

Pentahedral Volume, Chaos, and Quantum Gravity

Hal Haggard

May 30, 2012

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ



Polyhedral Volume (Bianchi, Doná and Speziale):

 $\hat{V}_{\mathsf{Pol}} = \mathsf{The}$ volume of a quantum polyhedron



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ



1 Pentahedral Volume

2 Chaos & Quantization

3 Volume Dynamics and Quantum Gravity





1 Pentahedral Volume

2 Chaos & Quantization

3 Volume Dynamics and Quantum Gravity





Minkowski's theorem: polyhedra

The area vectors of a convex polyhedron determine its shape:

$$\vec{A}_1 + \cdots + \vec{A}_n = 0.$$





◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ



Minkowski's theorem: polyhedra

The area vectors of a convex polyhedron determine its shape:

$$\vec{A}_1 + \cdots + \vec{A}_n = 0.$$



Only an existence and uniqueness theorem.



Minkowski's theorem: a tetrahdedron

Interpret the area vectors of tetrahedron as angular momenta:

$$\vec{A}_1 + \vec{A}_2 + \vec{A}_3 + \vec{A}_4 = 0 \quad \iff \quad \checkmark$$

For fixed areas A_1, \ldots, A_4 each area vector lives in S^2 .

Symplectic reduction of $(S^2)^4$ gives rise to the Poisson brackets:

$$\{f,g\} = \sum_{l=1}^{4} \vec{A}_{l} \cdot \left(\frac{\partial f}{\partial \vec{A}_{l}} \times \frac{\partial g}{\partial \vec{A}_{l}}\right)$$



Minkowski's theorem: a tetrahdedron

For fixed areas A_1, \ldots, A_4



 $p = |ec{A}_1 + ec{A}_2|$ q = Angle of rotation generated by p: $\{q,p\} = 1$

イロト イロト イヨト イヨト ヨー のへぐ



Take as Hamiltonian the Volume:

$$H = V^{2} = \frac{2}{9}\vec{A}_{1} \cdot (\vec{A}_{2} \times \vec{A}_{3})$$

$$p$$

$$F_{1}$$

$$E_{2}$$

$$E_{3}$$

$$E_{4}$$

$$E_{4}$$

$$E_{1}$$

$$E_{4}$$

$$E_{4}$$



Require Bohr-Sommerfeld quantization condition,

$$S=\oint_{\gamma}pdq=(n+rac{1}{2})2\pi\hbar.$$

Area of orbits given in terms of complete elliptic integrals,

$$S(E) = \left(\sum_{i=1}^{4} a_i K(m) + \sum_{i=1}^{4} b_i \Pi(\alpha_i^2, m)\right) E$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ



イロト (得) イヨト イヨト ヨー のなべ

Table			
j1 j2 j3 j4	Loop gravity	Bohr- Sommerfeld	Accuracy
6667	1.828	1.795	1.8%
	3.204	3.162	1.3%
	4.225	4.190	0.8%
	5.133	5.105	0.5%
	5.989	5.967	0.4%
	6.817	6.799	0.3%
$\frac{11}{2}$ $\frac{13}{2}$ $\frac{13}{2}$ $\frac{13}{2}$	1.828	1.795	1.8%
	3.204	3.162	1.3%
	4.225	4.190	0.8%
	5.133	5.105	0.5%
	5.989	5.967	0.4%
	6.817	6.799	0.3%

オロト オポト オヨト オヨト ヨー のへで



A pentahedron can be completed to a tetrahedron





A pentahedron can be completed to a tetrahedron





A pentahedron can be completed to a tetrahedron



The volume of the prism is then,

$$V = \frac{\sqrt{2}}{3} \left(\sqrt{\alpha \beta \gamma} - \sqrt{(\alpha - 1)(\beta - 1)(\gamma - 1)} \right) \sqrt{\vec{A}_1 \cdot (\vec{A}_2 \times \vec{A}_3)}$$



What's most difficult about Minkowski reconstruction? Adjacency!



Remarkable side effect of introducing α, β and γ : they completely solve the adjacency problem!

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ▶



Determining the adjacency

Let $W_{ijk} = \vec{A}_i \cdot (\vec{A}_j \times \vec{A}_k)$. Different closures imply,

$$\begin{aligned} \bullet & \alpha_1 \vec{A}_1 + \beta_1 \vec{A}_2 + \gamma_1 \vec{A}_3 + \vec{A}_4 = 0, \\ & \alpha \equiv \alpha_1 = -\frac{W_{234}}{W_{123}} \quad \beta \equiv \beta_1 = \frac{W_{134}}{W_{123}} \quad \gamma \equiv \gamma_1 = -\frac{W_{124}}{W_{123}} \\ \bullet & \alpha_2 \vec{A}_1 + \beta_2 \vec{A}_2 + \vec{A}_3 + \gamma_2 \vec{A}_4 = 0, \\ & \alpha_2 = \frac{W_{234}}{W_{124}} = \frac{\alpha}{\gamma} \quad \beta_2 = -\frac{W_{134}}{W_{124}} = \frac{\beta}{\gamma} \quad \gamma_2 = -\frac{W_{123}}{W_{124}} = \frac{1}{\gamma} \end{aligned}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ

They are mutually incompatible!





Smaller graphs and the associated observables can be consistently included into larger ones



Cylindrical consistency is non-trivially implemented for the polyhedral volume







1 Pentahedral Volume

2 Chaos & Quantization

3 Volume Dynamics and Quantum Gravity





Sommerfeld and Epstein extended Bohr's condition, $L = n\hbar$, as we have seen

$$S = \int_0^T p \frac{dq}{dt} dt = nh$$

and applied it to bounded, separable systems with d degrees of freedom,

$$\int_0^{T_i} p_i \frac{dq_i}{dt} dt = n_i h, \qquad i = 1, \dots, d$$

Here the T_i are the periods of each of the coordinates.

Einstein(!) was not satisfied. These conditions are not invariant under phase space changes of coordinates.



Motivating example: central force problems



In configuration space trajectories cross

・ロト ・ 理 ト ・ ヨト ・ ヨー ・ のへで



Motivating example: central force problems



Momenta are distinct at such a crossing



Motivating example: central force problems



In phase space the distinct momenta lift to the two sheets of a torus



Following Poincaré, Einstein suggested that we use the invariant

 $\sum_{i=1}^d p_i dq_i$

to perform the quantization.

The topology of the torus remains under coordinate changes, and so the quantization condition should be,

$$S_i = \oint_{C_i} \vec{p} \cdot d\vec{q} = n_i h.$$

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ● □ ● ● ● ●





Visualizing dynamics with a surface of section



KAM: Weak perturbation of an integrable system \rightarrow Break up of those tori foliated by trajectories with rational frequency ratios



KAM: Weak perturbation of an integrable system \rightarrow Break up of those tori foliated by trajectories with rational frequency ratios



Toroidal Islands and island chains are left within a sea of chaos



1 Pentahedral Volume

2 Chaos & Quantization

3 Volume Dynamics and Quantum Gravity





The pentahedron has two fundamental degrees of freedom,



The angles generated by $p_1 = |\vec{A}_1 + \vec{A}_2|$ and $p_2 = |\vec{A}_3 + \vec{A}_4|$.

イロン 不得 とうじょう イロン ほう ろくの



For fixed p_1 and p_2 these angles sweep out a torus.



The phase space consists of tori over a convex region of the p_1p_2 -plane.



The volume is a very nonlinear function of any of the variables we have considered:

$$V = \frac{\sqrt{2}}{3} \left(\sqrt{\alpha \beta \gamma} - \sqrt{(\alpha - 1)(\beta - 1)(\gamma - 1)} \right) \sqrt{|\vec{A}_1 \cdot (\vec{A}_2 \times \vec{A}_3)|}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへぐ

Recall,

$$\alpha = -\frac{\vec{A}_4 \cdot (\vec{A}_2 \times \vec{A}_3)}{\vec{A}_1 \cdot (\vec{A}_2 \times \vec{A}_3)}, \qquad \text{similarly for } \beta, \ \gamma$$

Forced to integrate it numerically.



Numerical integration

Fortunately, the angular momenta can be lifted into the phase space of a collection of harmonic oscillators. This allows the use of a geometric (i.e. symplectic) integrator.

Explicit Euler: $u_{n+1} = u_n + h \cdot a(u_n)$

Implicit Euler: $u_{n+1} = u_n + h \cdot a(u_{n+1})$

Symplectic Euler: $u_{n+1} = u_n + h \cdot a(u_n, v_{n+1})$ $v_{n+1} = v_n + h \cdot b(u_n, v_{n+1})$

Implementation: Symplectic integrator preserves face areas to machine precision and volume varies in 14th digit



A Schlegel diagram projects a 3D polyhedron into one of its faces (left panel):



A Schlegel move merges two vertices of the diagram and and splits them apart in a different manner. This is precisely how the volume dynamics changes adjacency.





Poincaré section of pentahedral volume dynamics

Guess: Analogy with billiards systems suggests that the dynamics will be mixed, containing chaos





- Minkowski reconstruction for 5 vectors solved
- There is cylindrical consistency in the polyhedral picture and it is non-trivial
- The classical polyhedral volume is only twice continuously differentiable
- Can explore the classical dynamics of the volume operator in the case of a polyhedron with 5 faces

◆□ ▶ ◆帰 ▶ ◆ ヨ ▶ ◆ ヨ ▶ ● の Q @

Does this dynamics exhibit chaos?