

# WAVE BREAKING IN INHOMOGENEOUS PLASMA. 2. PLASMA CHANNEL FORMATION.

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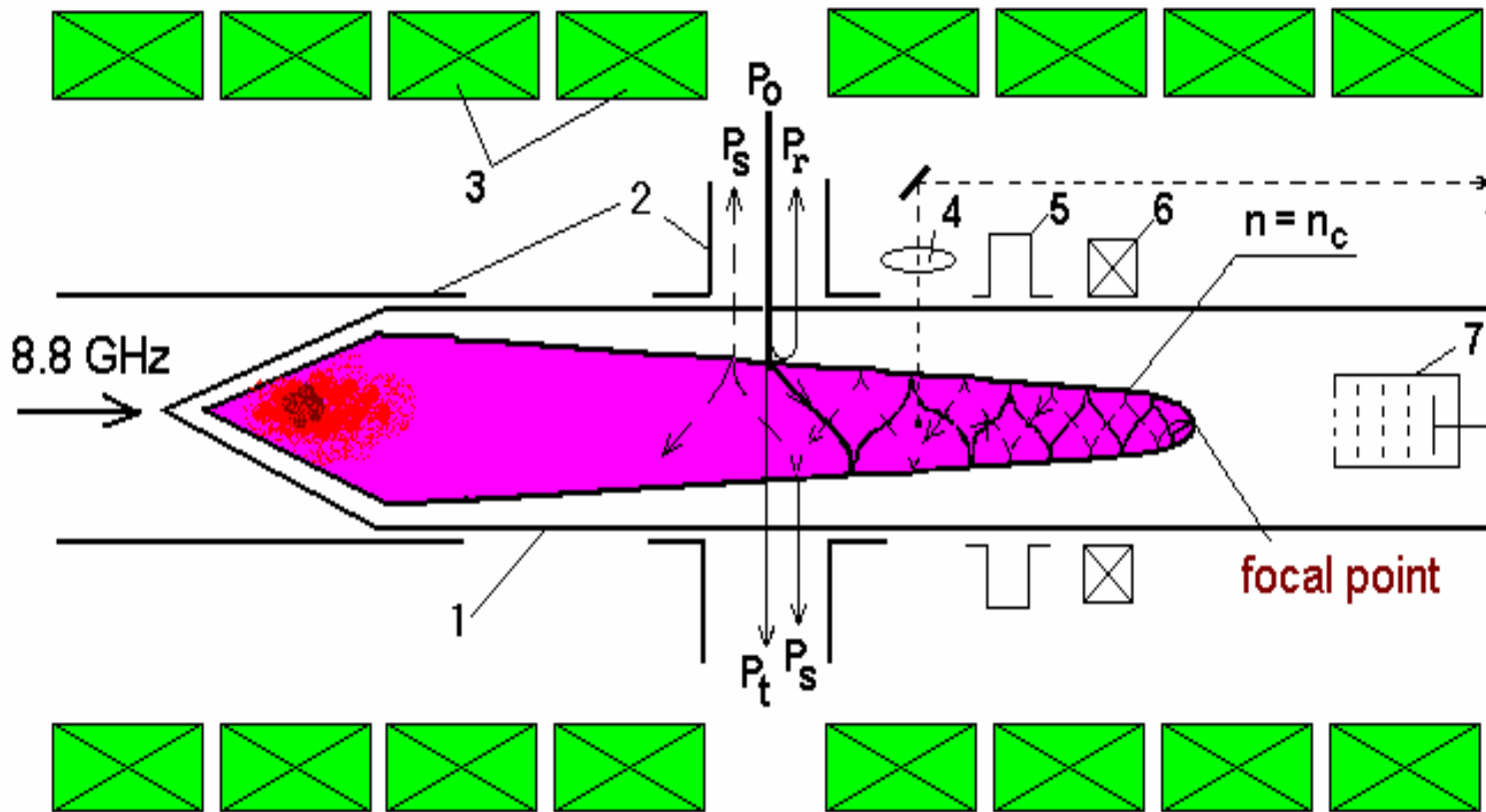
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## Outline

1. Introduction
2. Experimental situation
3. Plasma channel formation
4. Frequency up-shift
5. Conclusion

# Scheme of excitation and propagation of the EPW



- 1 – glass tube
- 2 – waveguides
- 3 – magnet coils
- 4 – optical system
- 5 – resonator
- 6 – Rogowski coil
- 7 – charge particle analyzer

$P_0$ ,  $P_t$ ,  $P_s$ ,  $P_r$  are incident, transient, scattered and reflected power correspondingly

$$n_e(r, z) = n_c \exp\left(-\frac{z}{l}\right) \left(1 - \frac{r^2}{r_0^2}\right)^\beta$$

Argon,  $p \sim 1 - 2 \text{ Pa}$ ,

$H \sim 3.5 \text{ kG}$

$l = 5 \text{ cm}$ ,  $\beta = 4 - 5$

$N_e \sim 10^{12} \text{ cm}^{-3}$ ,  $T_e \sim 1 - 2 \text{ eV}$

# EPW – fundamental Trivelpiece-Gould mode

$$k_{\perp}^2 = \left( \frac{\omega_{pe}^2(r, z)}{\omega_o^2} - 1 \right) \cdot k_{\parallel}^2, \quad k_{\perp} \text{ and } k_{\parallel} - \text{perpendicular and parallel wave vector components}$$

$$E_o = \left( \frac{2P'_0}{\omega_o} \right)^{1/2} \frac{k_o^{3/2}}{\left( 3r_d^2 b k_o^3 + 1 \right)^{1/2}} \exp \left[ i \int_{-\infty}^z (k_o + i k_o'') - \frac{k_o}{2b} r^2 - i \omega_o t \right] + c.c.$$

$$3r_d^2 (k_o + i k_o'')^2 - \frac{z}{a} - \frac{2}{(k_o + i k_o'')b} + i \eta'' = 0,$$

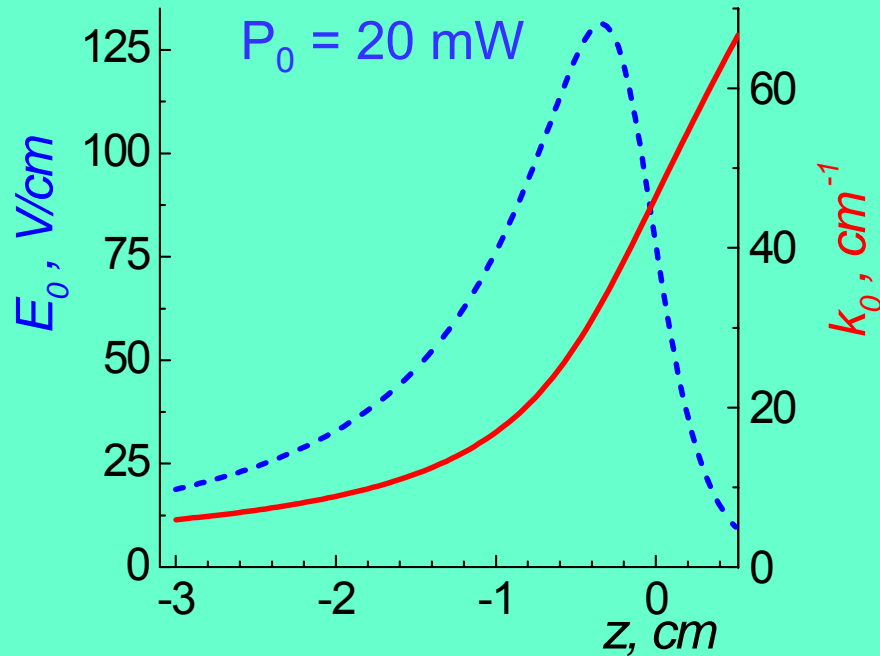
$$\eta = 1 - \frac{\omega_{pe}^2(r, z)}{\omega_o^2} (1 + 3r_d^2 k^2) + i \eta'',$$

$$\eta'' = \frac{v_{ea}}{\omega_o} - \pi \frac{\omega_{pe}^2}{k_o^2} \frac{\partial f_e}{\partial w} \pi \frac{\omega_{pe}^2}{k_o^2} \frac{\partial f_e}{\partial w} \Big|_{w=\frac{\omega_o}{k_o}},$$

$$\ln b_{\ell} = - \int_{-\infty}^z k'' dz' = - \frac{v_{ea}}{\omega_o} k_o a - \pi a \omega_o f_e(\omega_o/k_o).$$

$$a \approx l = 5 \text{ cm}, \quad b = \frac{r_o}{\sqrt{\beta}} \sim 0.4 \text{ cm}$$

# Nonlinear phenomena



## Landau damping

$$n(w) = n_c \left[ \left( \frac{m_e}{2\pi T_e} \right)^{1/2} \exp\left(-\frac{m_e w^2}{2T_e}\right) + \delta \left( \frac{m_e}{2\pi T_h} \right)^{1/2} \exp\left(-\frac{m_e w^2}{2T_h}\right) \right],$$

## Resonance nonlinearity

$$P_0 \geq 10 \text{ mW}$$

$l \rightarrow l' + s$  parametric instability of stimulated backscattering

## Nonresonance nonlinearity

$$P_0 \geq 1 \text{ W}$$

- charge separation under the influence of ponderomotive force
- wavebreaking

## Ionization nonlinearity

$W_~(\text{eV}) \approx 3.2P_0(\text{W})$  – oscillating energy of electrons  
of electrons

$P_0 = 5 \text{ W} \rightarrow W_~ \approx 16 \text{ eV} \geq E_i = 15.76 \text{ eV}.$

# Initial experimental results

$$f_0 = \omega/(2\pi) = \mathbf{2.84 \text{ GHz}}$$

incident pulse power is  $P \sim \mathbf{50 - 200 \text{ W}}$

pulse duration is up to  $\mathbf{2.5 \mu s}$

pulse rise time is  $t_f \sim \mathbf{40 \text{ ns}}$

repetition frequency is  $\mathbf{300 \text{ Hz}}$ .

## Oscillograms:

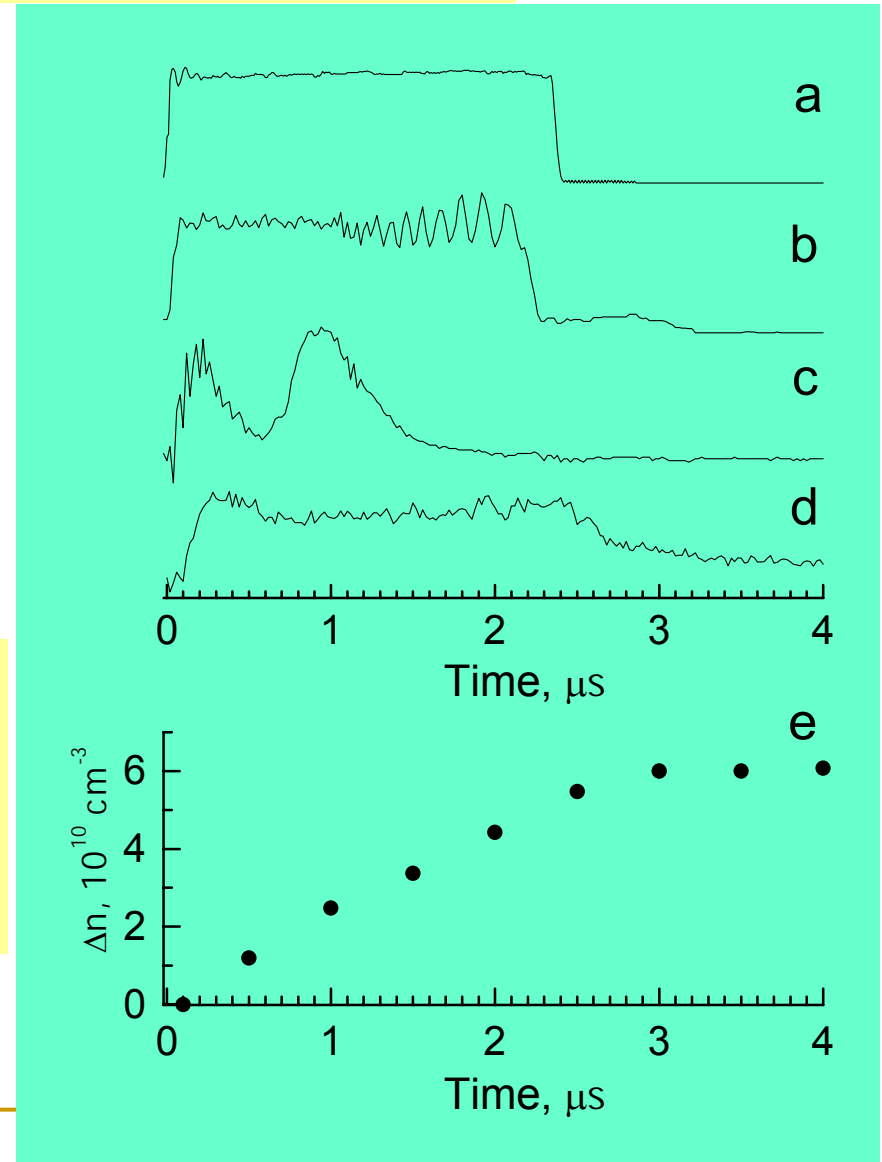
a – incident microwave pulse

b – pulse after an interaction with plasma

c – analyzer current

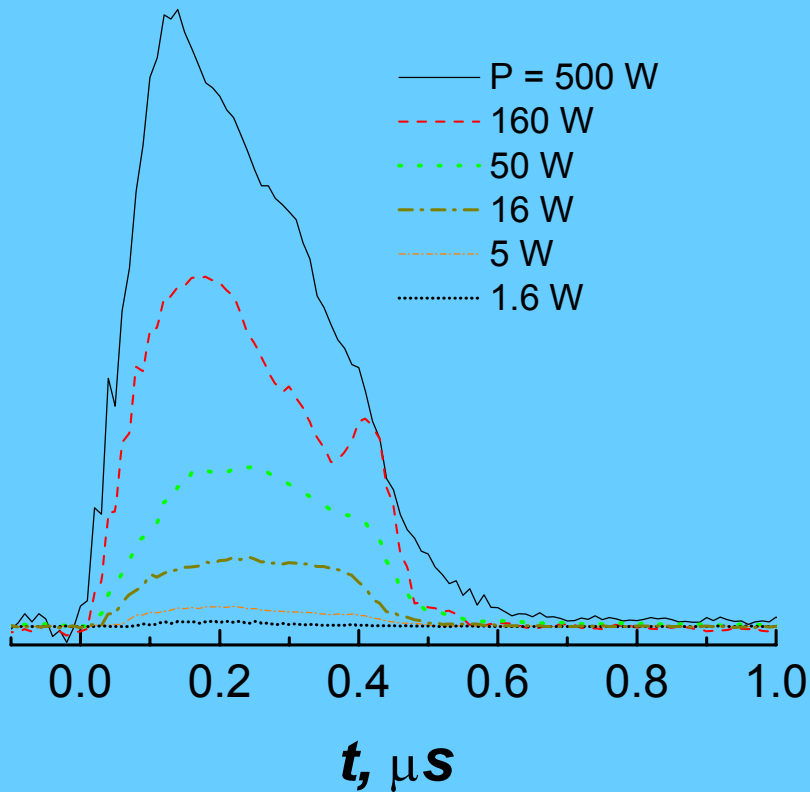
d – integral light from a focal region

e – electron density increase in focal region

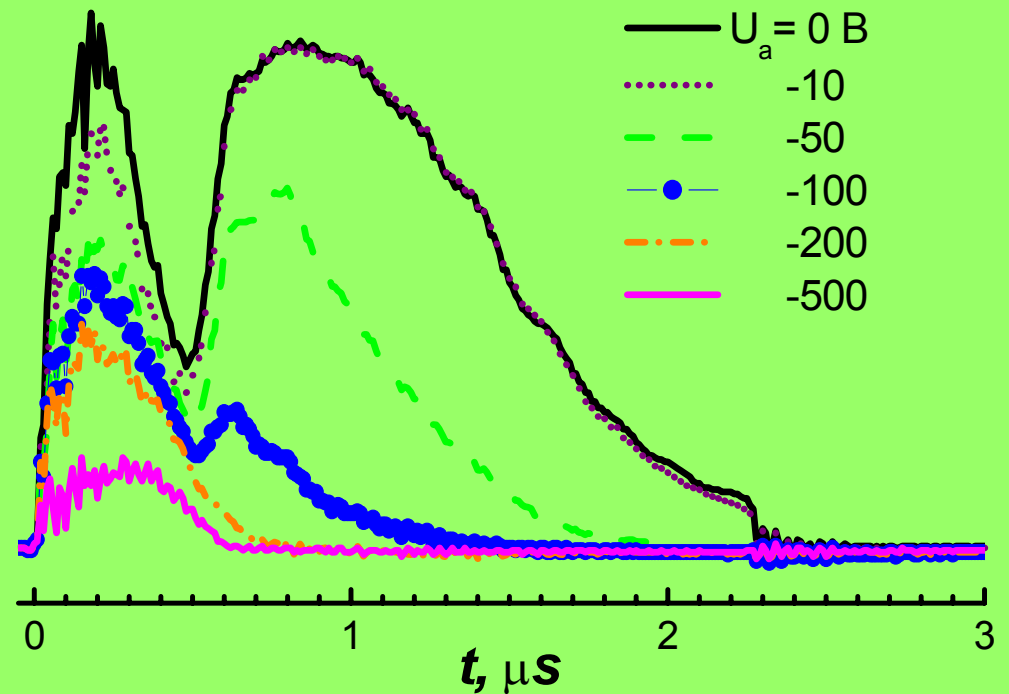


# Oscillograms of a multi-grid analyzer current

$U_a = -10 \text{ V}$

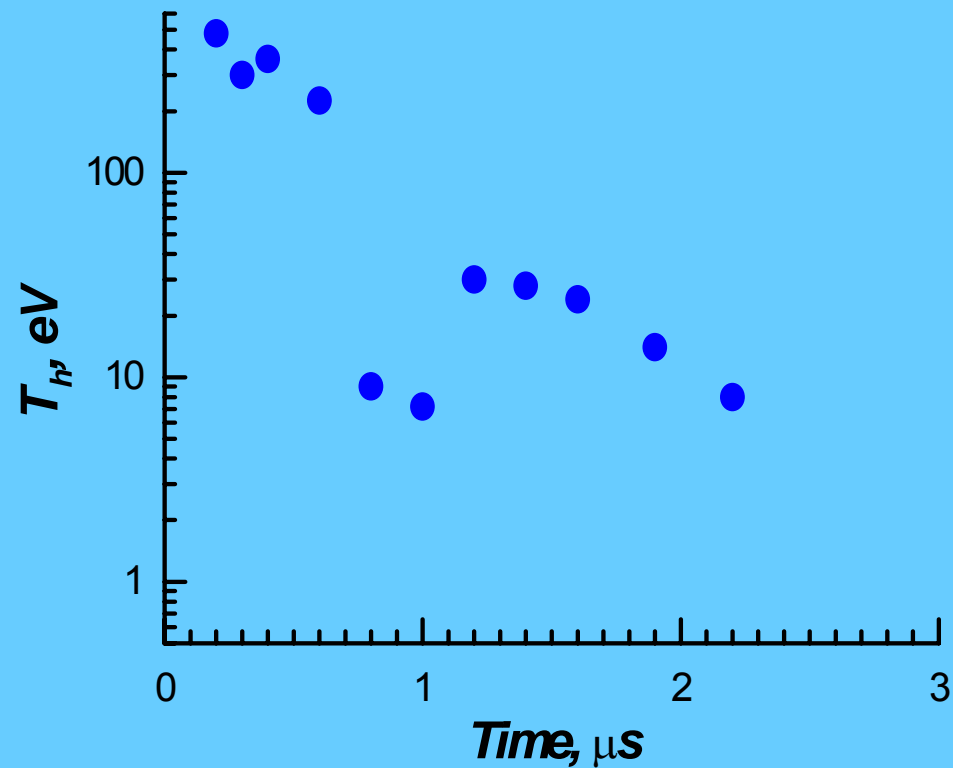
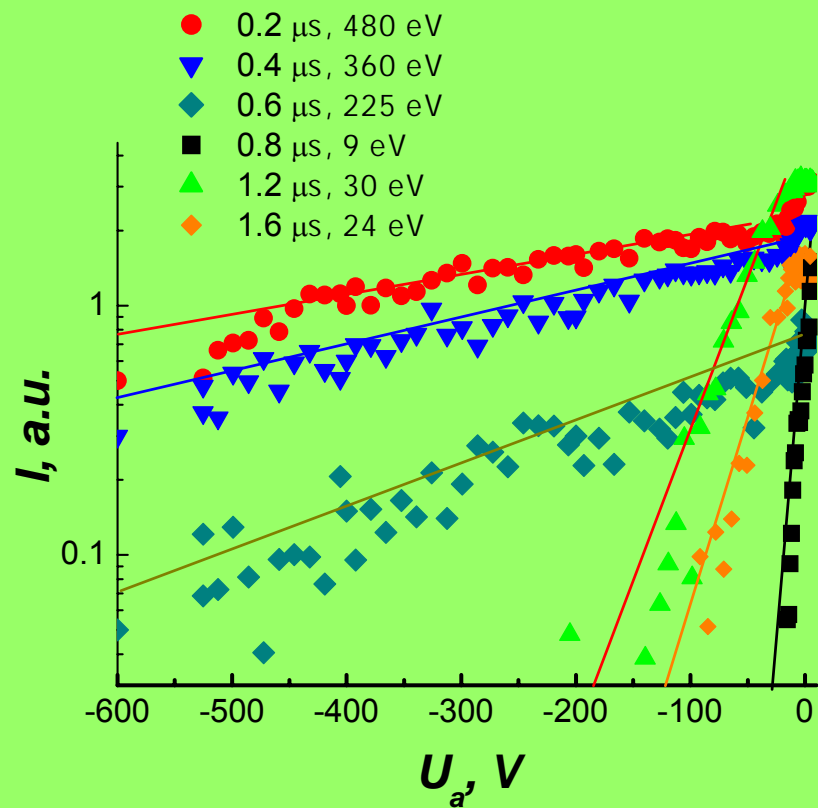


$P_0 = 50 \text{ W}$

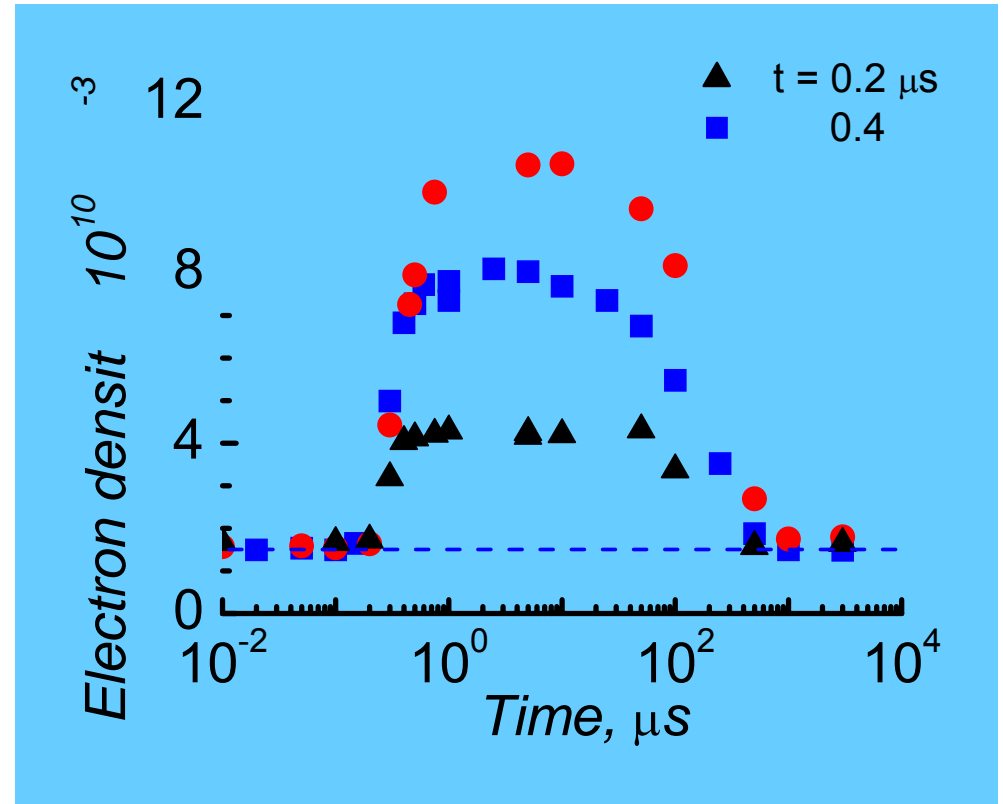
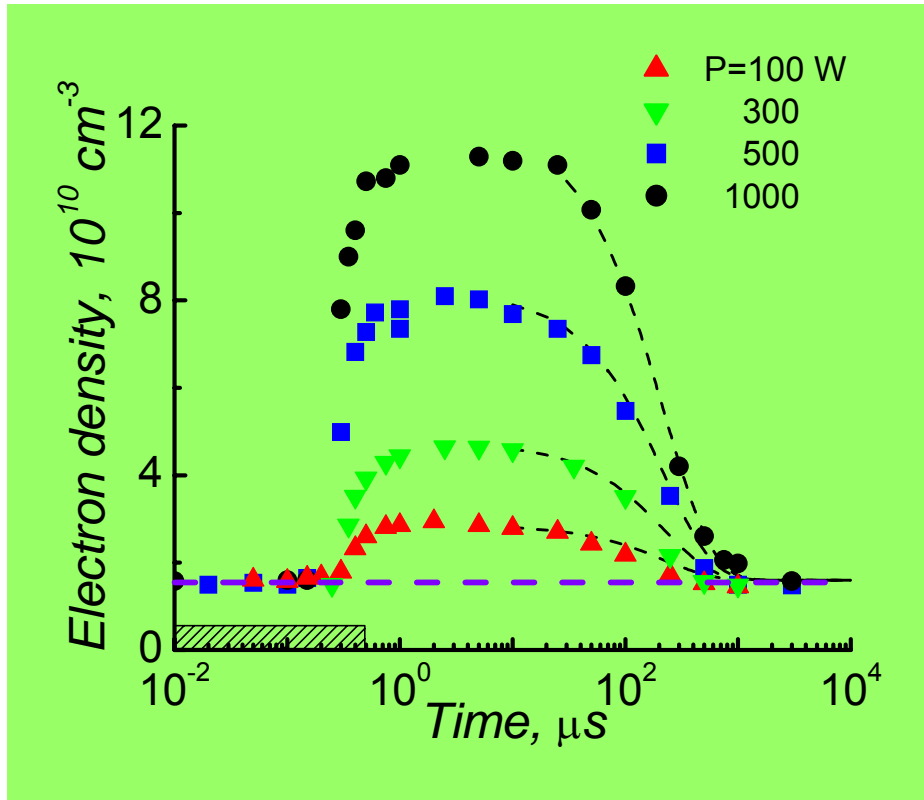


# Effective temperature of accelerated electrons

$P_0 = 50 \text{ W}$



# Electron density dynamics

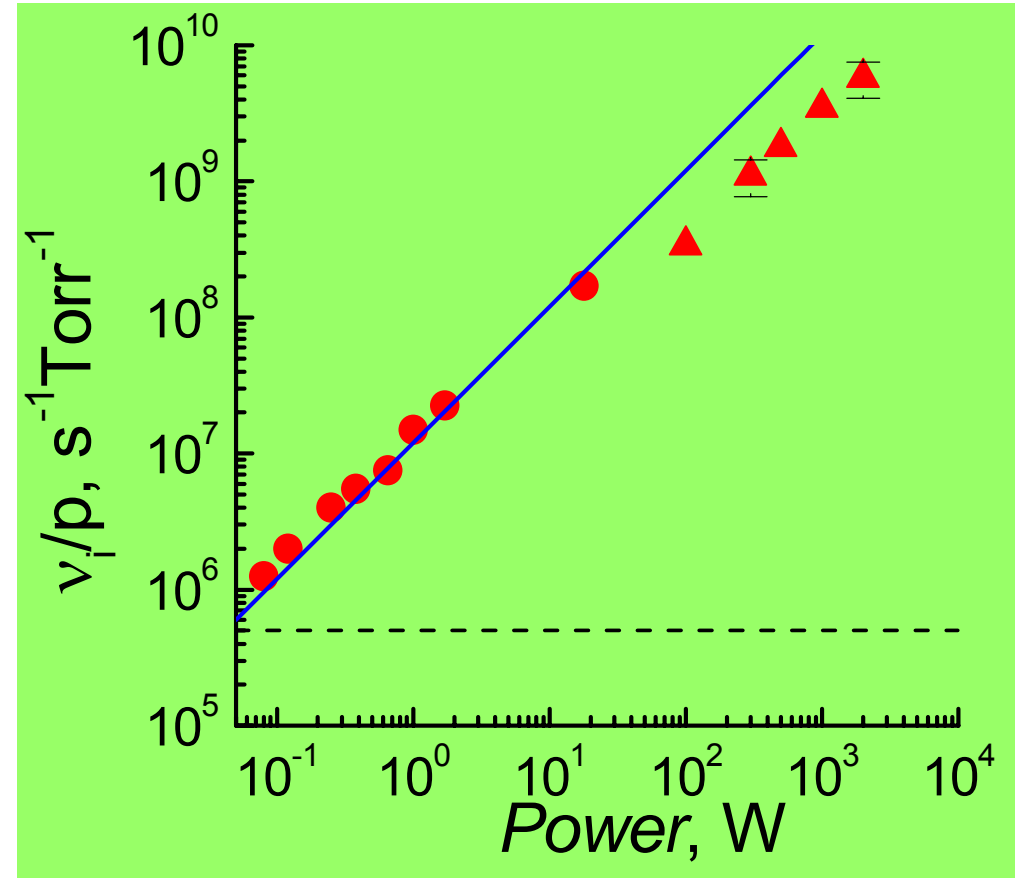
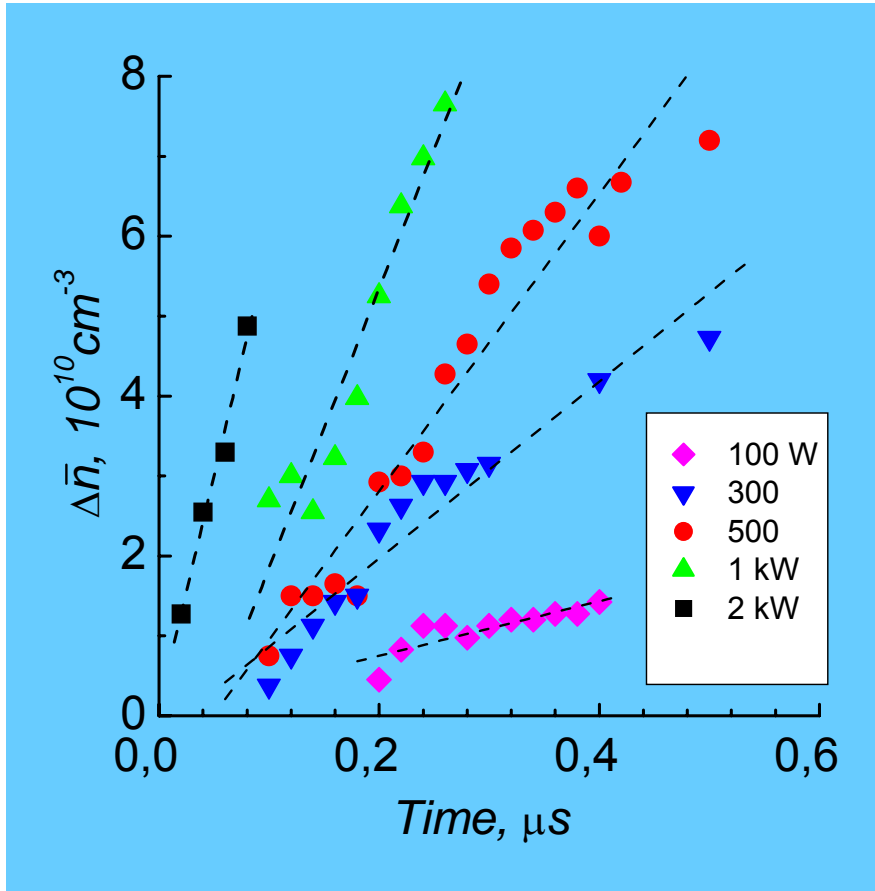


$$n = n_m \exp(-(t-t')/\tau_d)$$

$$\tau_d = 200 \mu\text{s}$$

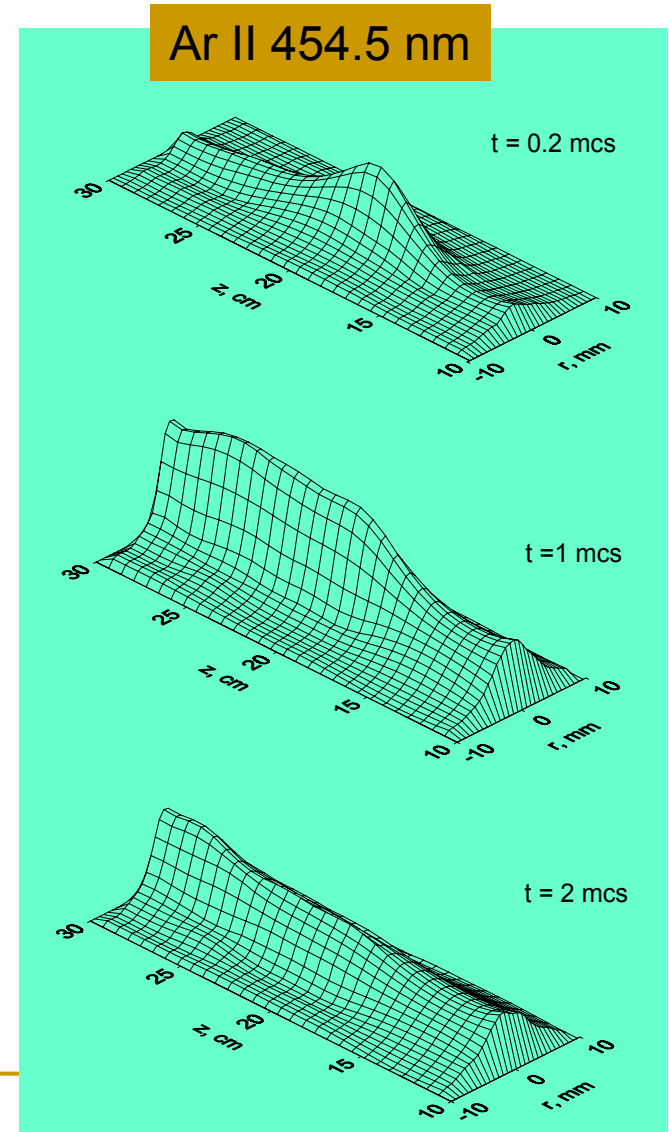
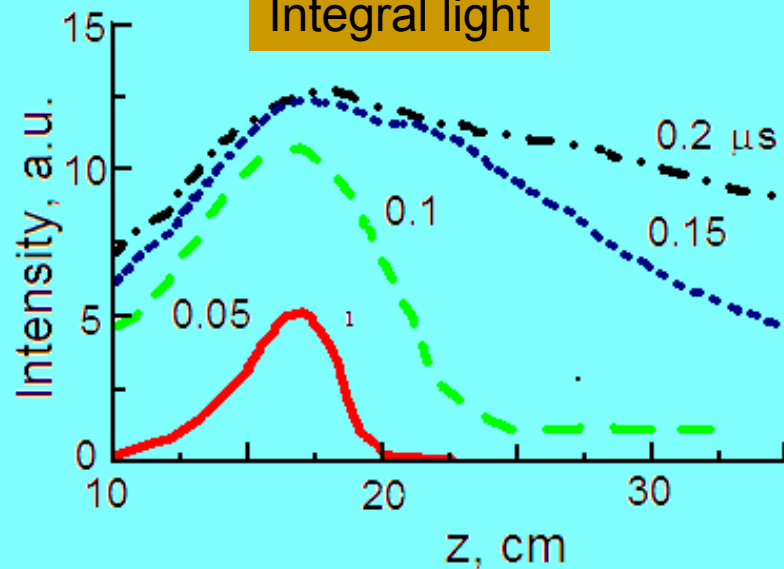
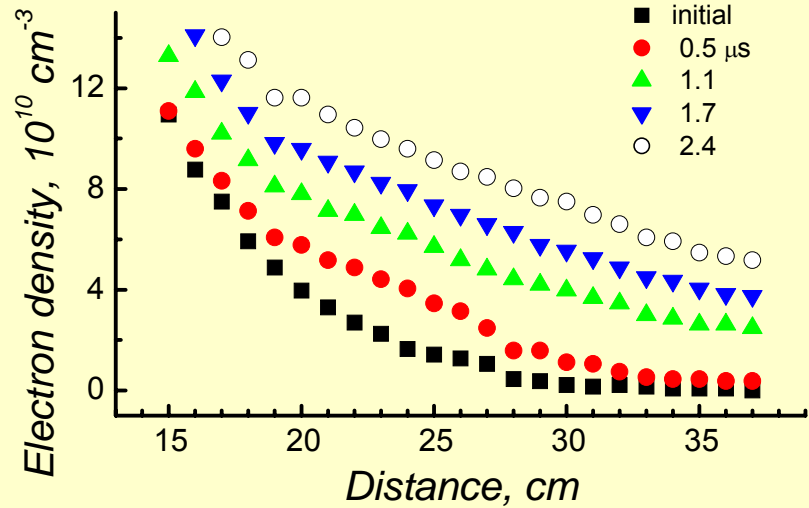


# *Ionization rate*

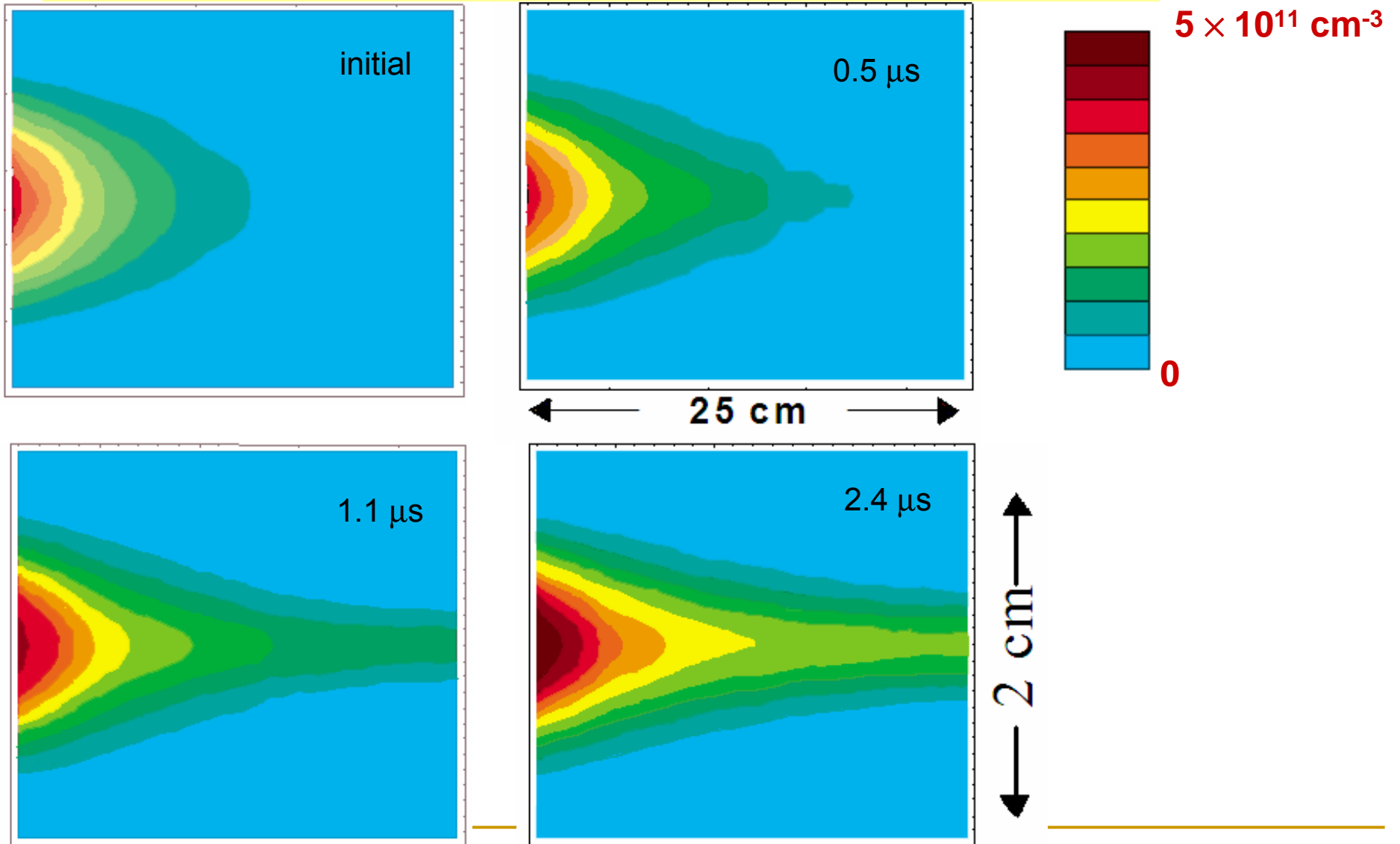


$$\frac{v_i}{p} \approx \frac{1}{pn_c} \frac{\partial n}{\partial t}$$

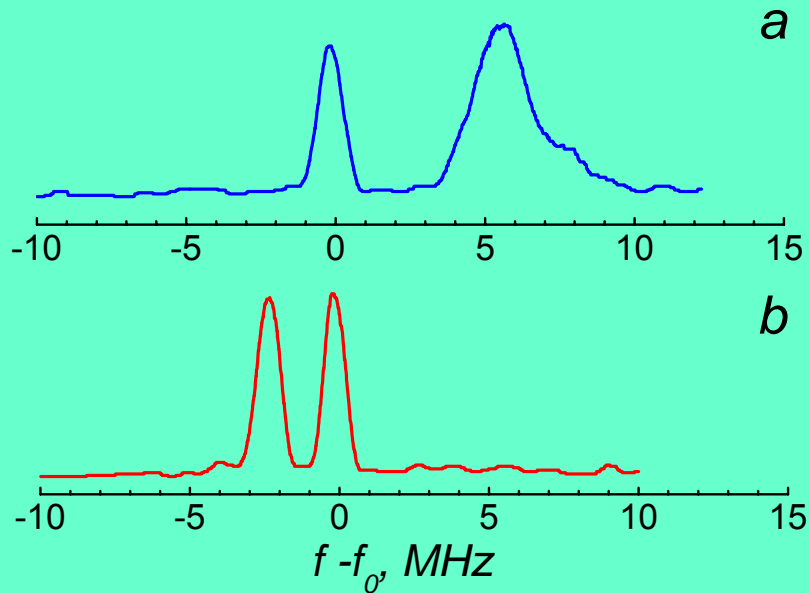
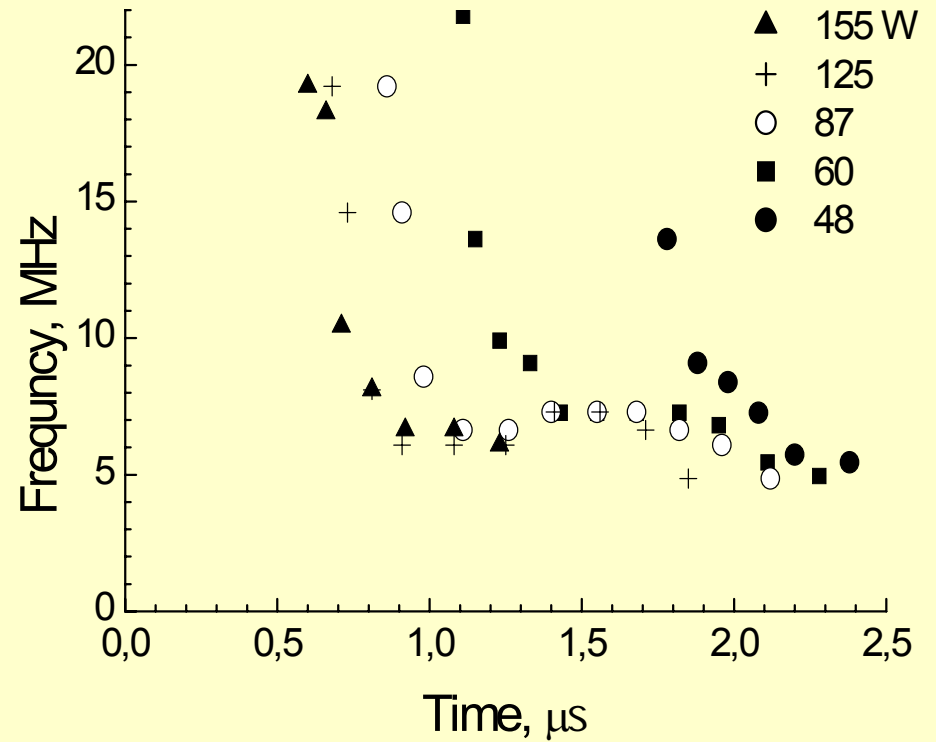
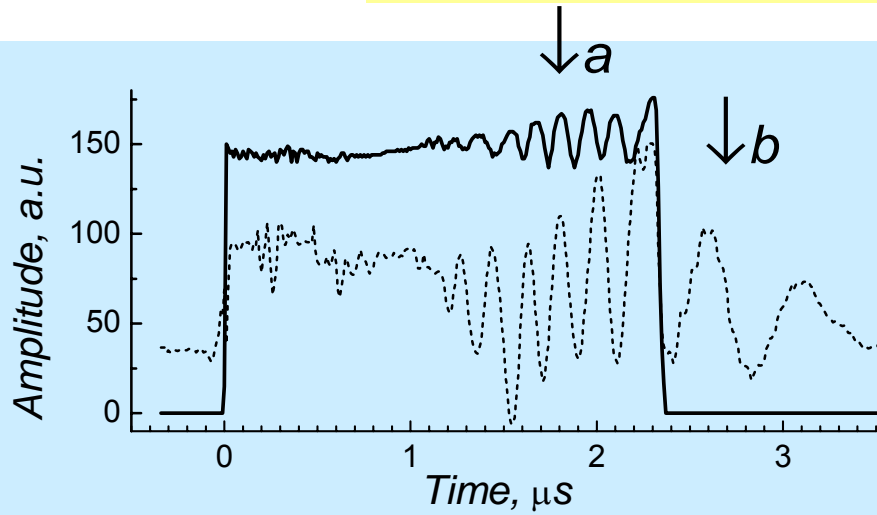
# Electron density and light distributions



# *Dynamics of a plasma waveguide channel*



# Low frequency oscillations



# Theoretical model

$$\delta\Phi = -2k_0L$$

$$\delta\Phi = -2 \frac{L}{b} \frac{n_c}{n_e(z,t) - n_c}$$

$$\Phi = \omega_0 t - 2 \int k_0 dz \quad \omega = \frac{d\Phi}{dt} = \omega_0 + \frac{d\delta\Phi}{dt}$$

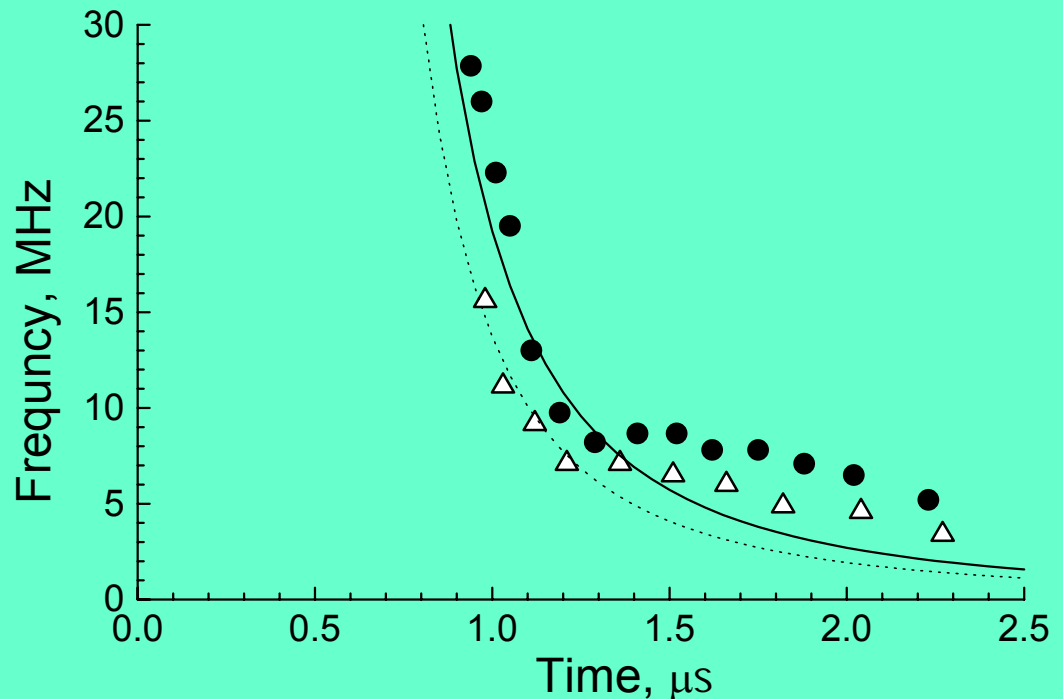
$$\delta f = \frac{\delta\omega}{2\pi} = \frac{1}{2\pi} \frac{L}{b} \frac{n_c}{(n_e(z,t) - n_c)^2} \frac{dn_e}{dt}$$

$$n_e = [1 + \alpha(t)] n_c$$

$$\delta f = \frac{L}{\pi b} \frac{dn_e/dt}{\alpha^2 n_c}$$

$$\frac{dn_e}{dt} = \frac{d\alpha}{dt} n_c = \text{const} \quad \alpha \cong \frac{dn_e}{dt} \frac{t - t_*}{n_c}$$

$$\delta f = \frac{1}{\pi} \frac{L}{b} \frac{n_c}{dn_e/dt (t - t_*)^2}$$



# Conclusion

The interaction of an electromagnetic pulse with inhomogeneous magnetized plasma results in acceleration of electrons due to the wave breaking in vicinity of a resonant point (focus). The acceleration period is less than  $0.5 \mu\text{s}$  at power about 50 W. At large times ( $t > 0.5 \mu\text{s}$ ) a narrow homogeneous plasma channel is created due the fast ionization caused by oscillations of electrons in a wave field. It results in both the increase of collision absorption of the EPW and suppression of electron acceleration effect.

The reflection of the EPW at the edge of the plasma waveguide takes place later ( $t > 1 \mu\text{s}$ ). The propagation of a reflected wave in the plasma waveguide with a increasing electron density results in a phase taper variation and frequency up-shift of the reflected wave.

The work was supported by RFBR-BRFBR collaborative grant (02-02-81033 Bel 2002\_a, F02P-092) and grant INTAS-01-0233.