

A photograph of a laser laboratory setup. The scene is dominated by a bright green laser beam that illuminates various optical components, including mirrors, lenses, and fiber optic cables. The background is dark, making the green light stand out. The text is overlaid on the image in a yellow, sans-serif font.

22. Lasers

Stimulated Emission: Gain

Population Inversion

Rate equation analysis

**Two-level, three-level, and
four-level systems**

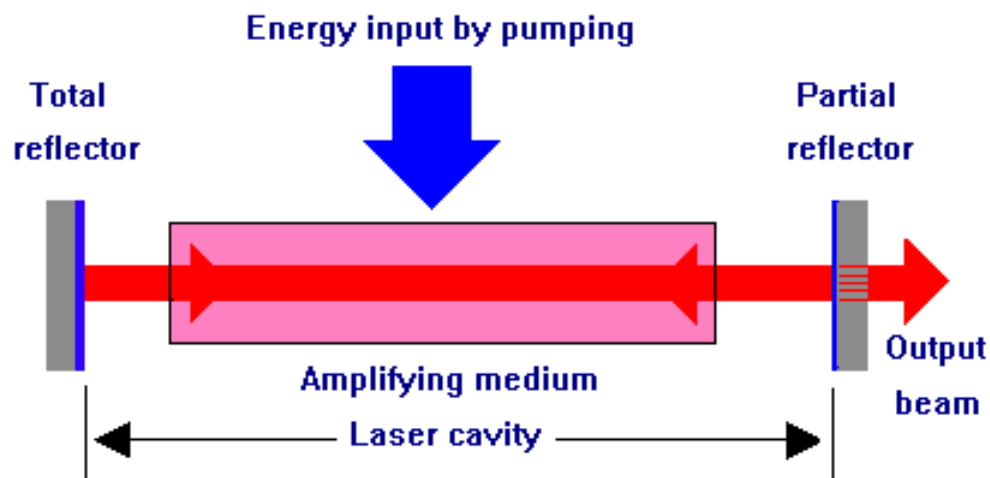
What is a laser?

LASER: **L**ight **A**mplification by **S**timulated **E**mission of **R**adiation

↙ "light" could mean anything from microwaves to x-rays

Essential elements:

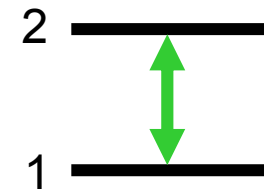
1. A laser medium - a collection of atoms, molecules, etc.
2. A pumping process - puts energy into the laser medium
3. Optical feedback - provides a mechanism for the light to interact (usually many times) with the laser medium



Stimulated emission causes the number of photons in the laser beam to grow.

The factor by which an input beam is amplified by a medium (during one pass through) is called the **gain** and is represented by G .

Reminder: Einstein A and B coefficients



In 1916, Einstein considered the various transition rates between two energy levels (say, 1 and 2) involving light of irradiance, I .

Absorption rate = $B N_1 I$

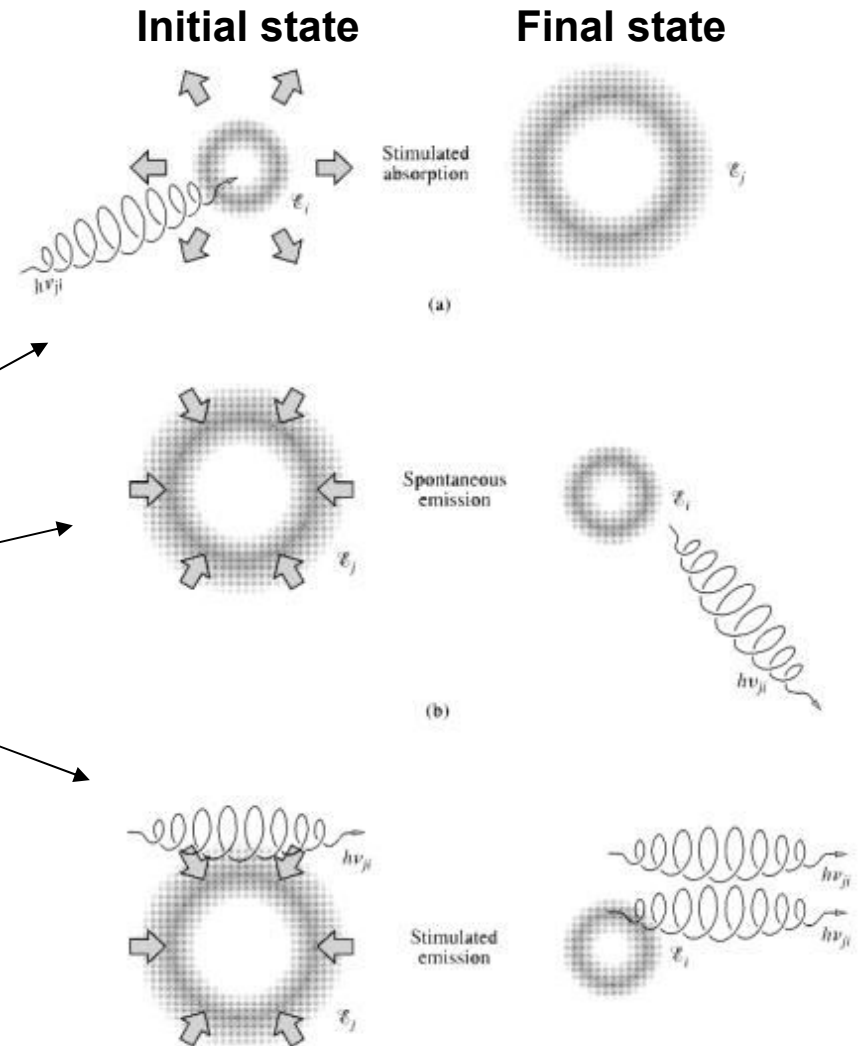
Spontaneous emission rate = $A N_2$

Stimulated emission rate = $B N_2 I$

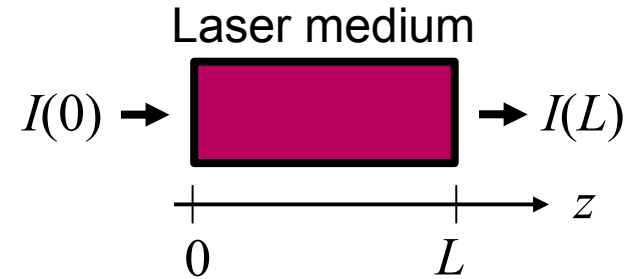
N_1, N_2 = populations of states 1 and 2

I = irradiance of the light field

A, B = Einstein coefficients



Laser gain



Neglecting spontaneous emission:

$$\frac{dI}{dt} = \frac{c}{n} \frac{dI}{dz} \propto BN_2I - BN_1I \quad [\text{Stimulated emission minus absorption}]$$
$$\propto B[N_2 - N_1]I$$

The solution is:

$$I(z) = I(0) \exp\{\sigma[N_2 - N_1]z\} \longrightarrow I(z) = I(0) \exp(-\alpha z)$$

Proportionality constant is called the **cross-section**

There can be exponential gain or loss in irradiance.

Normally, $N_2 < N_1$, and there is loss (absorption).

But if $N_2 > N_1$, there's gain, and we define the gain, G :

$$G \equiv \exp\{\sigma[N_2 - N_1]L\}$$

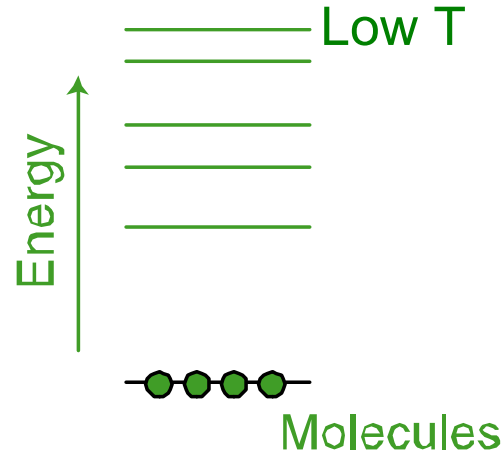
$$\text{If } N_2 > N_1: \quad g \equiv [N_2 - N_1]\sigma$$

$$\text{If } N_2 < N_1: \quad \alpha \equiv [N_1 - N_2]\sigma$$

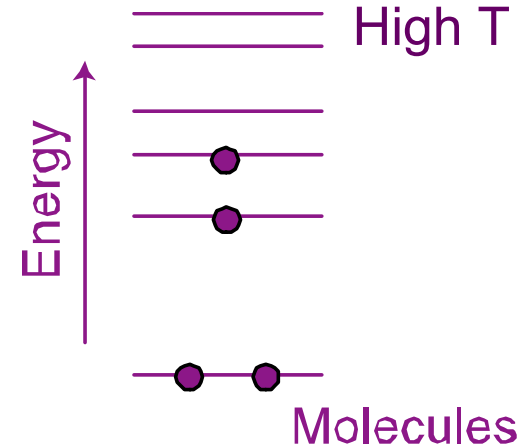
Reminder:

The Maxwell-Boltzmann Distribution

In the absence of collisions, molecules tend to remain in the lowest energy state available.



Collisions can knock a molecule into a higher-energy state. The higher the temperature, the more this happens.



In equilibrium at a temperature T , the ratio of the populations of any two states is given by:

$$\frac{N_2}{N_1} = \frac{\exp[-E_2 / k_B T]}{\exp[-E_1 / k_B T]} = \exp[-\Delta E / k_B T]$$

Since $T > 0$, we always find $N_2 < N_1$.

Populations and the cross-section σ

$$\frac{N_2}{N_1} = \exp[-\Delta E / k_B T]$$

At low temperatures, or for large values of ΔE , we find $N_2 \ll N_1$.

Example:

- The energy level separation is equal to the photon energy.
For visible light, e.g., $\lambda = 600 \text{ nm}$: $h\nu = hc/\lambda = 3.3 \times 10^{-19} \text{ joules}$.
- At room temperature: $k_B T = 4.1 \times 10^{-21} \text{ joules}$.

$$\frac{N_2}{N_1} = \exp[-\Delta E / k_B T] = e^{-79} \sim 10^{-35} \quad \text{So } N_2 = 0, \text{ effectively.}$$

$$\text{If } N_2 > N_1: \quad g \equiv [N_2 - N_1] \sigma$$

$$\text{If } N_2 < N_1: \quad \alpha \equiv [N_1 - N_2] \sigma$$

$$\sigma = \frac{\alpha}{N_1} \quad \text{units of } \sigma \text{ are } \text{m}^2 \text{ (i.e., area)}$$

σ is something like the absorption per molecule. It is the effective size (area) of a molecule as seen by an incoming photon.

Inversion

In order to achieve $G > 1$, stimulated emission must exceed absorption:

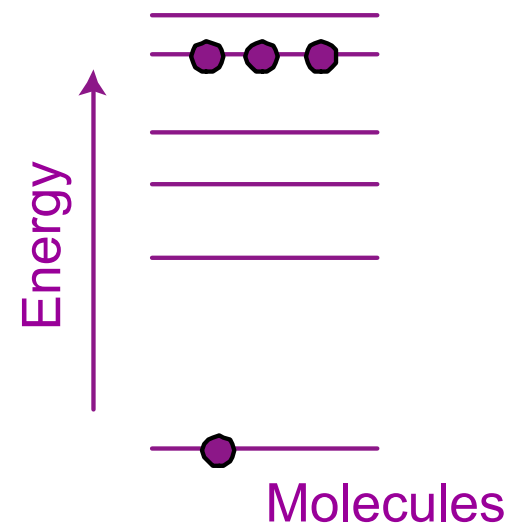
$$B N_2 I > B N_1 I$$

Or, equivalently,

$$N_2 > N_1$$

This condition is called **inversion**. It does not occur naturally in steady-state. It is inherently a non-equilibrium state.

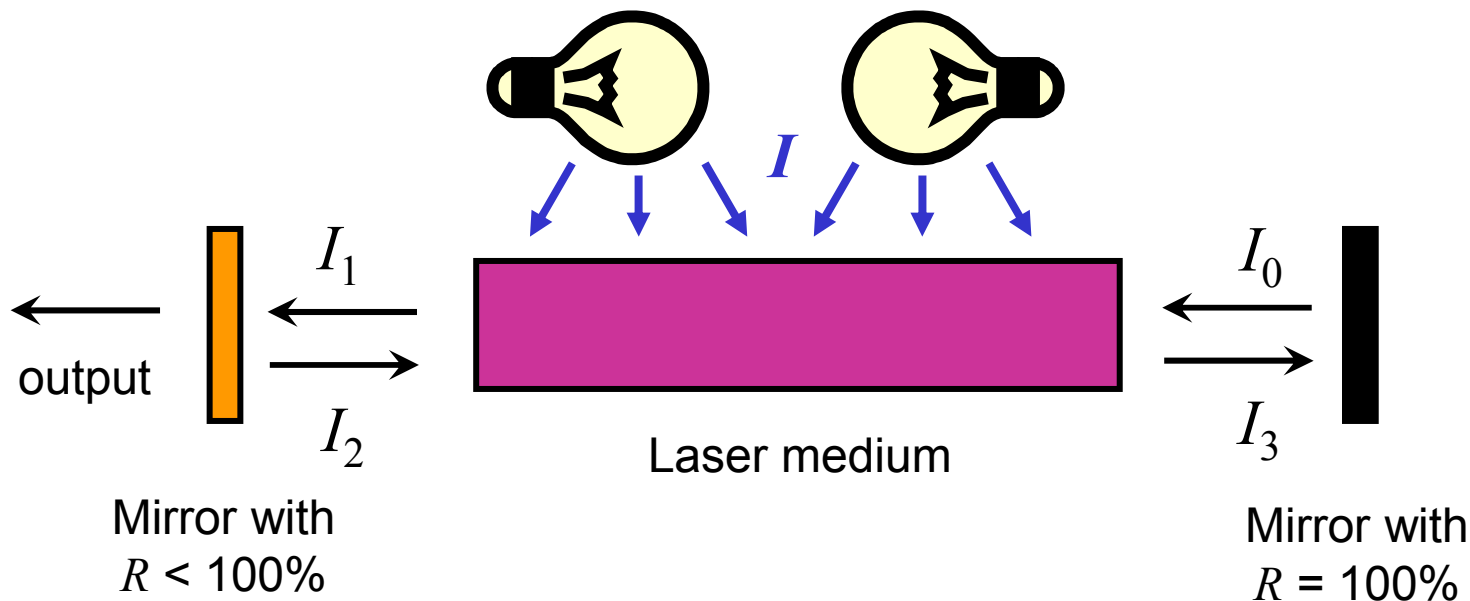
A population inversion



In order to achieve inversion, we must pump a lot of energy into the gain medium. And it needs to be the right medium.

Achieving inversion: Pumping the laser medium

Suppose we pump energy into the laser medium, using another light source with intensity I :



The key question for the remainder of today's lecture:

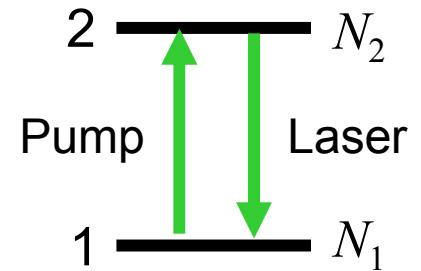
Will this intensity I be sufficient to achieve inversion, $N_2 > N_1$?

The answer depends on the laser medium's energy level configuration.

Rate equations for a two-level system

Earlier we neglected spontaneous emission. Let's look again, and be a bit more careful.

Rate equations for the population densities of the two states:



$$\frac{dN_2}{dt} = BI(N_1 - N_2) - AN_2$$

Absorption
Stimulated emission
Spontaneous emission

Pump intensity

$$\frac{dN_1}{dt} = BI(N_2 - N_1) + AN_2$$

If the total number of molecules is N :

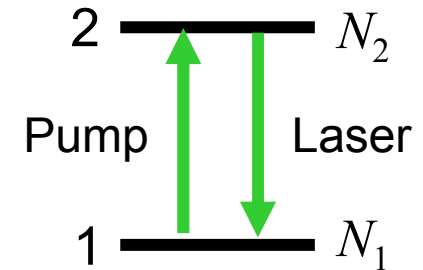
$$N \equiv N_1 + N_2$$

$$\Delta N \equiv N_1 - N_2$$

$$\Rightarrow \frac{d\Delta N}{dt} = -2BI\Delta N + 2AN_2 \quad \leftarrow \begin{aligned} 2N_2 &= (N_1 + N_2) - (N_1 - N_2) \\ &= N - \Delta N \end{aligned}$$

$$\Rightarrow \frac{d\Delta N}{dt} = -2BI\Delta N + AN - A\Delta N$$

How does the population difference depend on pump intensity?



$$\frac{d\Delta N}{dt} = -2BI\Delta N + AN - A\Delta N$$

In steady-state the time derivative is zero: $0 = -2BI\Delta N + AN - A\Delta N$

Solve for ΔN : $(A + 2BI)\Delta N = AN$

$$\Rightarrow \Delta N = AN / (A + 2BI) = \frac{N}{1 + \frac{2B}{A}I}$$

$$\Rightarrow \Delta N = \frac{N}{1 + 2I / I_{sat}}$$

where: $I_{sat} = A / B$

I_{sat} is called the **saturation intensity**, a unique parameter for any gain medium.

What is the saturation intensity?

$$I_{sat} = A / B$$

A is the excited-state relaxation rate due to spontaneous emission: $1/\tau$

B is the absorption cross-section, σ , divided by the energy per photon, $\hbar\omega$: $\sigma/\hbar\omega$

Both σ and τ depend on the frequency of the light. And they are different for each molecule.

The diagram shows the equation $I_{sat} = \frac{\hbar\omega}{\sigma\tau}$ enclosed in a pink rectangular box. Three arrows point from external text to parts of the equation: one to the numerator $\hbar\omega$, one to the denominator $\sigma\tau$, and one to the entire equation. Below the box, an arrow points to the I_{sat} term with the text 10^5 to 10^{13} W/cm². To the right of the box, three lines of text provide typical values: $\hbar\omega \sim 10^{-19}$ J for visible/near IR light, $\tau \sim 10^{-12}$ to 10^{-8} s for molecules, and $\sigma \sim 10^{-20}$ to 10^{-16} cm² for molecules (on resonance).

$$I_{sat} = \frac{\hbar\omega}{\sigma\tau}$$

10^5 to 10^{13} W/cm²

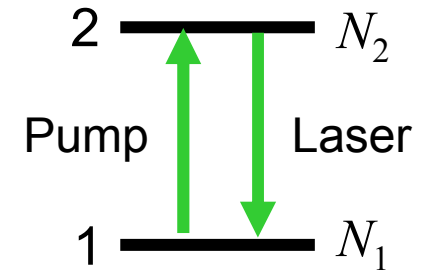
$\hbar\omega \sim 10^{-19}$ J for visible/near IR light

$\tau \sim 10^{-12}$ to 10^{-8} s for molecules

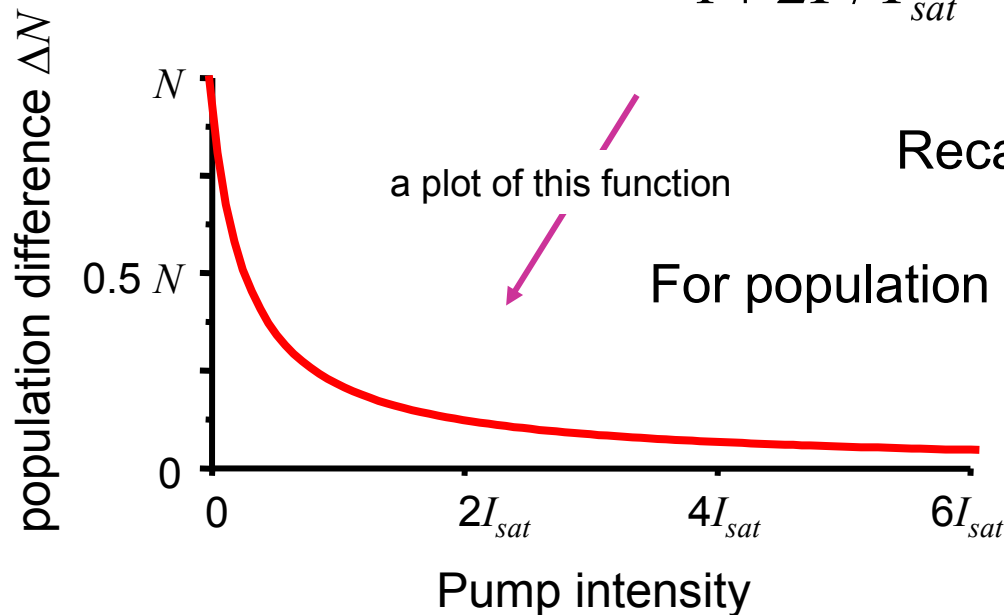
$\sigma \sim 10^{-20}$ to 10^{-16} cm² for molecules (on resonance)

The saturation intensity plays a key role in laser theory. It is the intensity which corresponds to one photon incident on each molecule, within its cross-section σ , per recovery time τ .

Why inversion is impossible in a two-level system



Population difference
$$\Delta N = \frac{N}{1 + 2I / I_{sat}}$$



Recall that $\Delta N \equiv N_1 - N_2$

For population inversion, we require $\Delta N < 0$

ΔN is **always** positive, no matter how hard we pump on the system!

It's impossible to achieve a steady-state inversion in a two-level system!

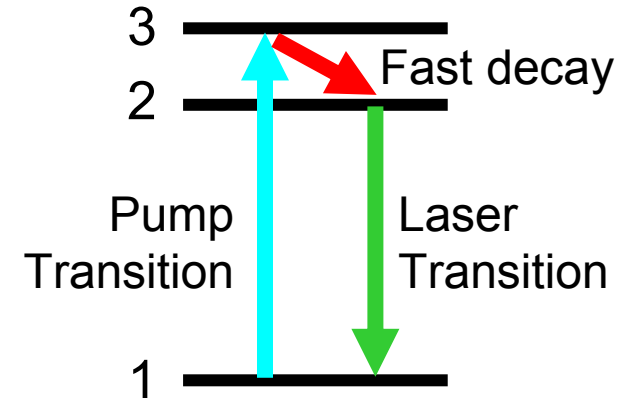
Why? Because absorption and stimulated emission are equally likely.

Even for an **infinite** pump intensity, the best we can do is $N_1 = N_2$ (i.e., $\Delta N = 0$)

Rate equations for a three-level system

So, if we can't make a laser using two levels, what if we try it with three?

Assume we pump to a state 3 that rapidly decays to level 2.



$$\frac{dN_2}{dt} = BIN_1 - AN_2$$

$$\frac{dN_1}{dt} = -BIN_1 + AN_2$$

The total number of molecules is N :

$$N \equiv N_1 + N_2$$

$$\Delta N \equiv N_1 - N_2$$

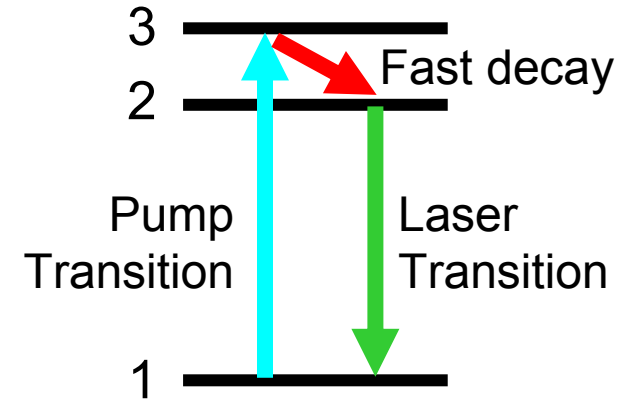
Level 3 decays fast and so $N_3 \approx 0$.

$$\frac{d\Delta N}{dt} = -2BIN_1 + 2AN_2$$

$2N_2 = N - \Delta N$
 $2N_1 = N + \Delta N$

Why inversion is possible in a three-level system

$$\frac{d\Delta N}{dt} = -BIN - BI\Delta N + AN - A\Delta N$$



In steady-state: $0 = -BIN - BI\Delta N + AN - A\Delta N$

Solve for ΔN :
$$\Delta N = N \frac{A - BI}{A + BI} = N \frac{1 - (B/A)I}{1 + (B/A)I}$$

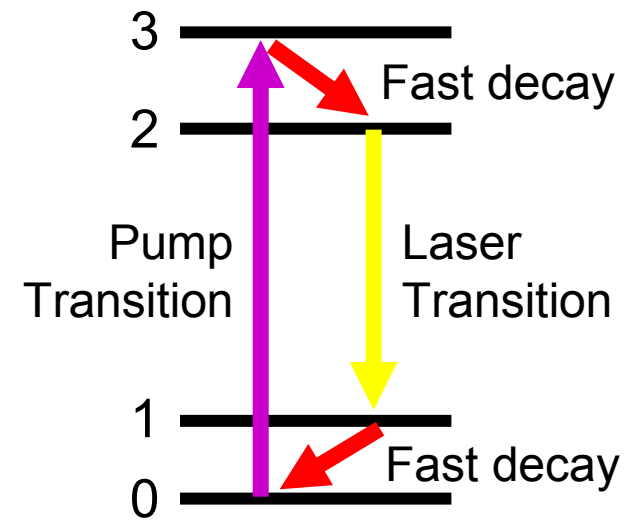
$$\Rightarrow \Delta N = N \frac{1 - I / I_{sat}}{1 + I / I_{sat}}$$

where, as before: $I_{sat} = A / B$
 I_{sat} is the **saturation intensity**.

Now if $I > I_{sat}$, ΔN is negative!

Rate equations for a four-level system

Now assume the lower laser level 1 also rapidly decays to a ground level 0.



As before:
$$\frac{dN_2}{dt} = BIN_0 - AN_2$$

$$\frac{dN_2}{dt} = BI(N - N_2) - AN_2$$

Because $N_1 \approx 0$, $\Delta N \approx -N_2$

$$-\frac{d\Delta N}{dt} = BIN + BI\Delta N + A\Delta N$$

At steady state: $0 = BIN + BI\Delta N + A\Delta N$

The total number of molecules is N :

$$N \equiv N_0 + N_2$$

$$N_0 = N - N_2$$

Why inversion is **easy** in a **four-level** system

$$0 = BIN + BI\Delta N + A\Delta N$$

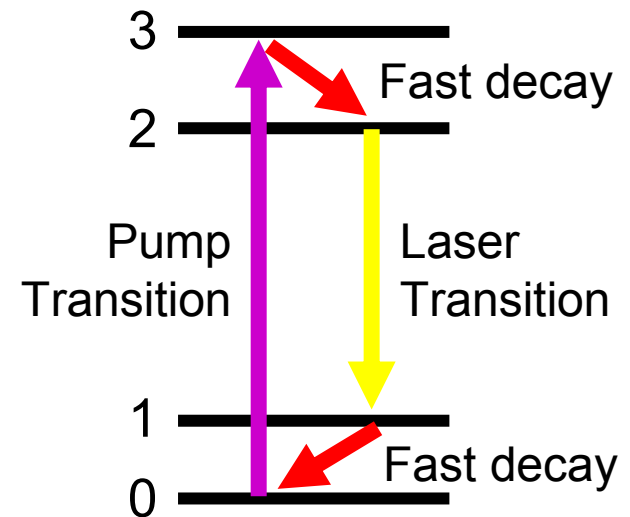
Solve for ΔN :

$$\Rightarrow \Delta N = -BIN / (A + BI) = -N \frac{(B/A)I}{1 + (B/A)I}$$

$$\Rightarrow \boxed{\Delta N = -N \frac{I / I_{sat}}{1 + I / I_{sat}}} \quad \text{where: } I_{sat} = A / B$$

I_{sat} is the **saturation intensity**.

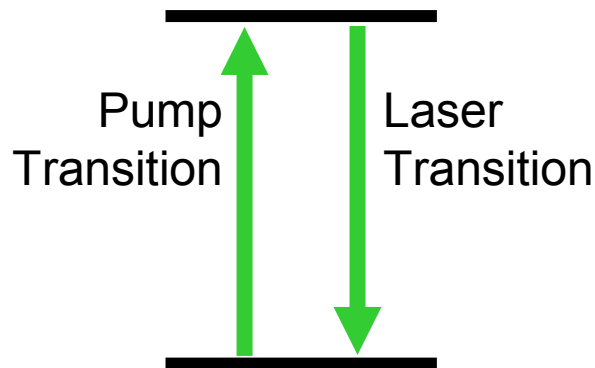
Now, ΔN is negative—for **any** non-zero value of I !



Two-, three-, and four-level systems

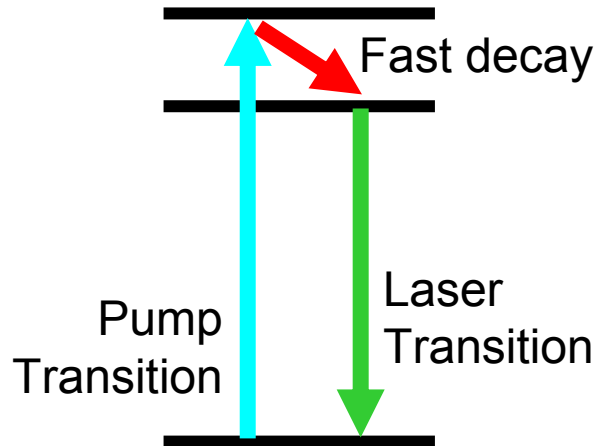
It took laser physicists a while to realize that four-level systems are best.

Two-level system



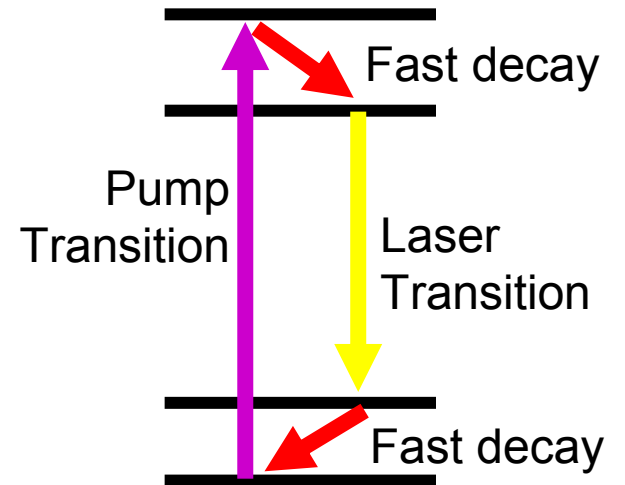
At best, you get equal populations. No lasing.

Three-level system



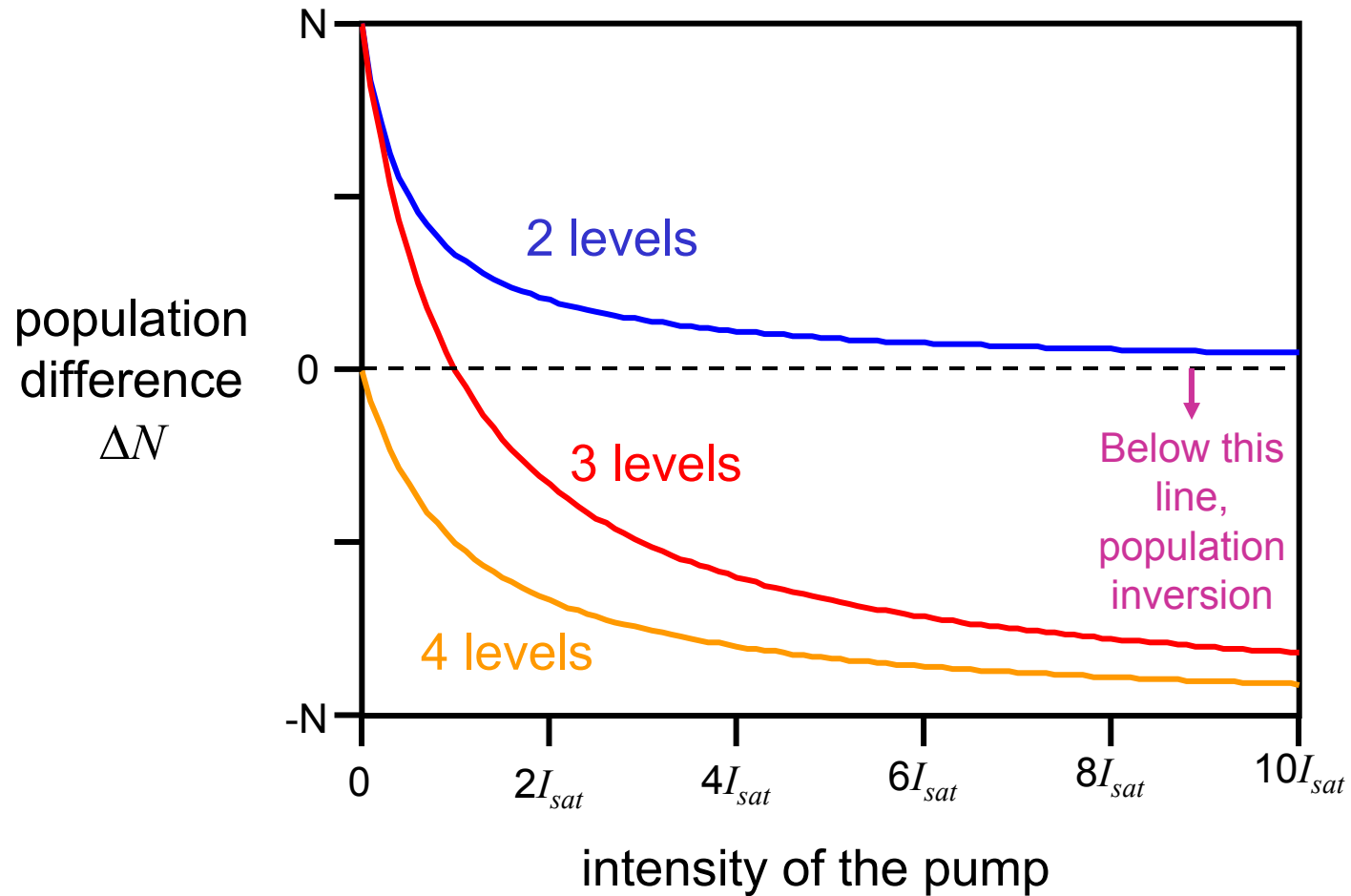
If you hit it hard, you get lasing.

Four-level system

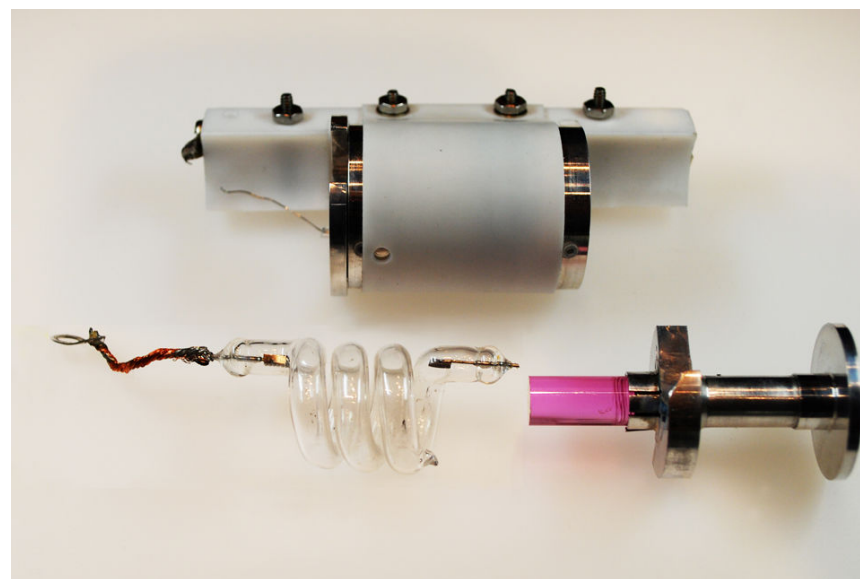
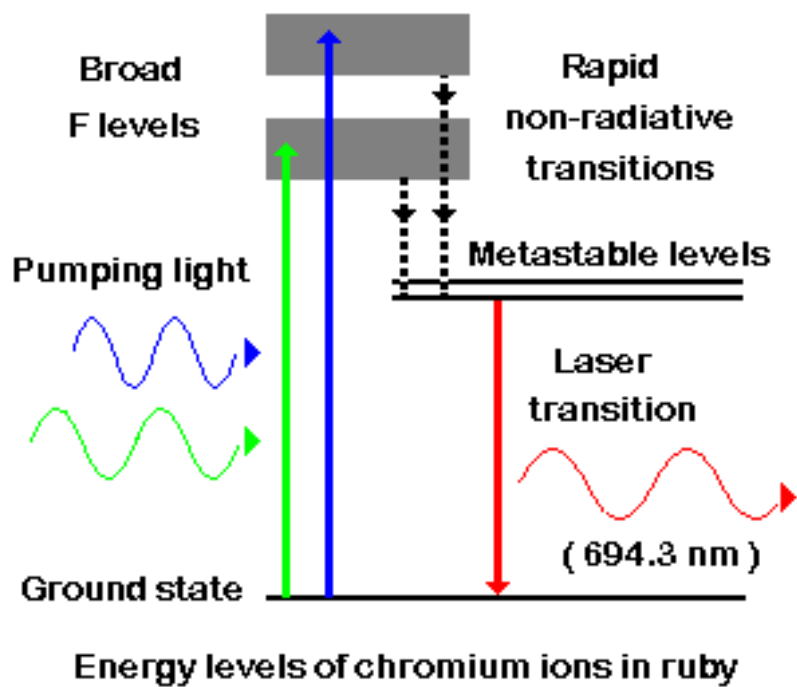


Lasing is easy!

Population inversion in two-, three-, and four-level systems



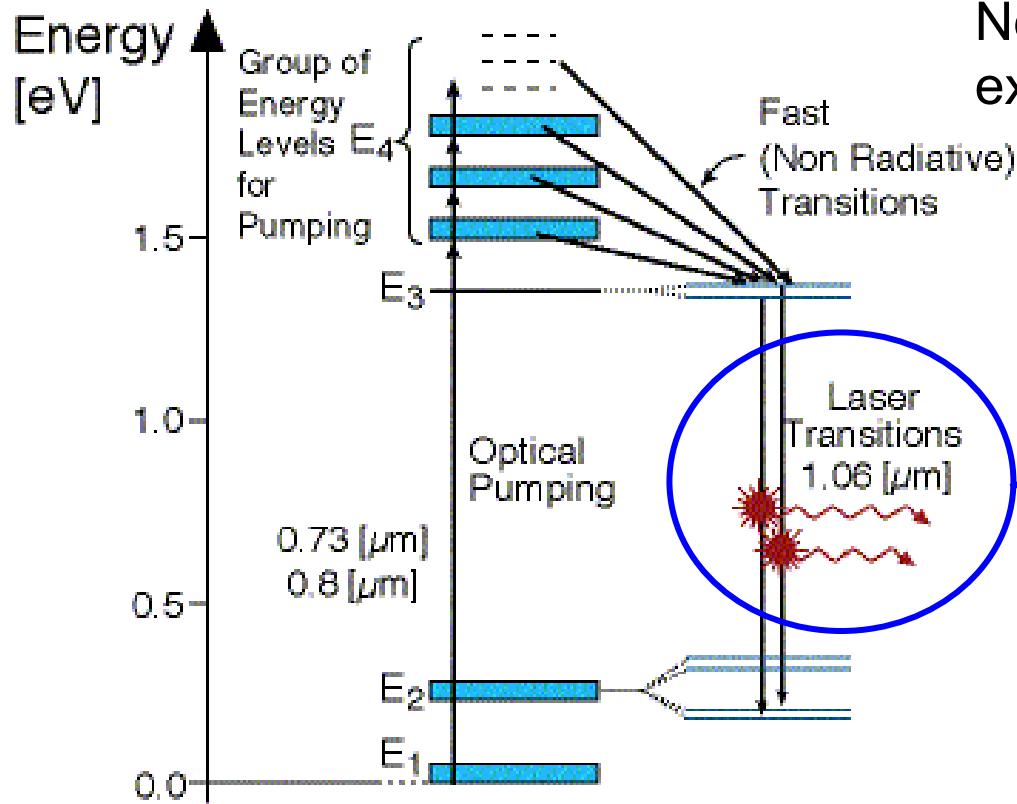
The first laser ever built was a three-level system



Many lasers are almost ideal 4-level systems

A periodic table of elements with a red arrow pointing to Neodymium (Nd), atomic number 60, in the f-block.

For example, the green laser we use in class is an $\text{Nd}^{3+}:\text{YAG}$ laser. The energy levels of Nd^{3+} constitute a classic example of a four-level system.



Infrared laser emission is converted to green light by a nonlinear optical process: **frequency doubling!** (see lecture 36)

The HeNe laser is a 4-level system

Lasing via a three-step process:

1. Ne excited by collisions with He atoms
2. lasing from 5s to 3p energy levels of Ne
3. fast decay to from 3p levels to ground state

There are nine possible 5s-3p transitions, ranging from green (543nm) to red (730nm).

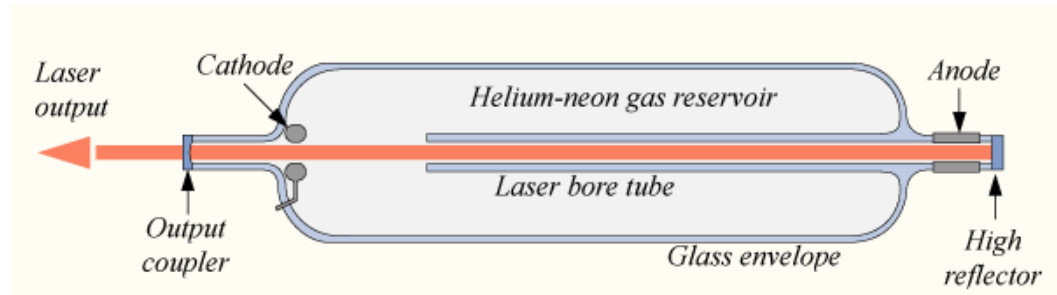
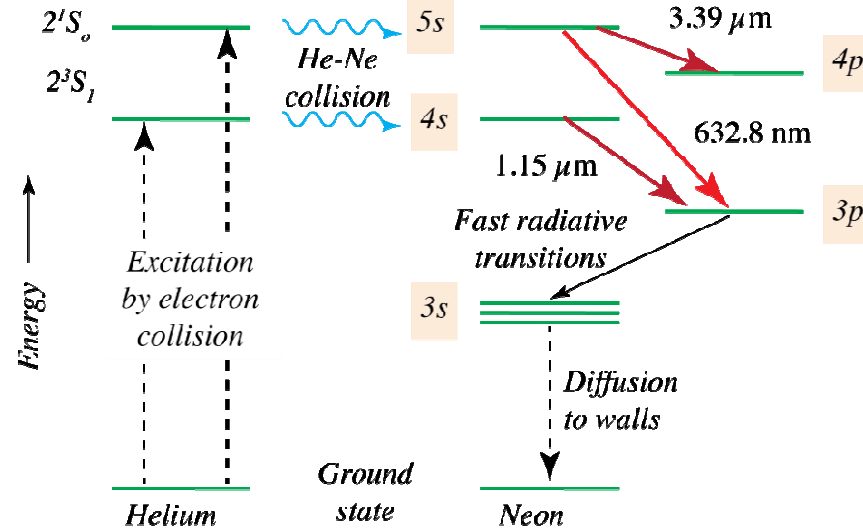


TABLE II. The absolute transition probabilities for the $3s_2-2p_1$ transitions of Ne I.

Transitions	λ (nm)	This expt.	(10^6 sec^{-1})		Decomps and Dumont ^a (Ref. 12)
			Hansch and Toschek (Ref. 1)	Bychkova <i>et al.</i> (Ref. 2)	
$3s_2-2p_1$	730.483	0.255	0.37		
$2p_2$	640.107	1.39			
$2p_3$	635.185	0.345	0.52	0.70	
$2p_4$	632.816	3.39	5.1	6.56	3.24
$2p_5$	629.374	0.639		1.35	
$2p_6$	611.801	0.609	0.93	1.28	
$2p_7$	604.613	0.226	0.39	0.68	
$2p_8$	593.931	0.200		0.56	
$2p_9$	forbidden				
$2p_{10}$	543.365	0.283	0.52	0.59	
Estimated errors		4%	14-21% ^b	17-24% ^b	19%

^a This experiment was based on the Hanle-effect measurements, while the others used the same method as described in the text.
^b These two data sets were based on one laser line at 632.8 nm, and the value of $A_{420,3} = 23.8 \times 10^6 \text{ sec}^{-1}$ (Ladenberg, Ref. 8) was used to convert relative values to absolute values. The error included in this value, which is 30%, is not included in the error estimations given here.

Transition Probabilities for the $5s[1/2]_1-3p$ Transitions of Ne I^{†*}

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 Department of Physics and Astronomy, University of Hawaii, Honolulu, Hawaii 96822
 (Received 12 March 1973)

The HeNe laser is a 4-level system

Question: if there are nine different transitions here, why does the laser produce just one color?

