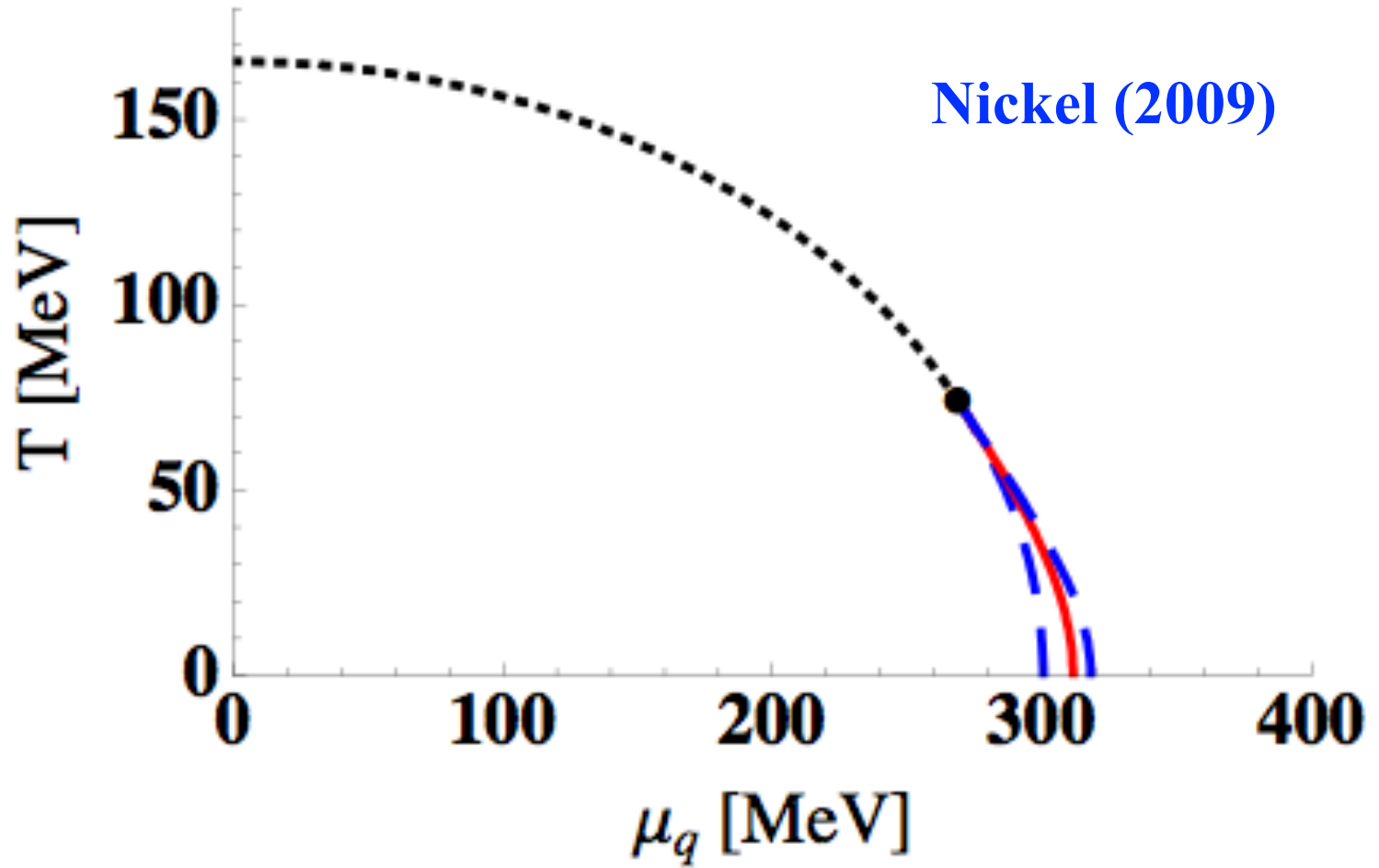
Revival of Interest



BNL-Kyoto-... Diagram

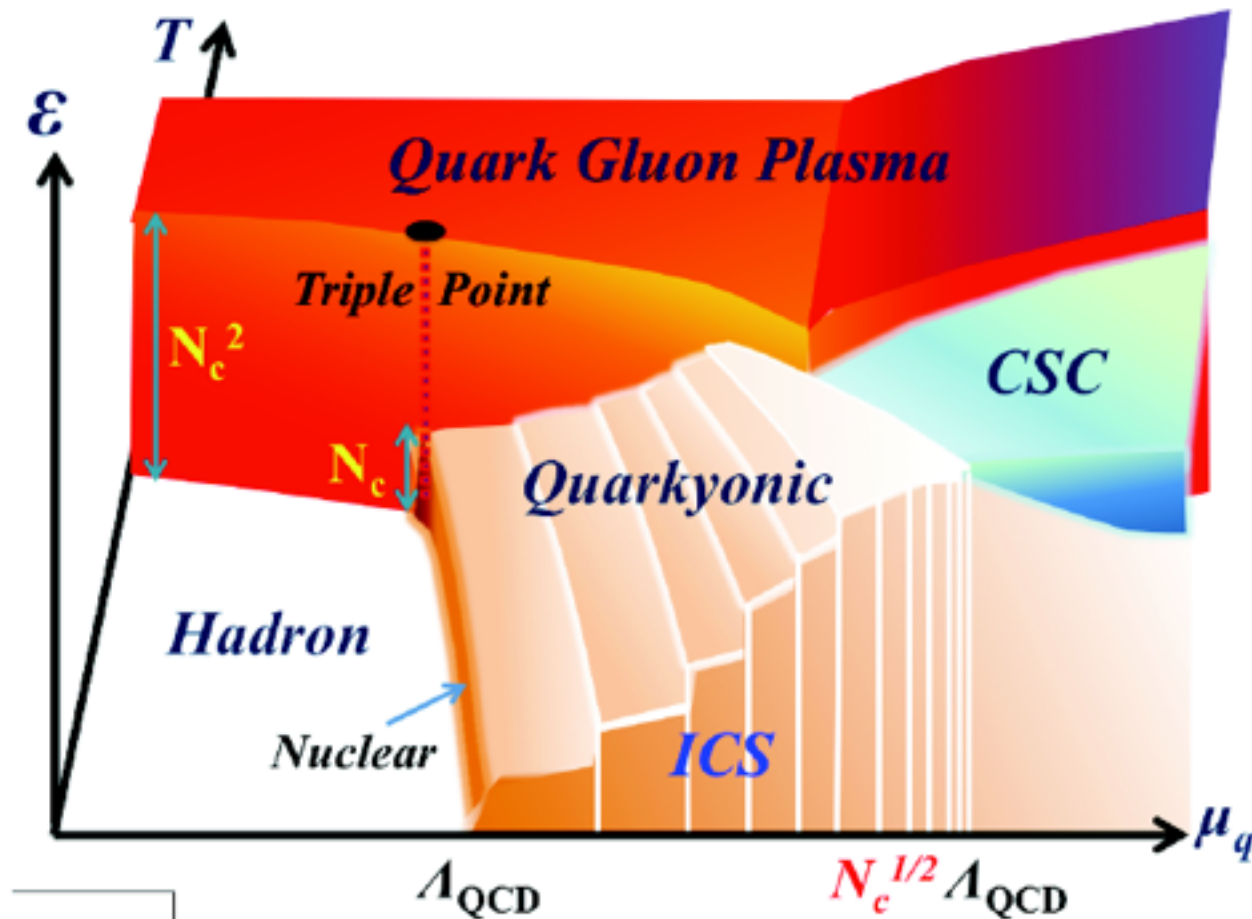


Inhomogeneity

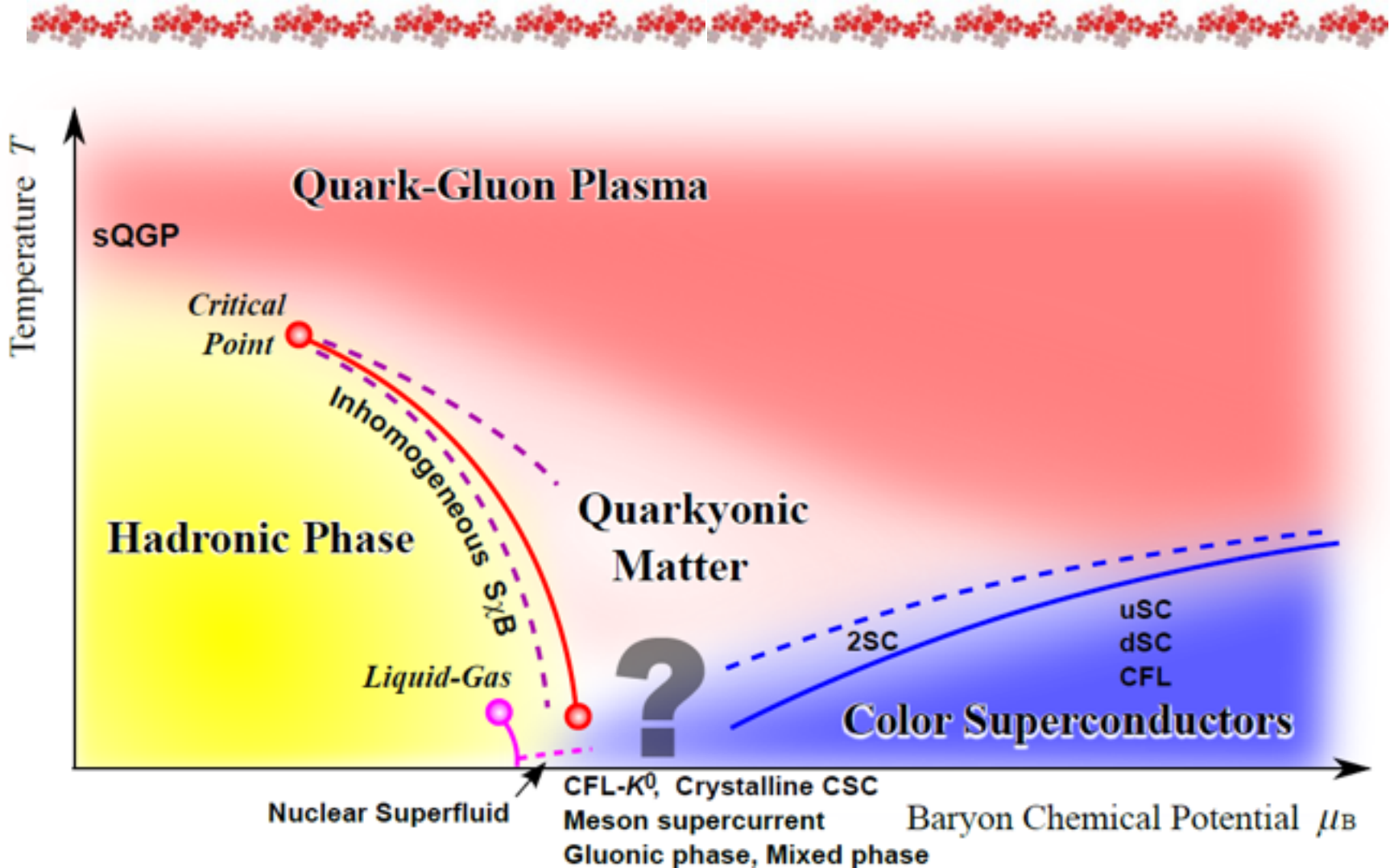


Quarkyonic Chiral Spirals

Kojo-Hidaka-Fukushima-McLerran-Pisarsky (2010)

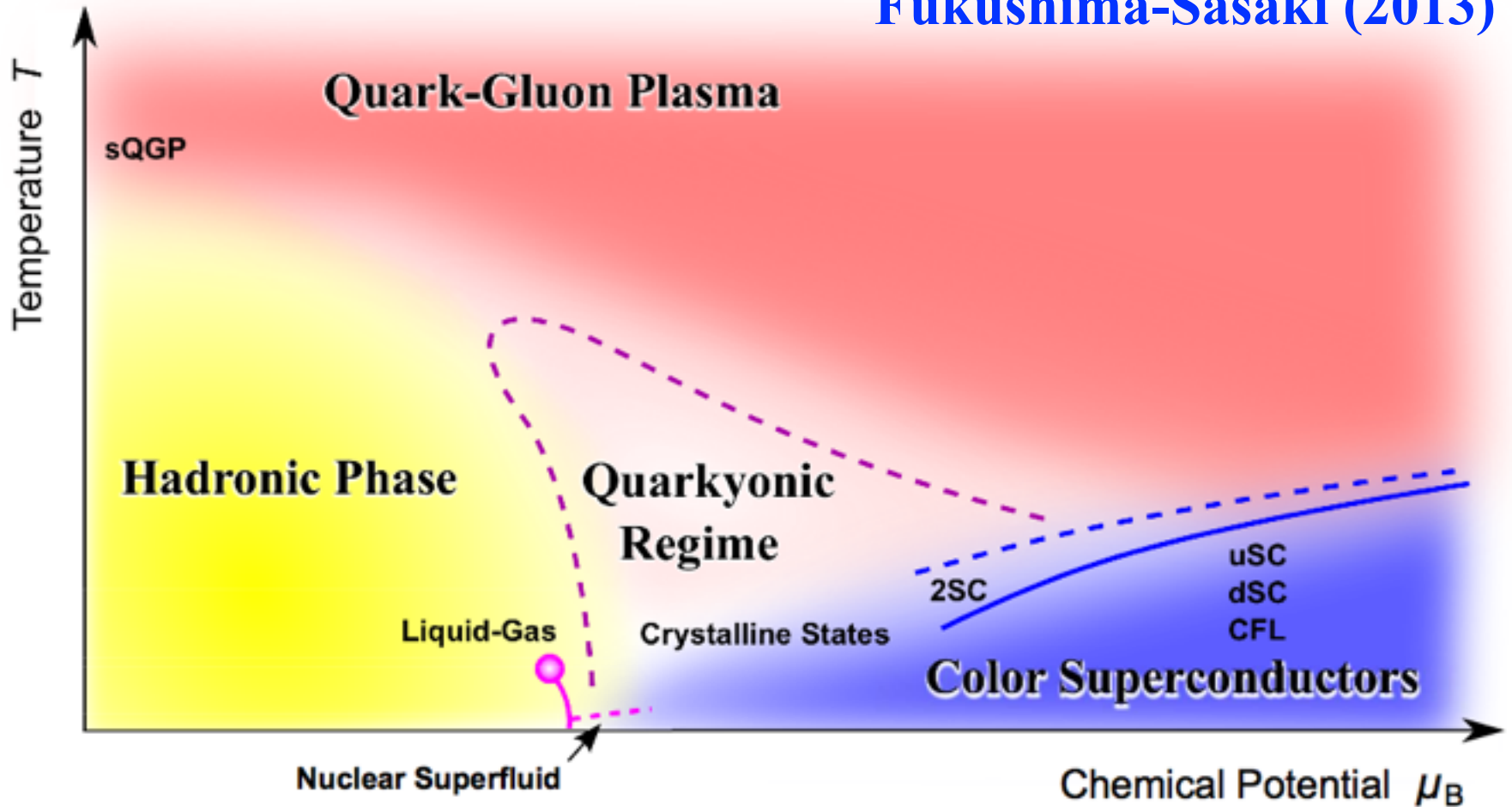


Summary



Alternative Summary

Fukushima-Sasaki (2013)



Interesting Topics Ongoing

■ Finite isospin / chiral / chemical potentials

- First-order phase transition? Confirmable with lattice
- Pion superfluid / LOFF (*Ask Pengfei for details*)

■ Finite B (not E) with/without T and density

- (Inverse) Magnetic catalysis
- CP in anisotropic pure Yang-Mills theory

■ Finite rotation (angular momentum)

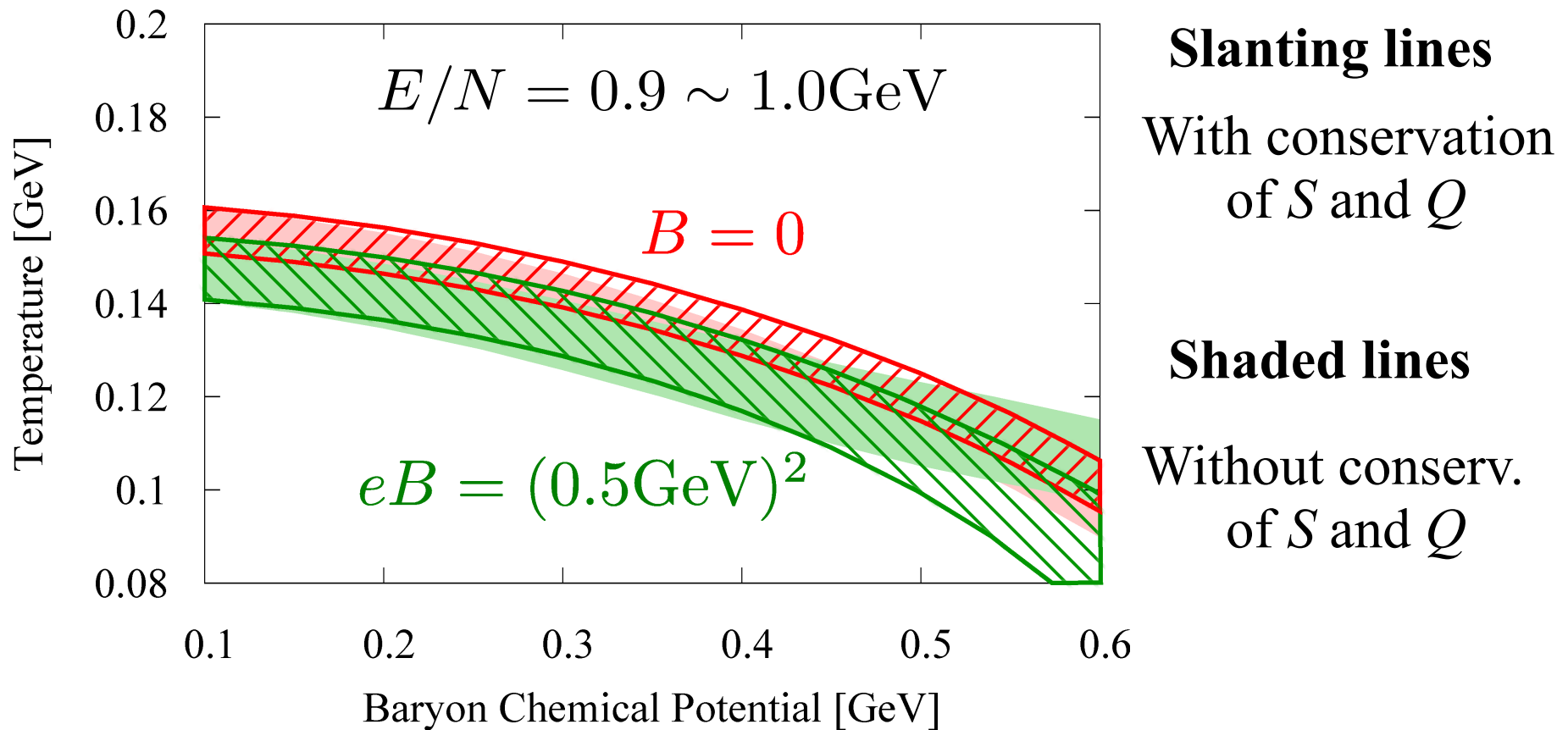
- Similar to B / similar to finite density
- Topological currents

*Xu-guang
Jinfeng...*

■ Finite curvature (curved space-time)

- Chiral symmetric mass gap / Early Universe

Magnetic Shift of Chemical Freezeout



KF-Hidaka (2016)

More about the Diquark

Never ending project with Jan...

Biggest Question Mark???



**Nuclear
Matter**



**Quark
Matter**

Diquarks

Bare vs Constituent



$$\text{Meson} \sim q\bar{q} + q\bar{q}q\bar{q} + q\bar{q}q\bar{q}q\bar{q} + \dots$$

(Vacuum Re-organized)

$$\sim q_{\text{con}}\bar{q}_{\text{con}} + (\text{Bag Constant})$$

How can we be so sure about

$$M \sim q\bar{q}$$
$$B \sim qqq$$

beyond quantum num of Quark Model ?

A SCHEMATIC MODEL OF BARYONS AND MESONS *

M. GELL-MANN

California Institute of Technology, Pasadena, California

Received 4 January 1964

6) James Joyce, *Finnegan's Wake* (Viking Press, New York, 1939) p.383.

Primeval expression of diquarks

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" 6) q and the members of the anti-triplet as anti-quarks \bar{q} . Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.

Who was the First?



Perhaps these missing states really do not exist. If baryons were diquark–quark systems, Fig. 2, as Lichtenberg and Tassie noted more than 40 years ago [4], the number of states would be restricted and in fact be very like that currently observed.

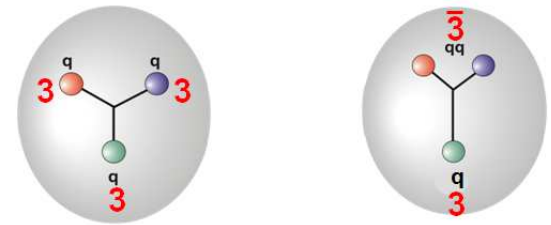
Pennington (2011)

Prog. Theor. Phys. Vol. 36 (1966), No. 4

Baryon Resonances in a Quark Model

Masakuni IDA and Reido KOBAYASHI

We suppose that baryons consist of a qq pair (or a diquark) and another quark moving around it with orbital angular momentum L . In order that for $L=0$ our model can produce the $1/2^+$ octet and the $3/2^+$ decuplet, which belong to the “56” of $SU(6)$, the qq pair must be in a 3S_1 state and form an $SU(3)$ sextet. Unwanted



If you know, let me know!

Exotica (scalar nonet)

Volume 60B, number 2

PHYSICS LETTERS

5 January 1976

UNCONVENTIONAL STATES OF CONFINED QUARKS AND GLUONS[☆]

R.L. JAFFE^{*} and K. JOHNSON

*Laboratory for Nuclear Science and Department of Physics,
Massachusetts Institute of Technology, Cambridge, Mass. 02139, USA*

mentally observed “resonances”. In particular we are led to classify the 0^{++} enhancements known as the ϵ , S^* and δ as $QQ\bar{Q}\bar{Q}$ states. If correct, this assignment disrupts further the already uneasy state of the P-wave mesons in the quark model.

broad exotic $QQ\bar{Q}\bar{Q}$ states and the P-wave baryons states overlap broad $4Q\bar{Q}$ states. In such cases one might expect that mixing effects will play an essential role in an unravelling of partial widths. This may provide a clue to an understanding of some of the elusive P-wave states such as the A_1 .

**Can explain why $a_0(980)$ heaviest without strangeness
Exotic component is to be mixed (via instanton int.)**

Classification

$$\mathbf{3} \otimes \mathbf{3} = \bar{\mathbf{3}} \oplus \mathbf{6}$$

Quantum numbers and operators

J^P	Color	Flavor	Operator
0^+	$\bar{\mathbf{3}}$	$\bar{\mathbf{3}}$	$\bar{\psi}_C \gamma_5 \psi$, $\bar{\psi}_C \gamma_0 \gamma_5 \psi$
1^+	$\bar{\mathbf{3}}$	$\mathbf{6}$	$\bar{\psi}_C \gamma_i \psi$, $\bar{\psi}_C \sigma_{0i} \psi$
0^-	$\bar{\mathbf{3}}$	$\mathbf{6}$	$\bar{\psi}_C \psi$, $\bar{\psi}_C \gamma_0 \psi$
1^-	$\bar{\mathbf{3}}$	$\bar{\mathbf{3}}$	$\bar{\psi}_C \gamma_i \gamma_5 \psi$, $\bar{\psi}_C \sigma_{ij} \psi$

Diquark Phenomenology

$\Delta I=1/2$ rule in non-leptonic weak decay

$$\Delta I = 1/2 \gg \Delta I = 3/2$$
$$\times \sim 20$$

Stech, Neubert, Xu, Dosch (1987~)

Fierz transformed interaction:

$$V_{\text{eff}} = \frac{G_F}{\sqrt{2}} V_{ud} V_{us} \left[c_- (ud)_{\bar{3}}^\dagger (su)_{\bar{3}} + c_+ (ud)_6^\dagger (su)_6 + \dots \right]$$

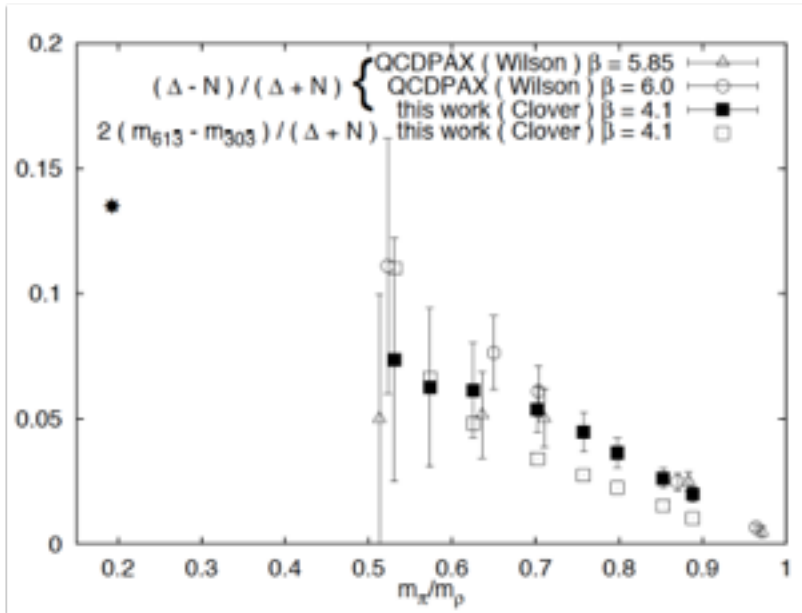
$\Delta I = 1/2$ $\Delta I = 1/2, 3/2$

Enhanced by diquarks

Diquark on Quenched Lattice

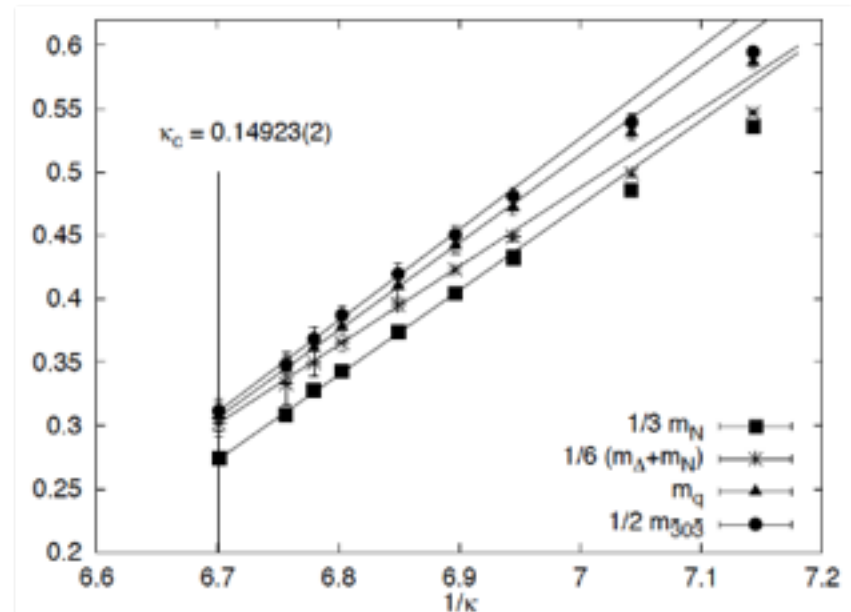


Lattice in Landau gauge (Hess-Karsch-Laermann-Wetzorke 1998)



Mass splitting from $S=1$ and $S=0$

$$m_{\text{bad}} - m_{\text{good}} \approx \frac{1}{2} (m_{\Delta} - m_N)$$



$$m_{\text{good}} = 694(22) \text{ MeV}$$

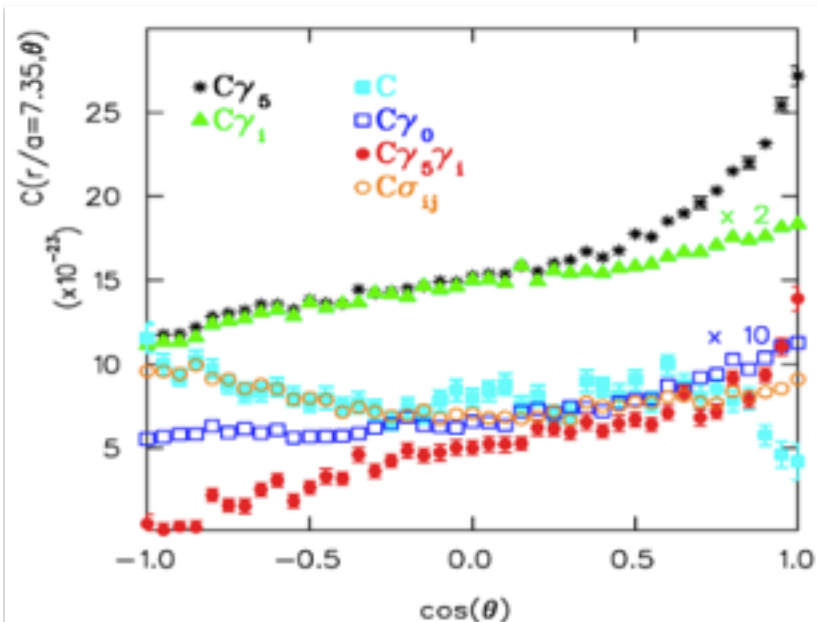
$$(m_{\pi} \simeq 350 \text{ MeV})$$

Diquark on Lattice

Density-density correlator (Alexandrou-de Forcrand-Lucini 2005)

Idealized : static-light-light baryon system

$$C(\mathbf{r}_u, \mathbf{r}_d) = \langle N | \rho^u(\mathbf{r}_u) \rho^d(\mathbf{r}_d) | N \rangle$$



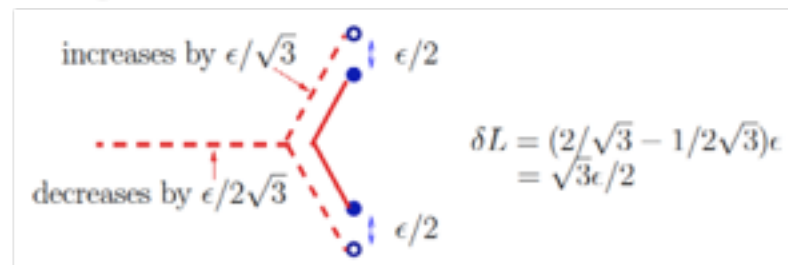
Lucini et al. (2006)

Correlation $0^+ > 1^+ \gg 0^-$

Characteristic diquark size

$\sim 1.1 \pm 0.2$ fm Leinweber (1993)

Larger than meson size



Diquark and Deconfinement

Deconfinement in pQCD



pQCD justifies itself at high T

All gluons are screened by gT or g^2T

pQCD does not justify itself at high μ

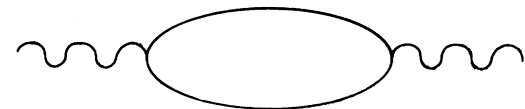
Magnetic gluons never screened

Superdense Matter: Neutrons or Asymptotically Free Quarks?

J. C. Collins and M. J. Perry

*Department of Applied Mathematics and Theoretical Physics, University of Cambridge,
Cambridge CB3 9EW, England*

(Received 6 January 1975)



Insufficient justification

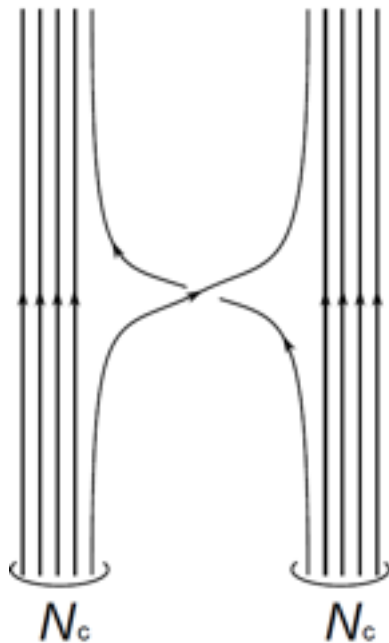
Color super justifies pQCD at high μ

All gluons are screened by $g\mu$

No Deconfinement?

Large- N_c Baryonic (Quark) Matter

$$\mathcal{O}(1) \longrightarrow \mathcal{O}(N_c)$$



$$N_c \Lambda_{\text{QCD}}^4$$

Quarkyonic Matter

McLerran, Pisarski
Hidaka, Kojo
Fukushima, Sasaki

Deconfinement Revisited



CFL (3-flavor CSC) reads:

SU(3)_c broken completely

All 8 gluons get massive (Meissner effect)

No confinement remains

Can this be a “definition” of deconfinement ?

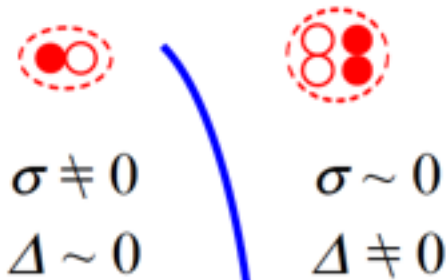
Private communications with Gordon Baym

If so, quark matter is realized only through diquarks

Continuity from NM to QM

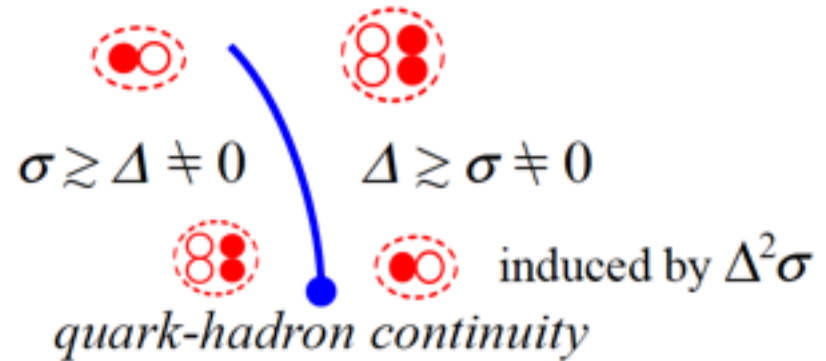


U(1)_A Symmetric

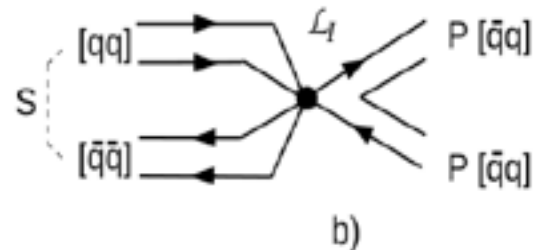
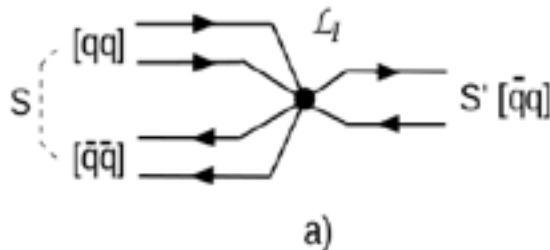


μ_q \longrightarrow

U(1)_A Broken



$$\pi^+ \sim (u)(\bar{d}) \sim (\bar{d}\bar{s})(su)$$

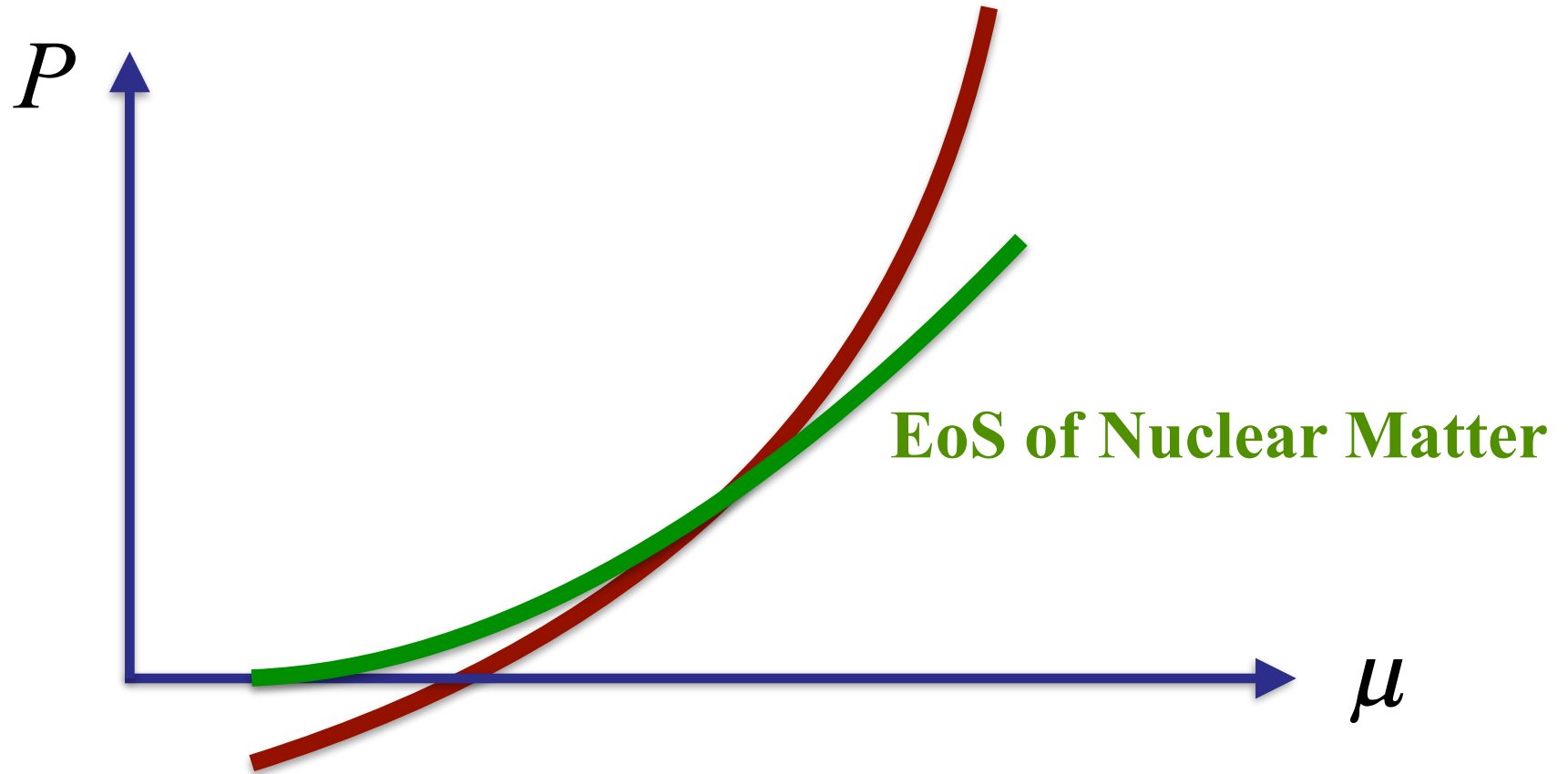


('t Hooft-Isidori-Maiani-Polosa-Riquer 2008)

Diquark Continuity



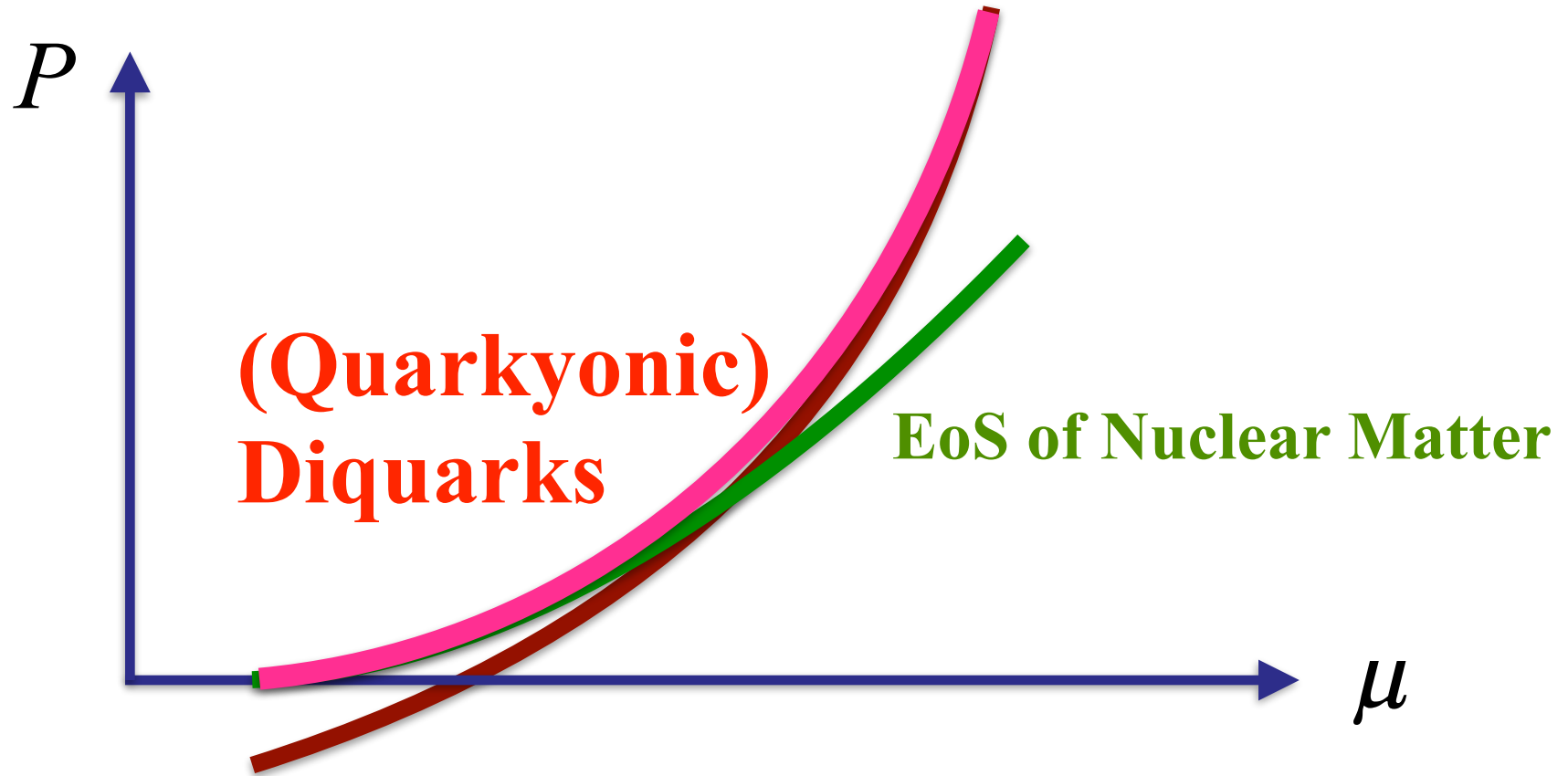
EoS of Quark Matter



Diquark Continuity



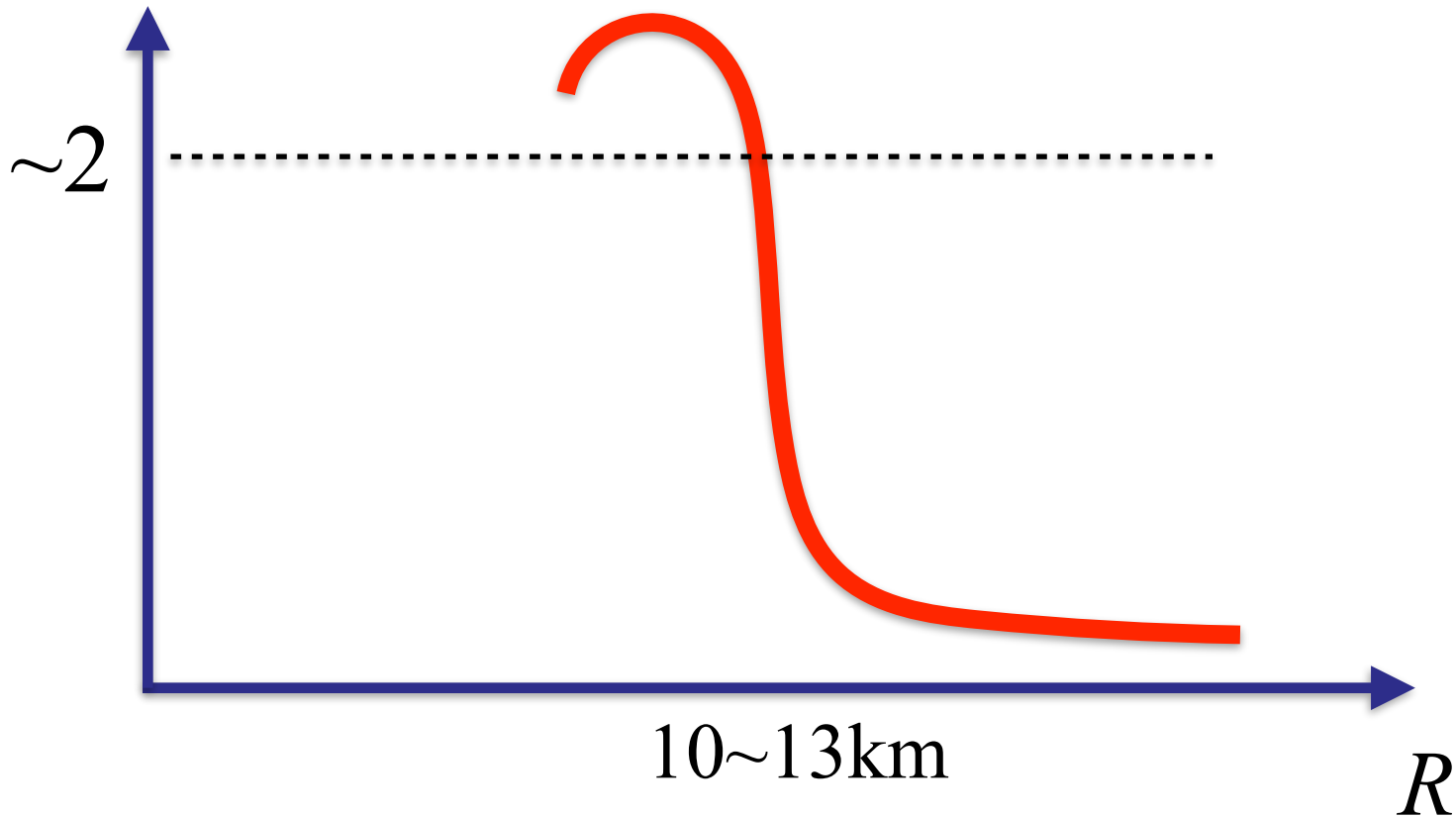
EoS of Quark Matter



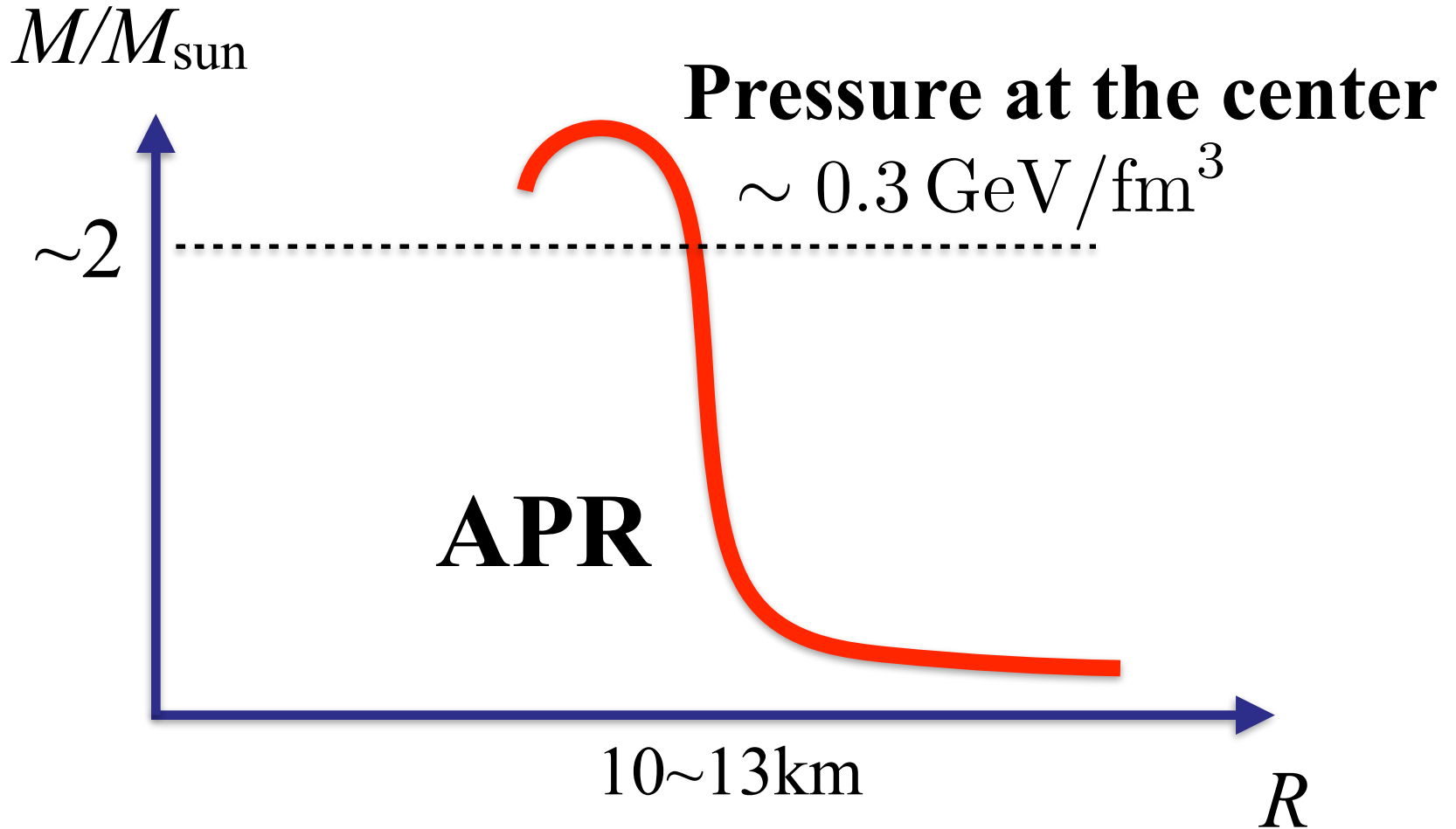
Implication to the M - R Relation



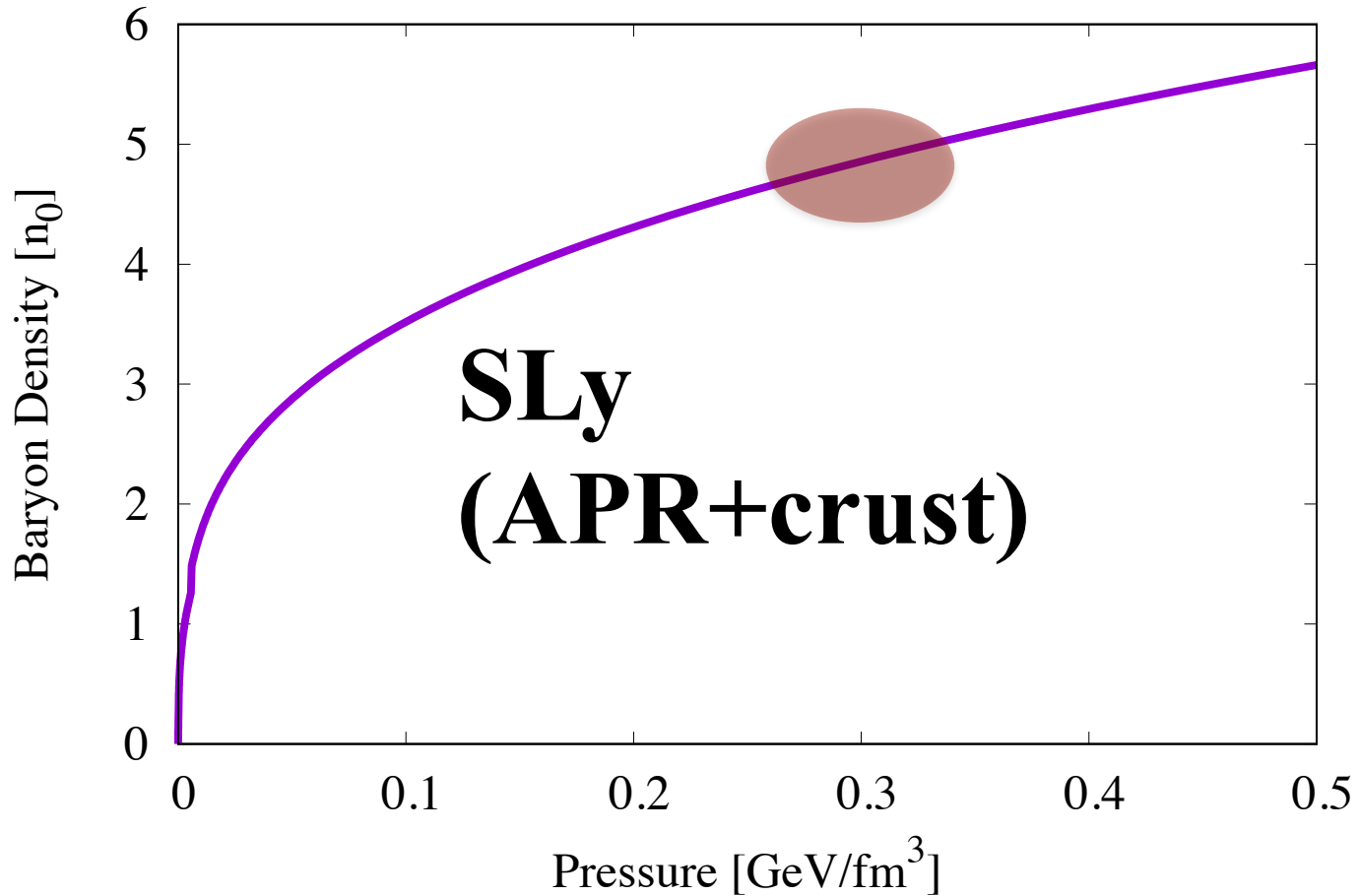
M/M_{sun}



Implication to the M - R Relation



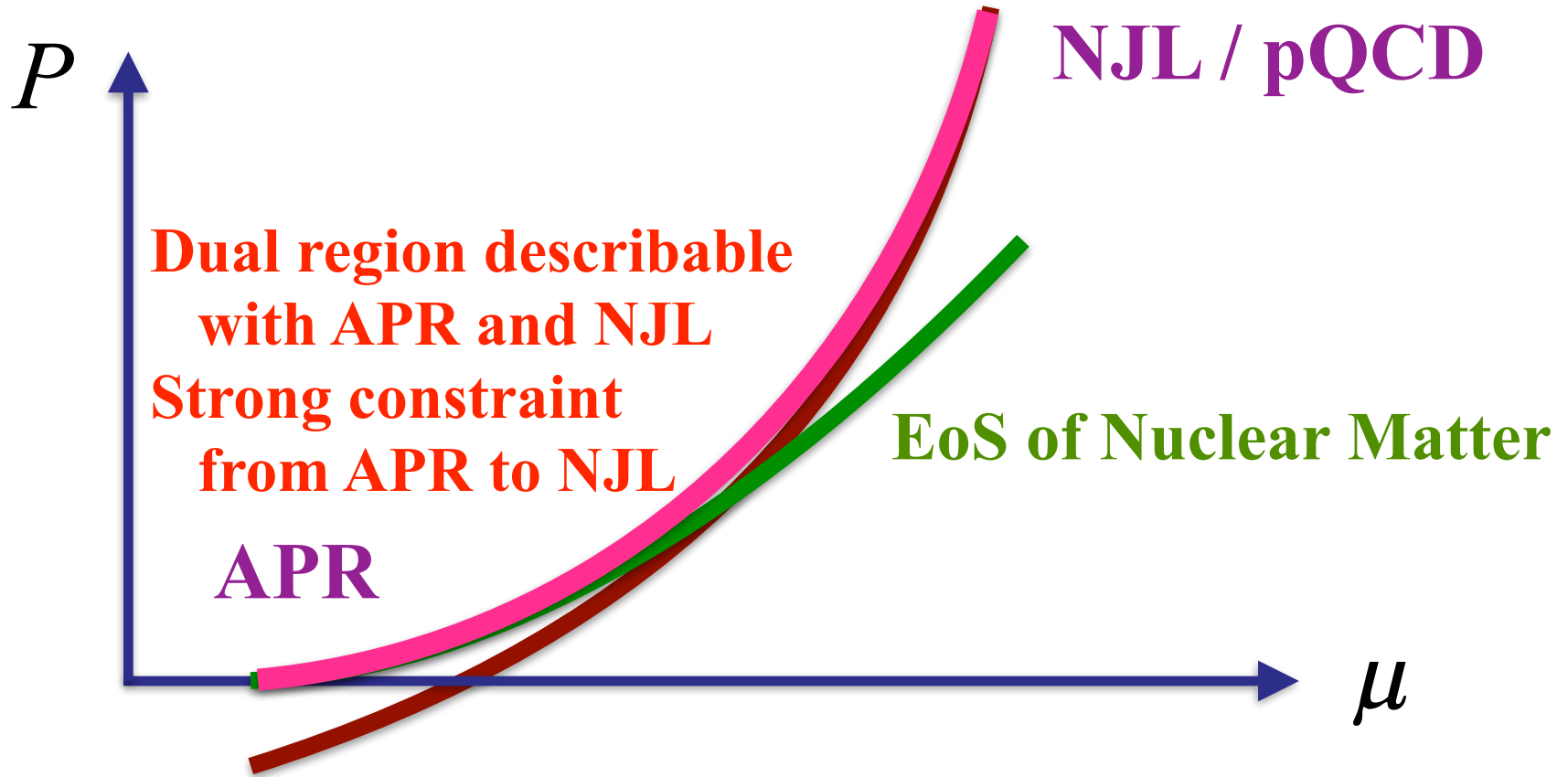
Cannot be right!?



Duality Hypothesis



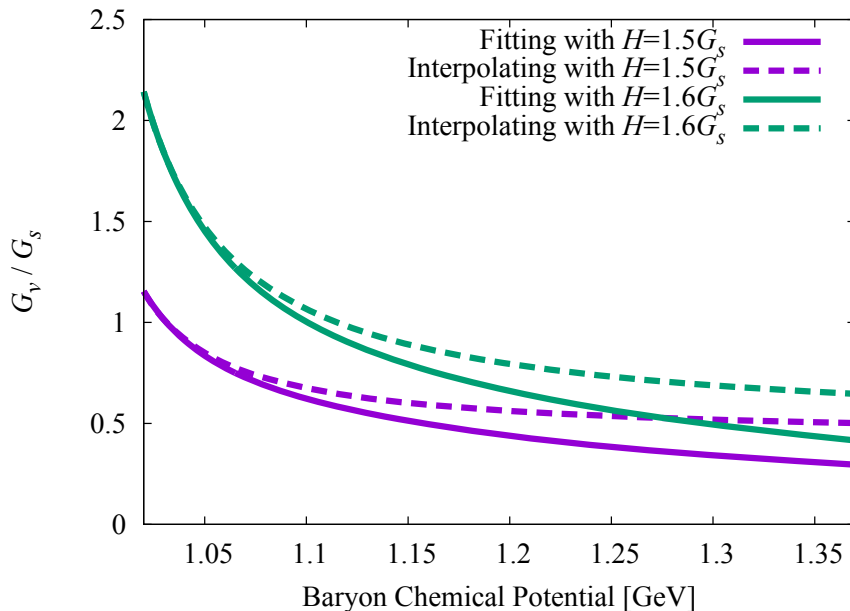
EoS of Quark Matter



APR-constrained NJL

All non-perturbative effects renormalized in G_V

APR can be reproduced with “running” vector interaction



Best fit function \sim inverse log

$$\text{cf. } \alpha_s(\mu) = \frac{1}{b \log(\mu/\Lambda)}$$

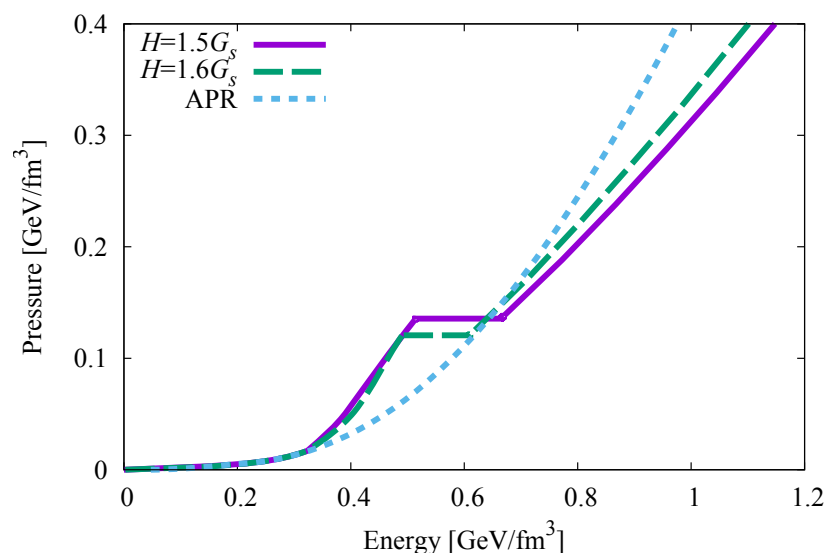
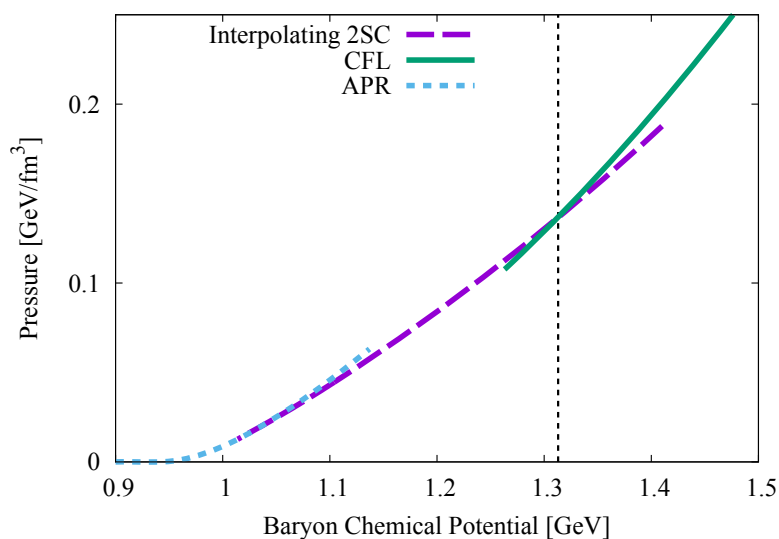
Suggestive!!!

Nuclear matter knows the running coupling?

APR-constrained NJL



Fukushima-Kojo (2015)

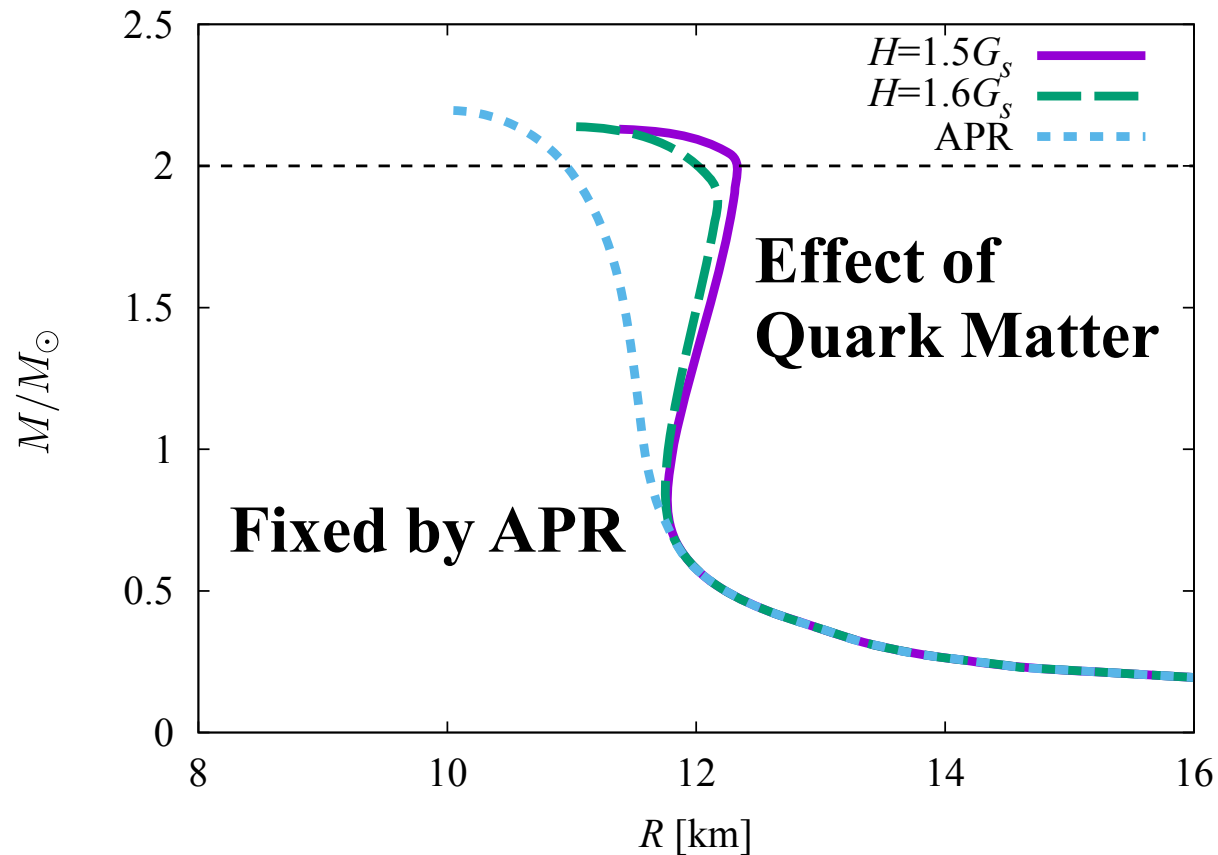


Weak 1st-order Phase Transition (2SC-CFL)

Single unified theory covering all the densities!

M-R Relation

Fukushima-Kojo (2015)



Summary and Speculation

Summary

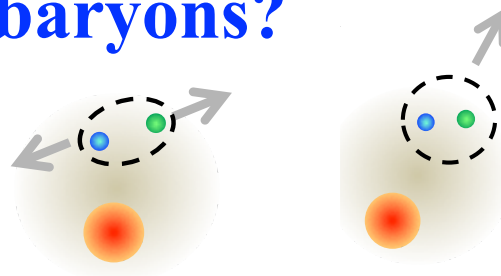
■ Nuclear Matter = 2SC + Chiral Cond. + 6-diquark

- More reasonable than CFL-NM continuity
- Chance to access the diquark superfluid phase in HIC?
- Enhanced fluctuations from (critical) diquarks
- Refined HRG with diquarks?

■ New Model = APR-constrained NJL

- Microscopic information superseding parametrization
- Less crazier than using APR to $\sim 5 n_0$! (Ask Toru!)

■ Spectroscopy of Qqq baryons?



Provocative Speculation

