



Mass effect, broadening and coherence in medium-induced QCD radiation off a $q\bar{q}$ antenna

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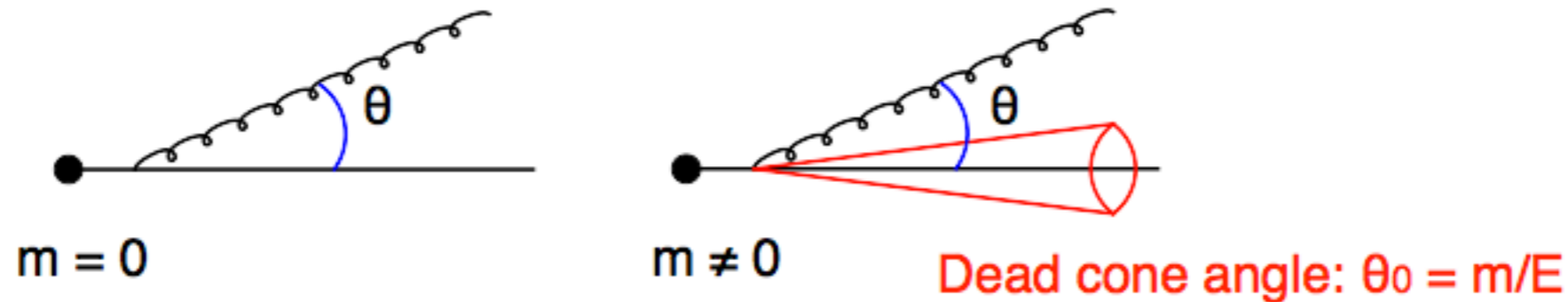
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Motivation

- Dead cone effect suppresses radiation in vacuum



⇒ Gluon radiation inside the dead cone is suppressed.

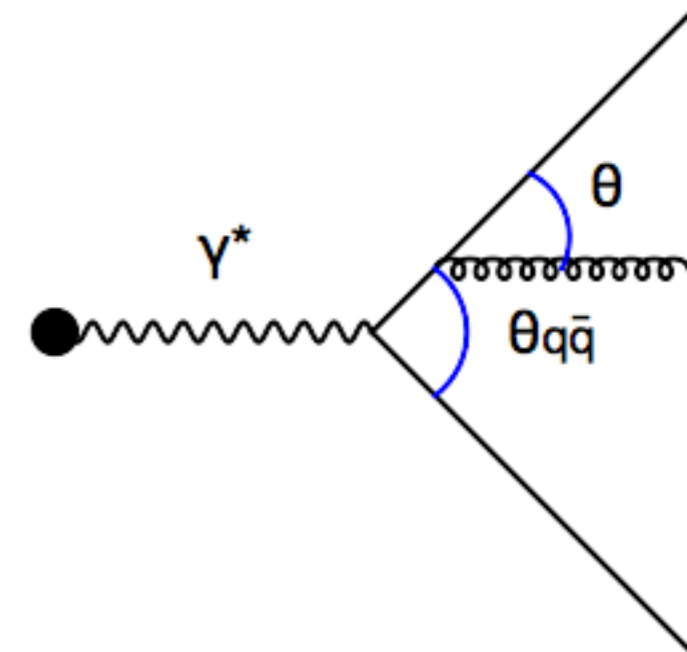
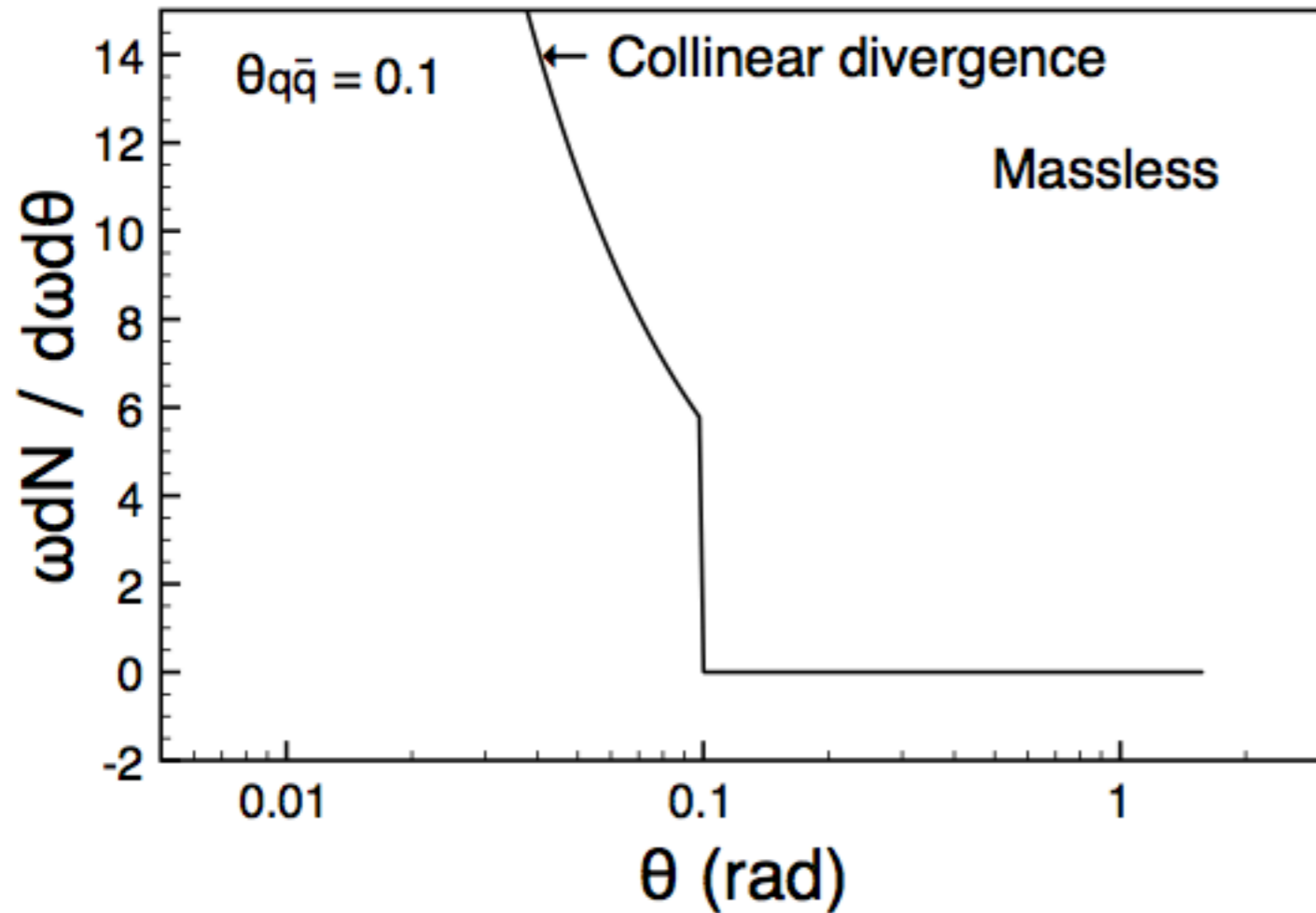
- Mass effects also suppress radiation in medium

⇒ There are remaining puzzles in RHIC and LHC.

- Study the properties of jets originated by heavy quarks

⇒ LHC measures the heavy quark jets.

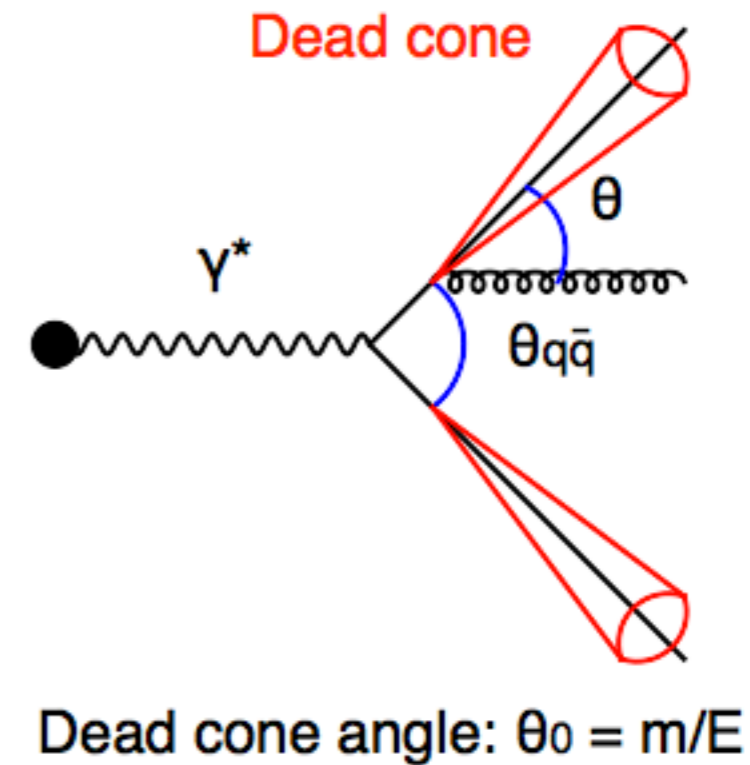
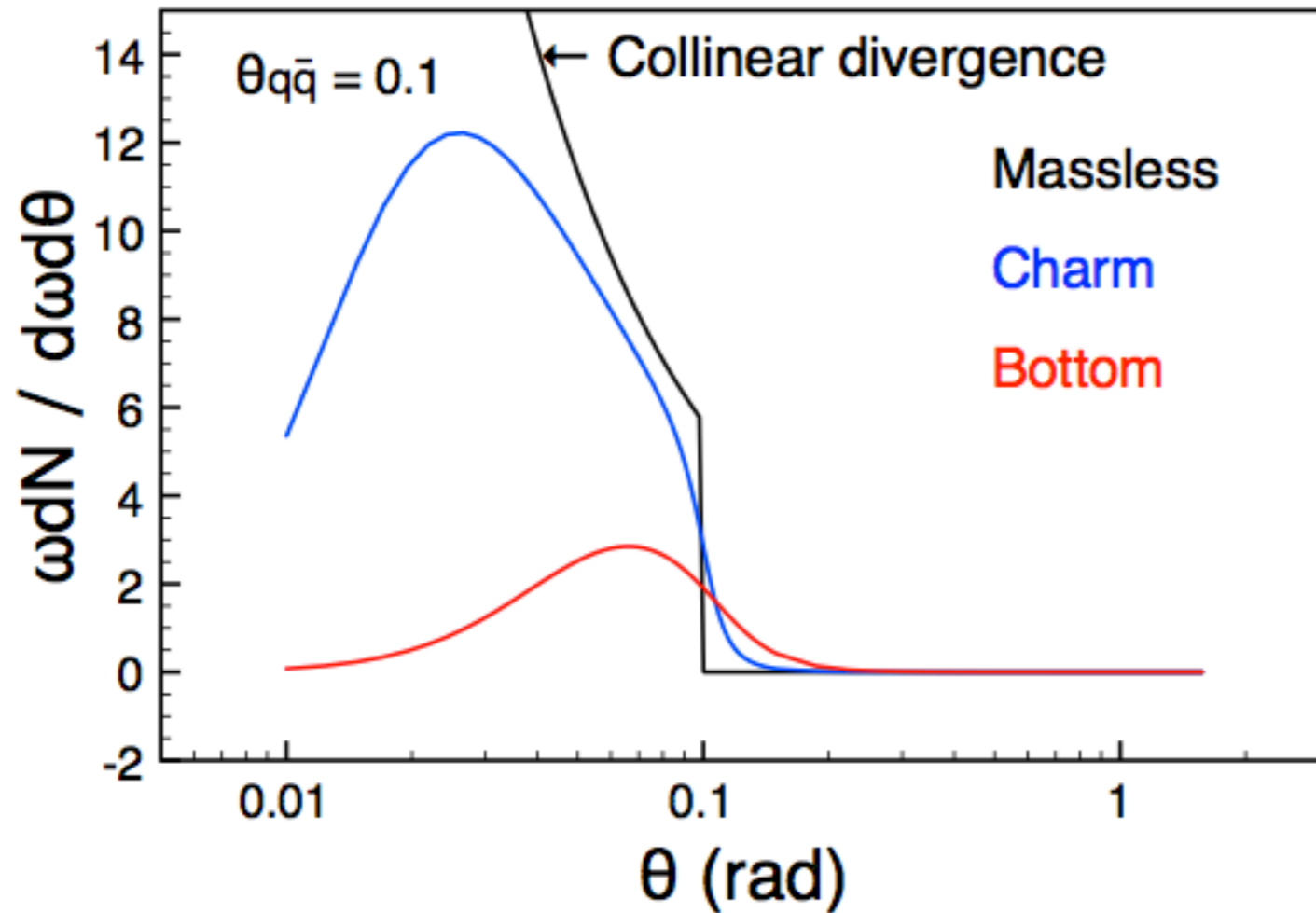
Angular ordering in the soft limit in vacuum



$$\omega \frac{dN_{\text{vac}}}{d\omega d\theta} \propto \frac{1}{\theta} \Theta(\theta_{q\bar{q}} - \theta)$$

Collinear divergence Angular ordering: $\theta_{q\bar{q}} > \theta$

Angular ordering in the soft limit in vacuum

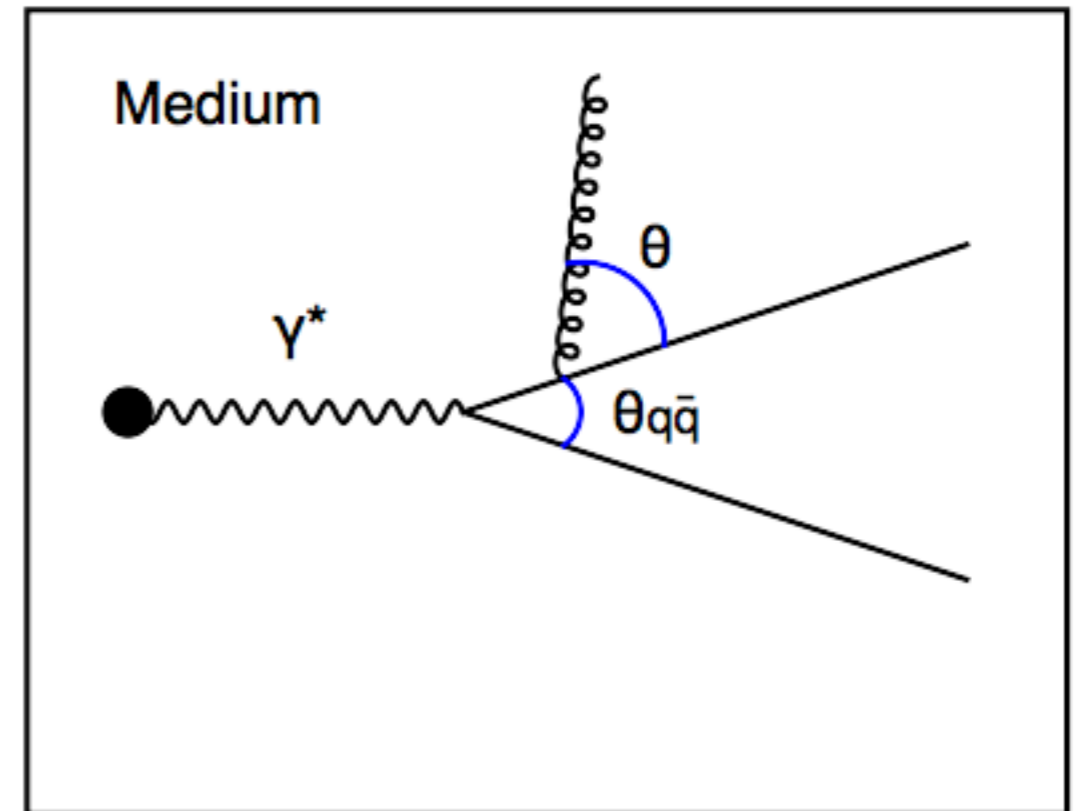
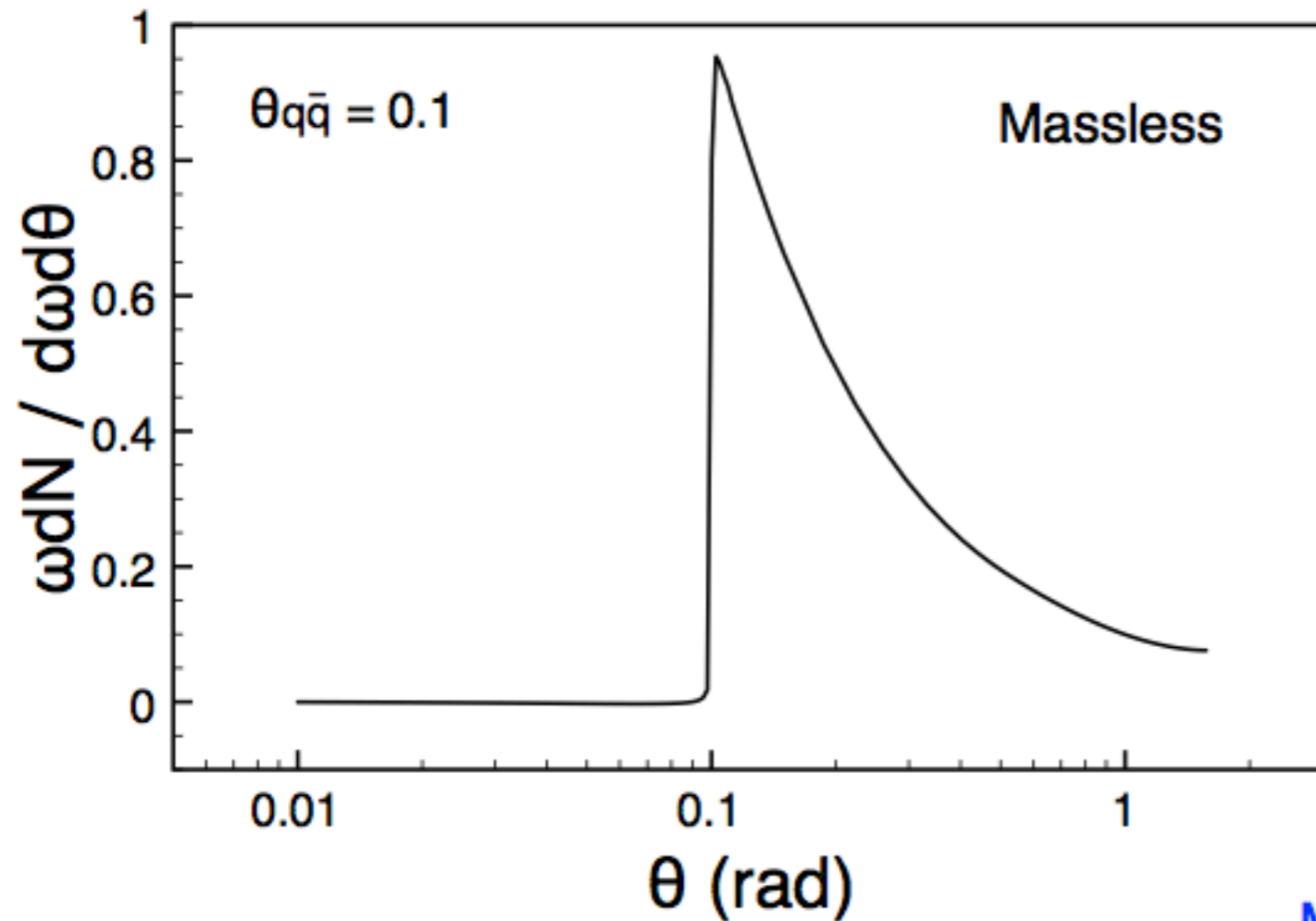


$$\omega \frac{dN_{\text{vac}}}{d\omega d\theta} \propto \frac{\theta}{\theta^2 + \theta_0^2} H_{\text{vac}}(\theta_{q\bar{q}}, \theta_0, \theta)$$

⇒ Both **collinear divergence** and **angular ordering** are modified by the dead cone angle θ_0 .

Medium-induced antiangular ordering in the soft limit

Mehtar-Tani, Salgado and Tywoniuk, *Phys. Rev. Lett.* 106 (2011) 122002



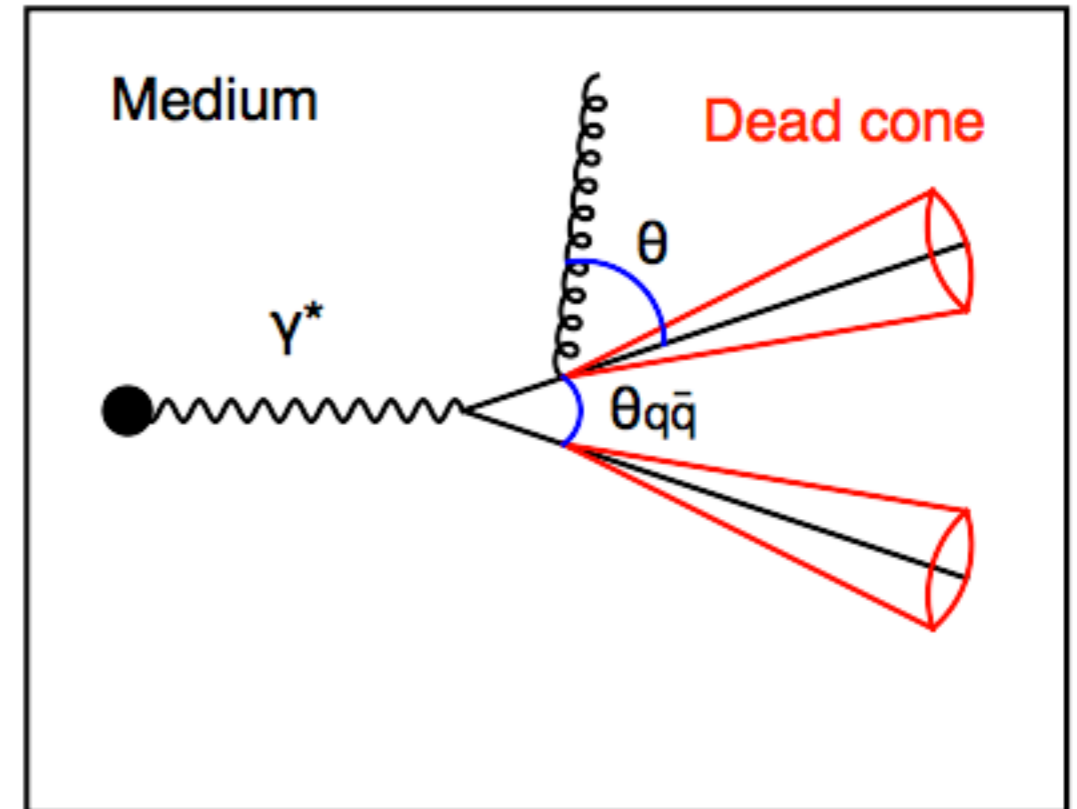
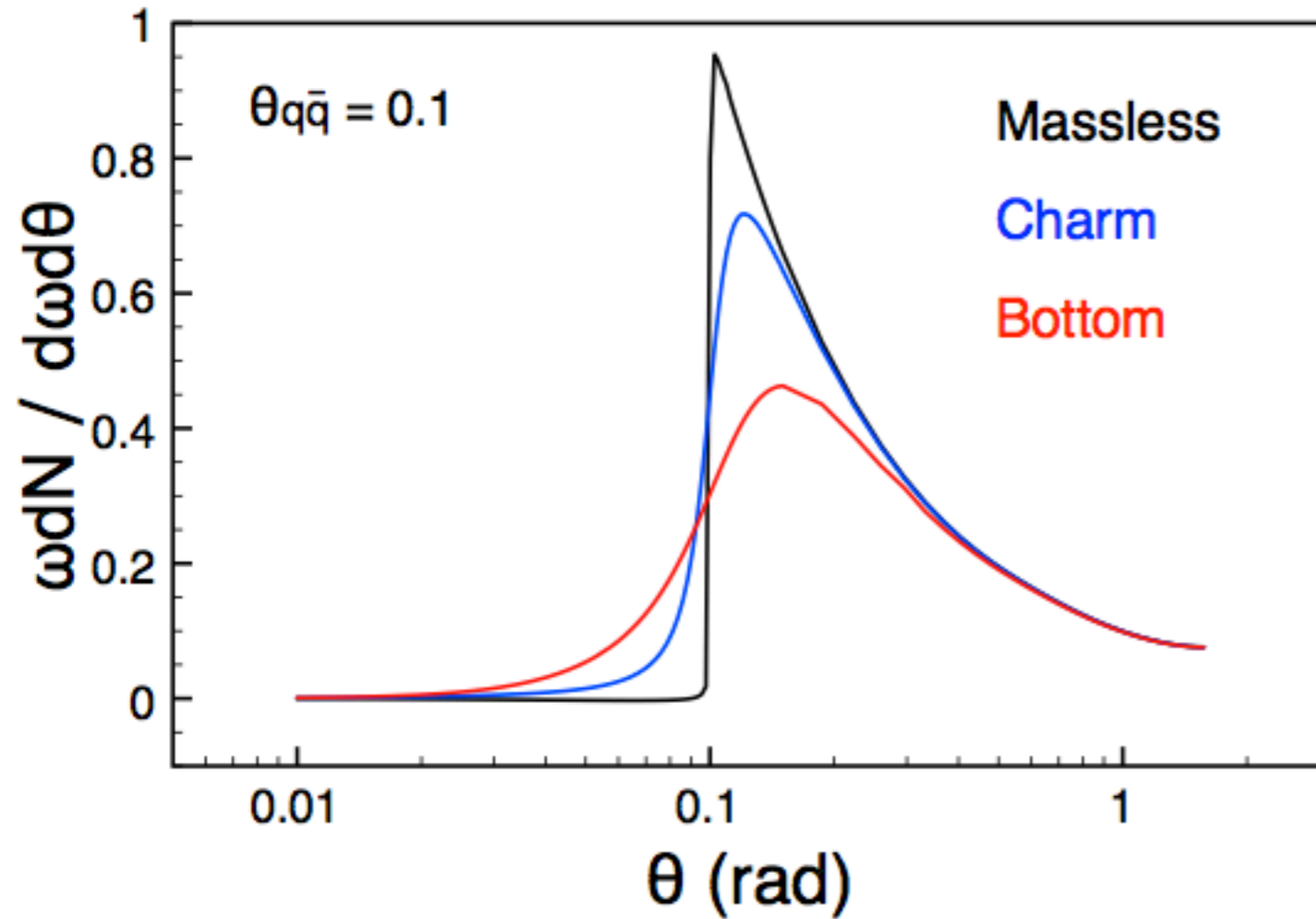
Medium decoherence parameter ($\mu_D \theta_{q\bar{q}} L_+ \ll 1$):

$$\omega \frac{dN}{d\omega d\theta} \propto \frac{1}{\theta} \Theta(\theta - \theta_{q\bar{q}}) \Delta_{\text{med}}(\theta_{q\bar{q}}, L_+)$$

$$\Delta_{\text{med}}(\theta_{q\bar{q}}, L_+) \propto \hat{q} L_+ |\mathbf{r}_\perp|^2 \sim \hat{q} L_+^3 \theta_{q\bar{q}}^2$$

Collinear convergence \Leftarrow **Antiangular ordering:** $\theta > \theta_{q\bar{q}}$

Medium-induced antiangular ordering in the soft limit



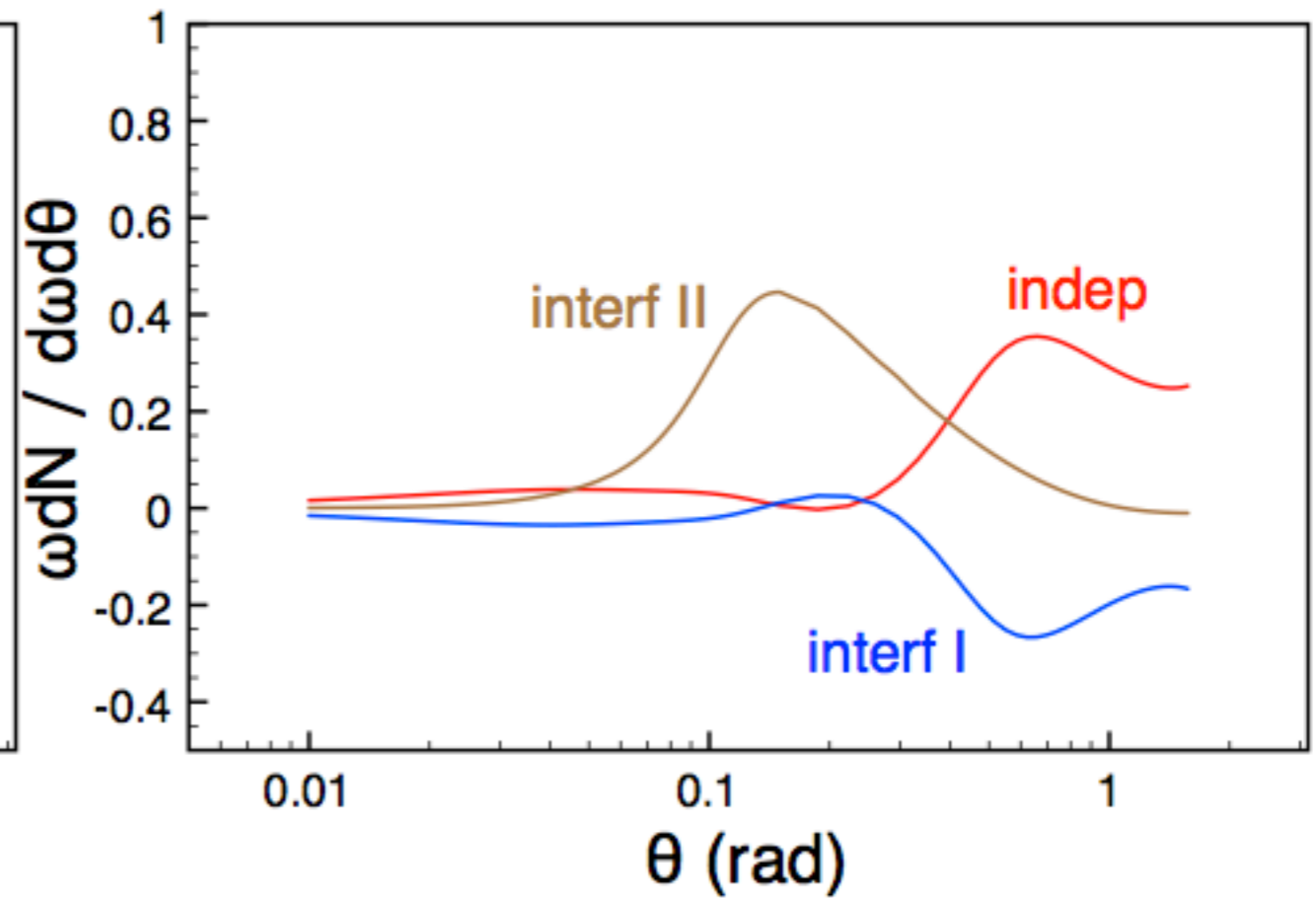
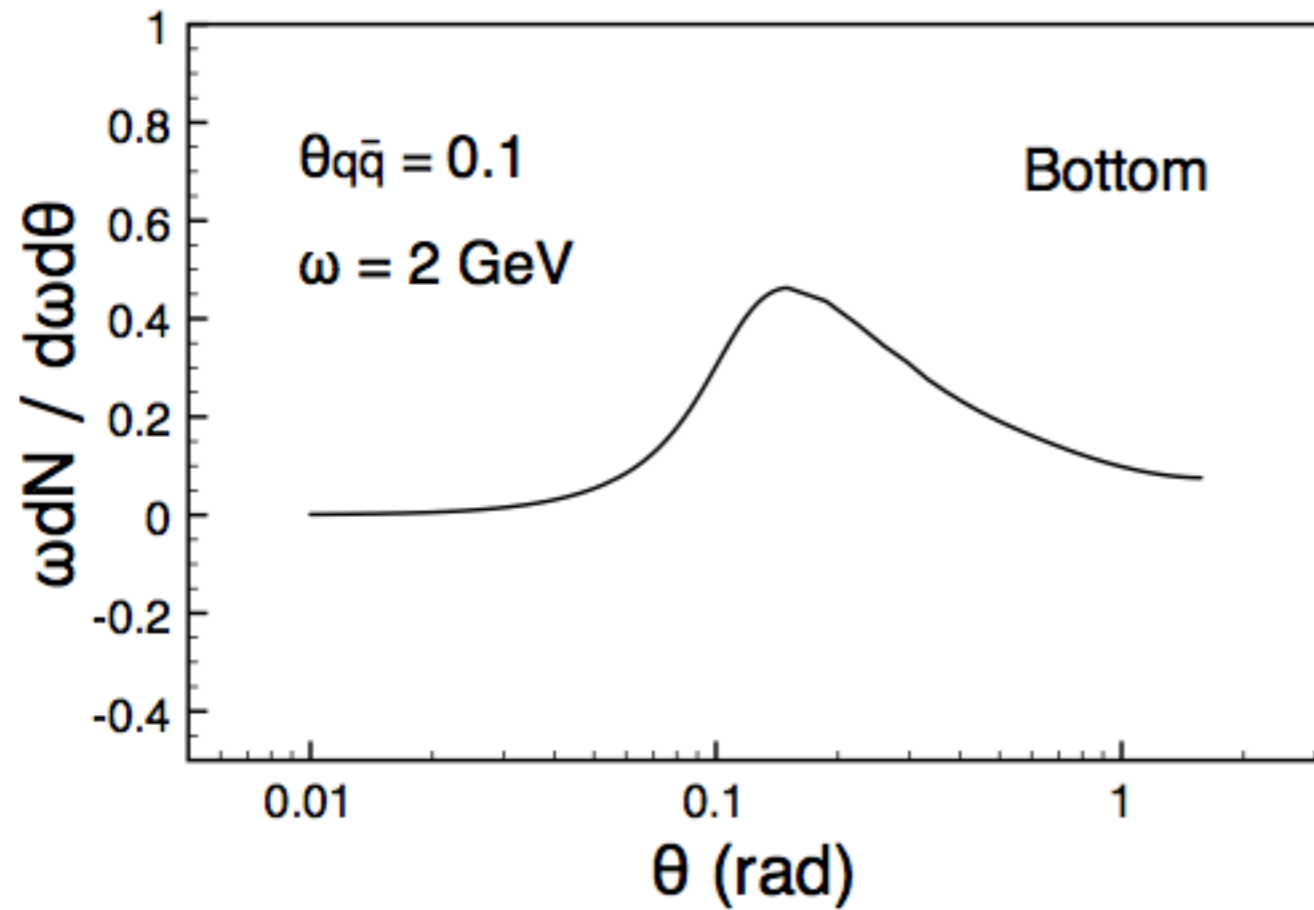
$$\omega \frac{dN}{d\omega d\theta} \propto \frac{\theta}{\theta^2 + \theta_0^2} H(\theta_{q\bar{q}}, \theta_0, \theta) \Delta_{\text{med}}(\theta_{q\bar{q}}, \theta_0, L_+)$$

Medium decoherence parameter

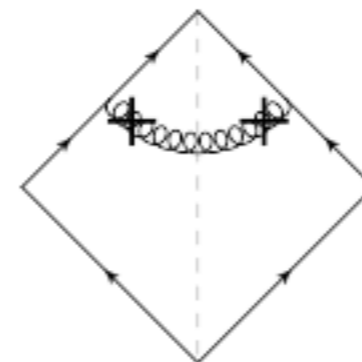
⇒ **Antiangular ordering** is modified by the dead cone angle θ_0 .

Cancellation in the soft limit extended to the massive case

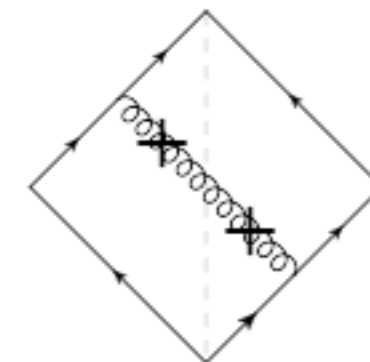
Massless case: Mehtar-Tani, Salgado and Tywoniuk, *Phys. Rev. Lett.* 106 (2011) 122002



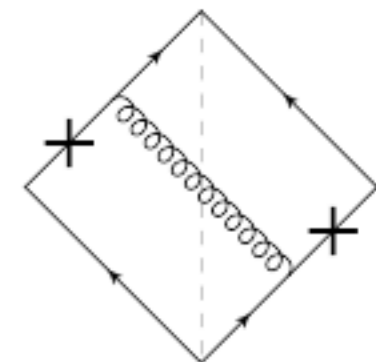
$$\omega \frac{dN}{d\omega d\theta} \sim \mathcal{I}_{q\bar{q}}^{\text{indep}} + \mathcal{I}_{q\bar{q}}^{\text{interf I}} + \mathcal{I}_{q\bar{q}}^{\text{interf II}}$$



independent

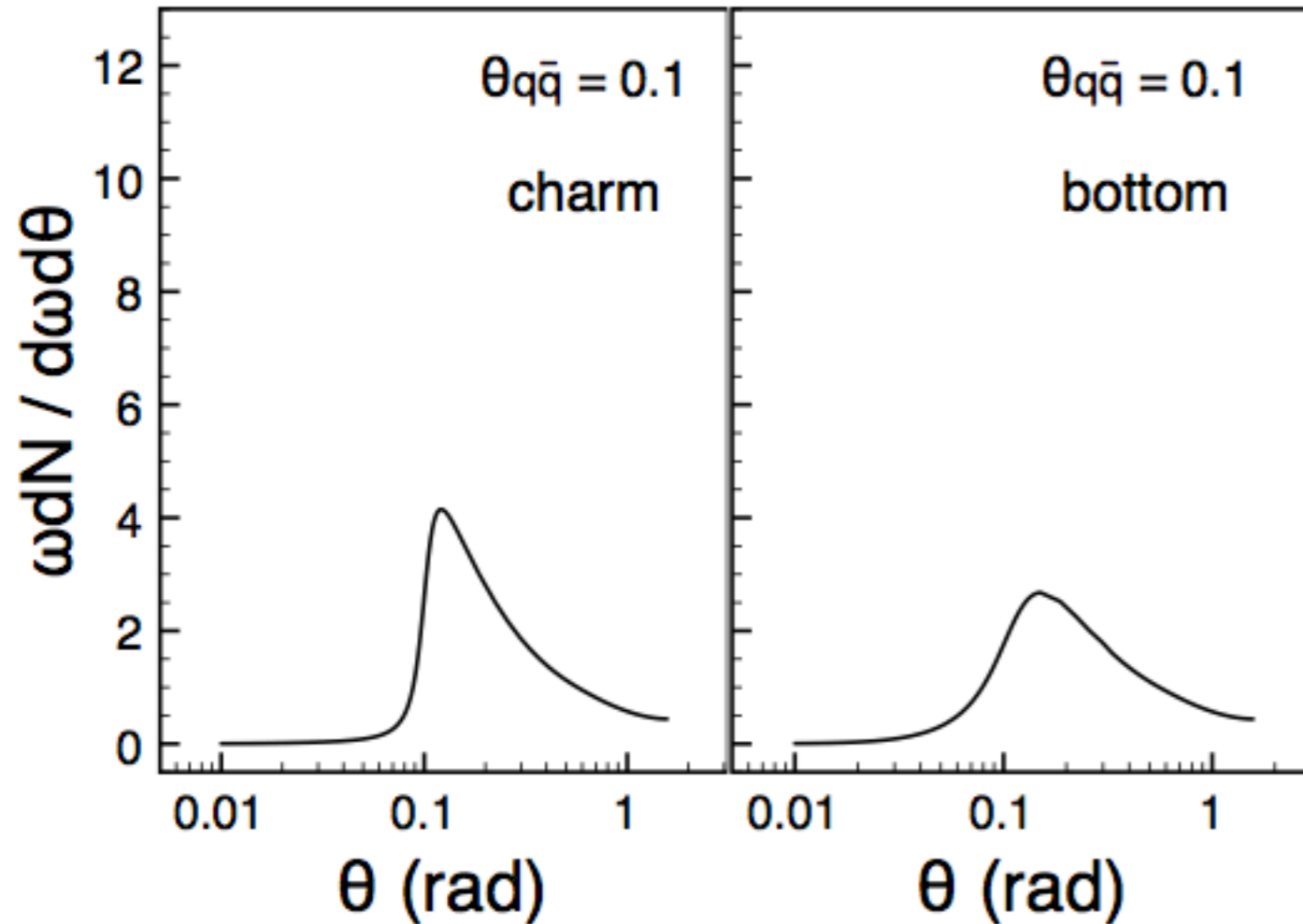


interference I



interference II

Decoherence in the soft limit extended to the massive case

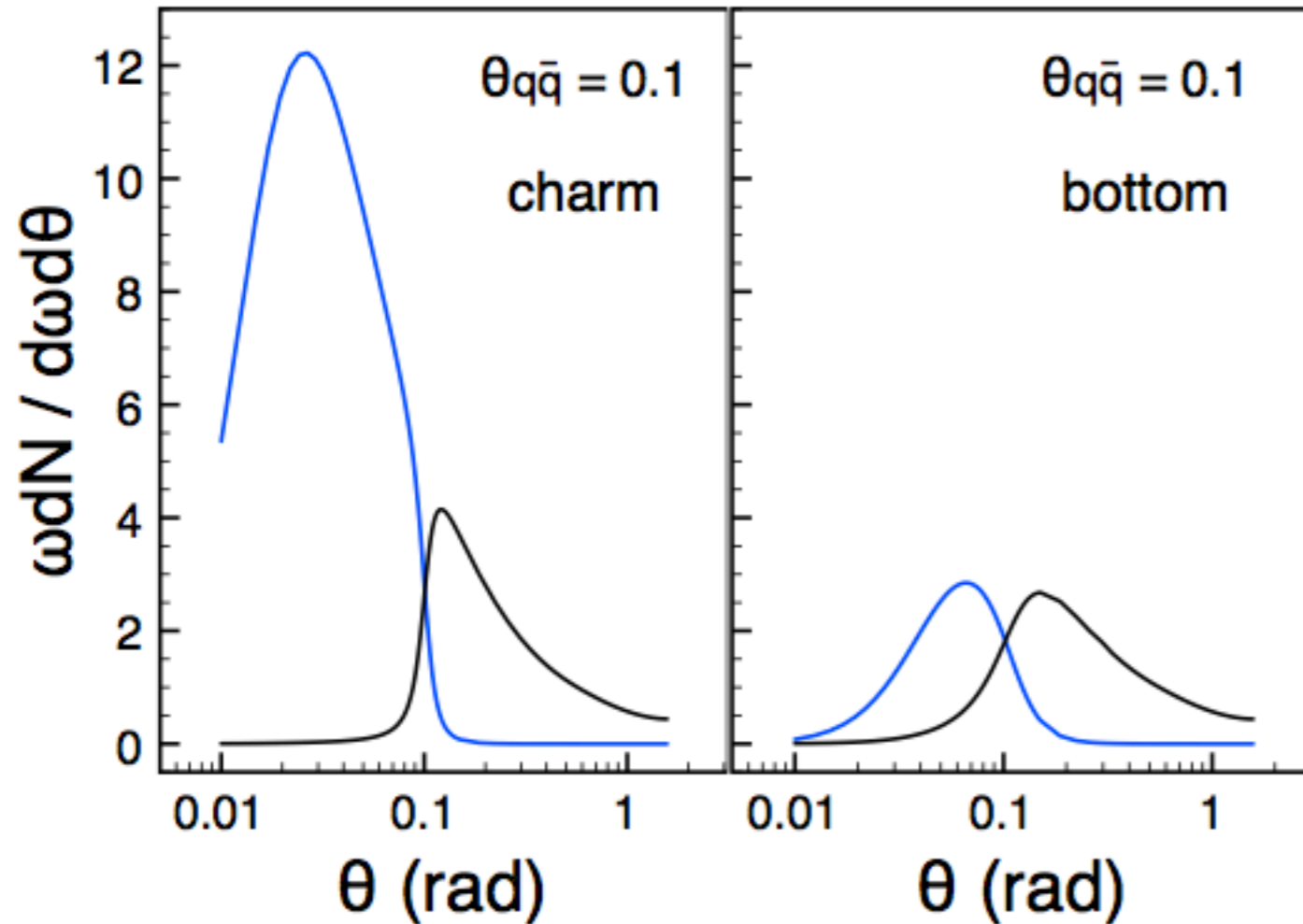


Medium-induced gluon radiation off a massive antenna in the opaque medium limit:

$$\omega \frac{dN}{d\omega d\theta}$$

[Massless \Rightarrow Antiangular ord.]

Decoherence in the soft limit extended to the massive case



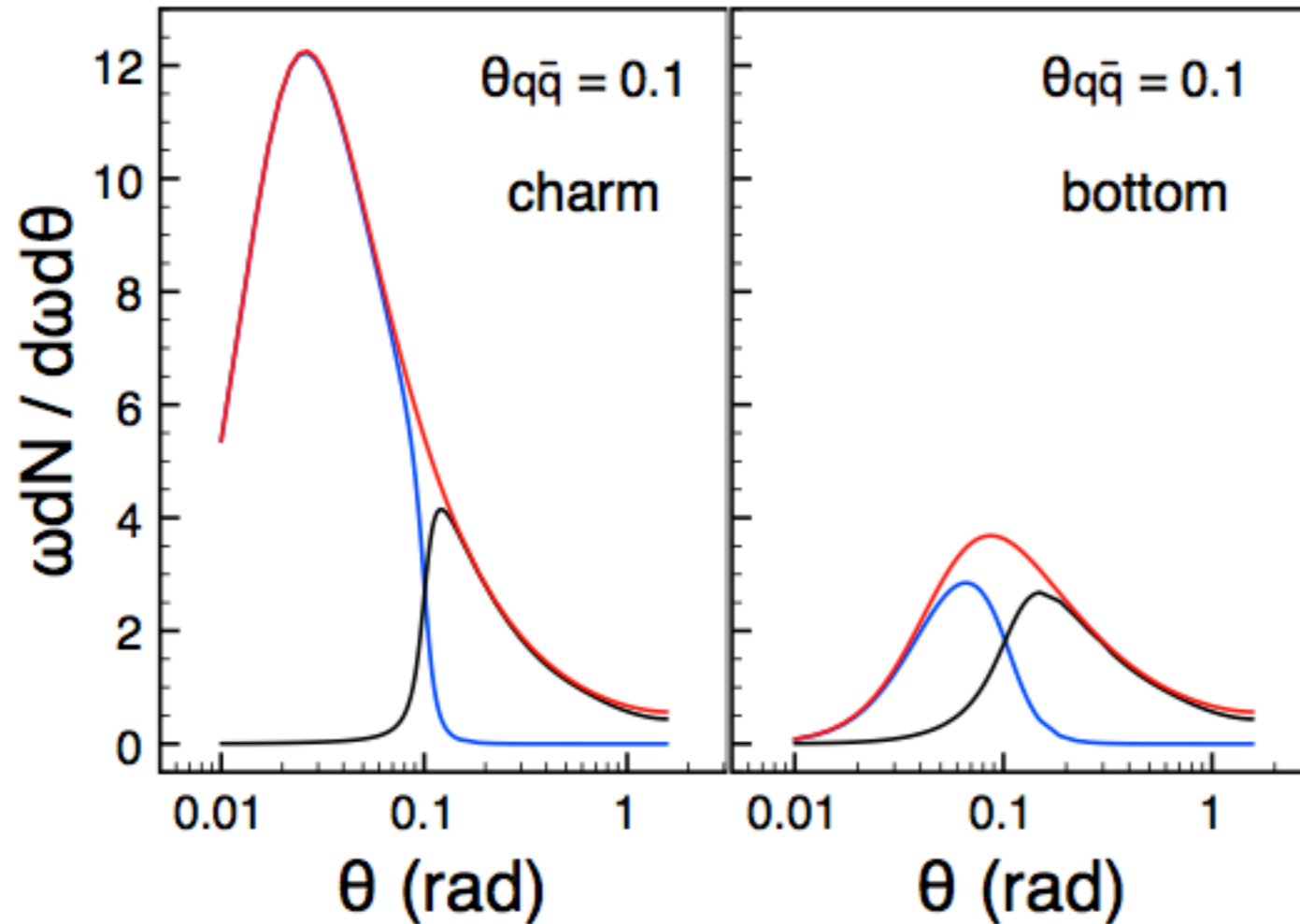
Medium-induced gluon radiation off a massive antenna in the opaque medium limit:

$$\omega \frac{dN}{d\omega d\theta} \quad [\text{Massless} \Rightarrow \text{Antiangular ord.}]$$

Gluon radiation off a massive antenna in vacuum:

$$\omega \frac{dN_{\text{vac}}}{d\omega d\theta} \quad [\text{Massless} \Rightarrow \text{Angular ord.}]$$

Decoherence in the soft limit extended to the massive case



Massless case: Mehtar-Tani, Salgado and Tywoniuk,
arXiv:1102.4317 [hep-ph]

$$\omega \frac{dN_{\text{vac}}^{\text{decoh}}}{d\omega d\theta} \approx \omega \frac{dN_{\text{vac}}}{d\omega d\theta} + \omega \frac{dN}{d\omega d\theta}$$

Medium-induced gluon radiation off a massive antenna in the opaque medium limit:

$$\omega \frac{dN}{d\omega d\theta} \quad [\text{Massless} \Rightarrow \text{Antiangular ord.}]$$

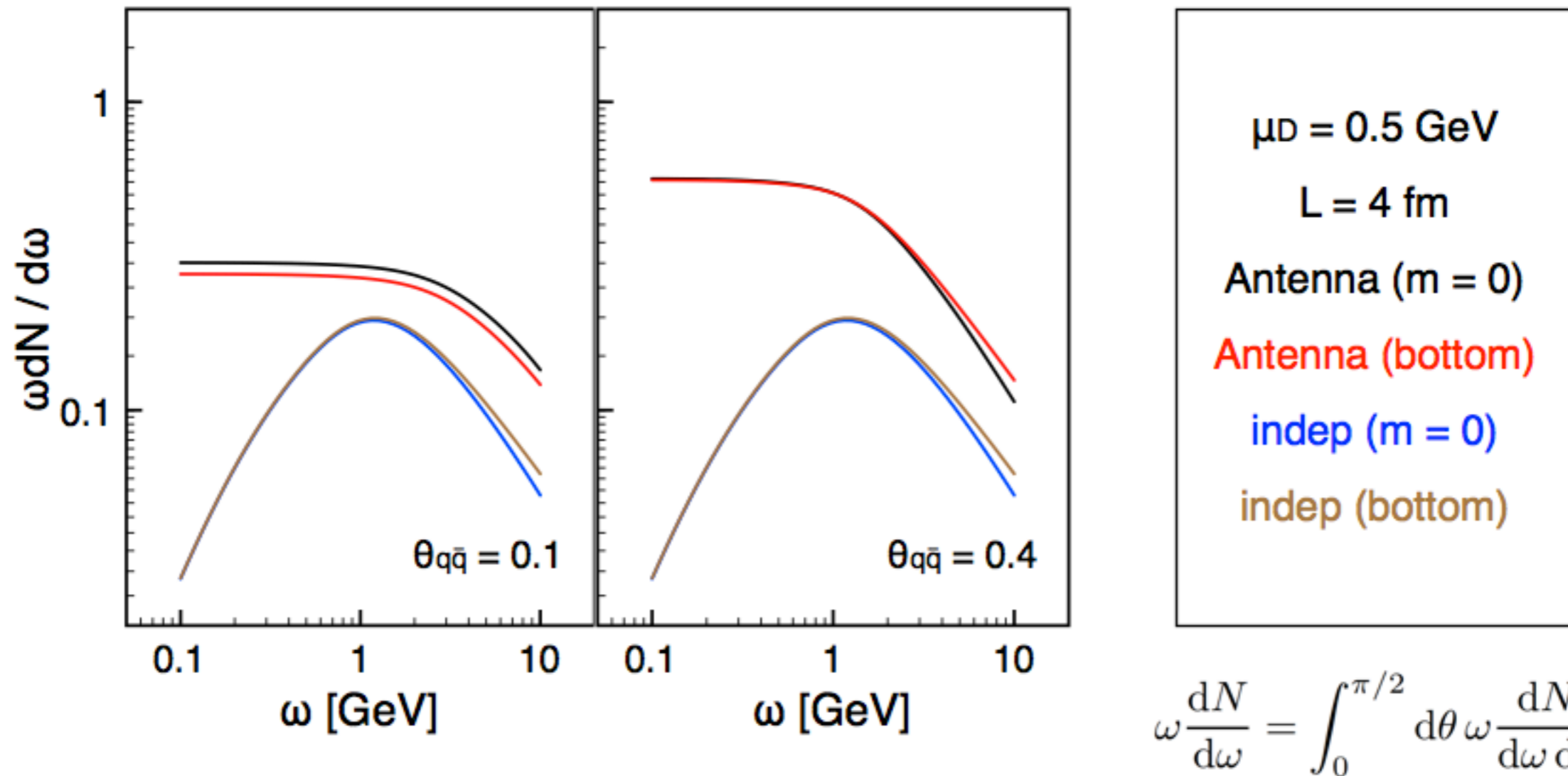
Gluon radiation off a massive antenna in vacuum:

$$\omega \frac{dN_{\text{vac}}}{d\omega d\theta} \quad [\text{Massless} \Rightarrow \text{Angular ord.}]$$

Gluon hard vacuum radiation multiplied by 2:

$$\omega \frac{dN_{\text{vac}}^{\text{decoh}}}{d\omega d\theta}$$

Medium-induced gluon energy spectrum

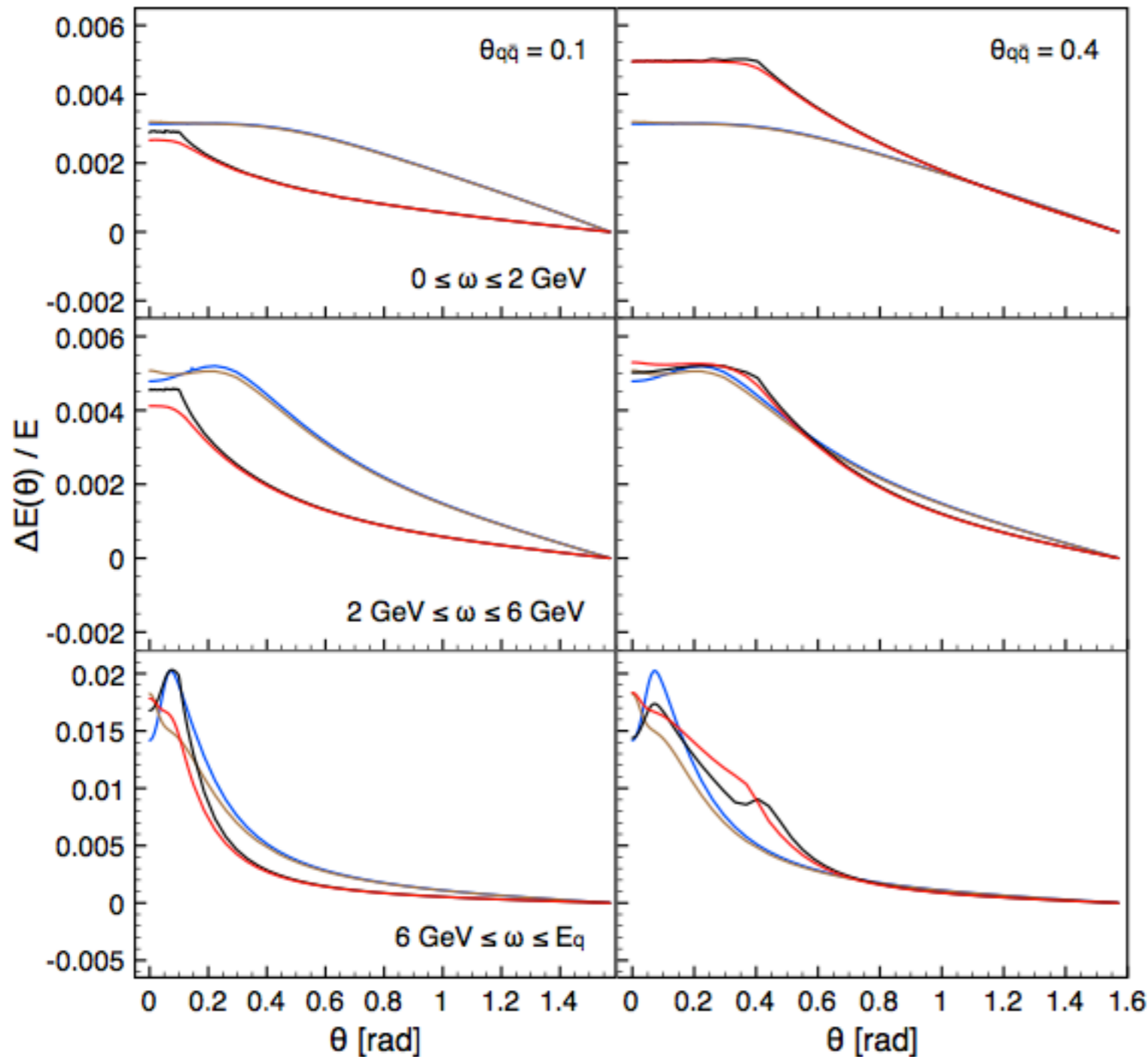


- Antenna opens phase space for soft gluon radiation at relatively large opening angles

⇒ indep is suppressed when $\omega \rightarrow 0$.

Cut-off scale: $\omega_{\text{coh}} \sim (\theta_{q\bar{q}}^2 L)^{-1}$

Angular dependence of gluon energy distribution

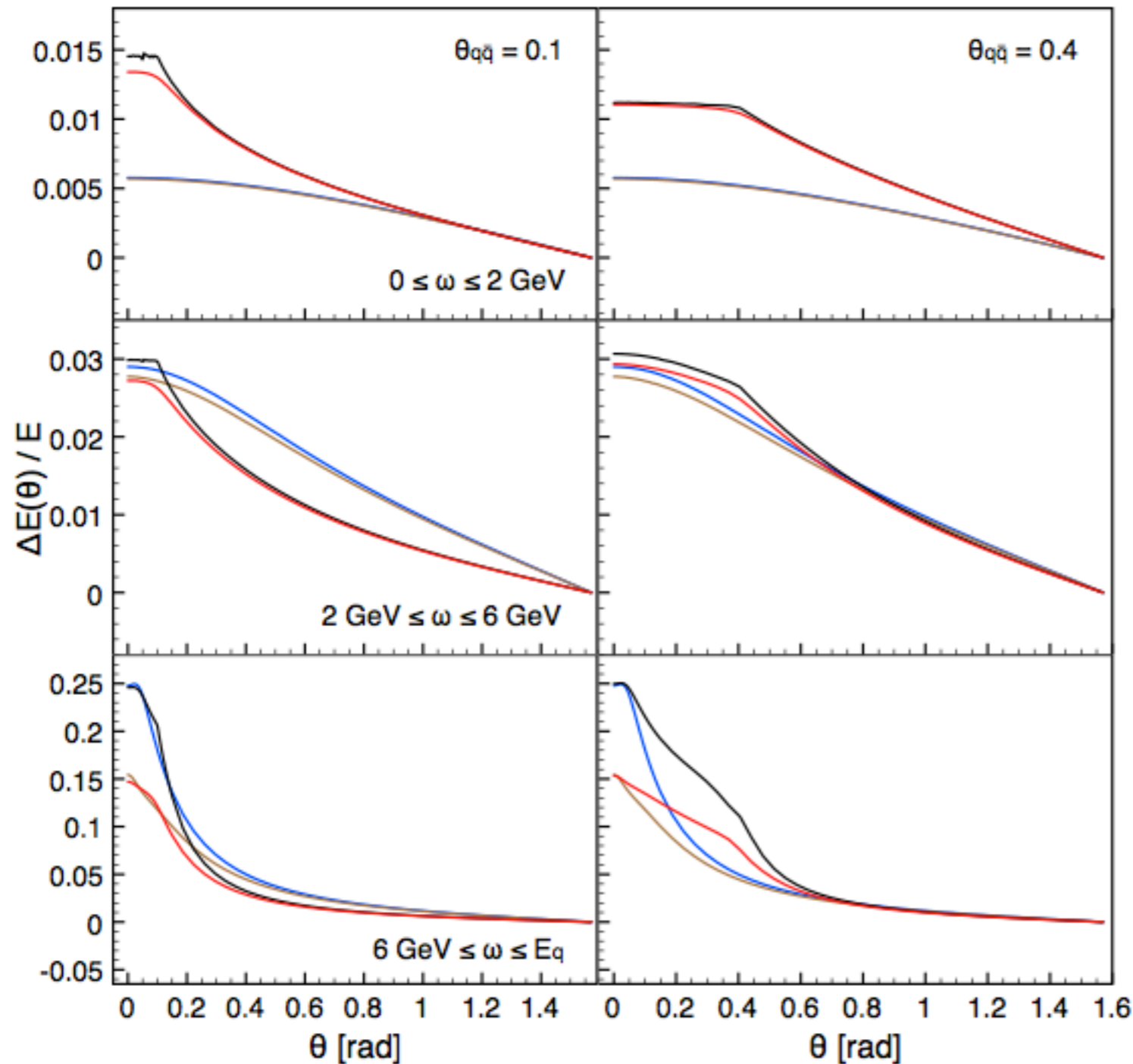


E = 100 GeV
 $\mu_D = 0.5$ GeV
L = 4 fm
1/2 Antenna (m = 0)
1/2 Antenna (bottom)
indep (m = 0)
indep (bottom)

$$\Delta E(\theta) = \int_{\omega_{\min}}^{\omega_{\max}} d\omega \int_{\theta}^{\pi/2} d\theta' \omega \frac{dN}{d\omega d\theta'}$$

\Rightarrow No \mathbf{k}_{\perp} -broadening in Antenna
 as in indep for the soft sector.

Angular dependence of gluon energy distribution

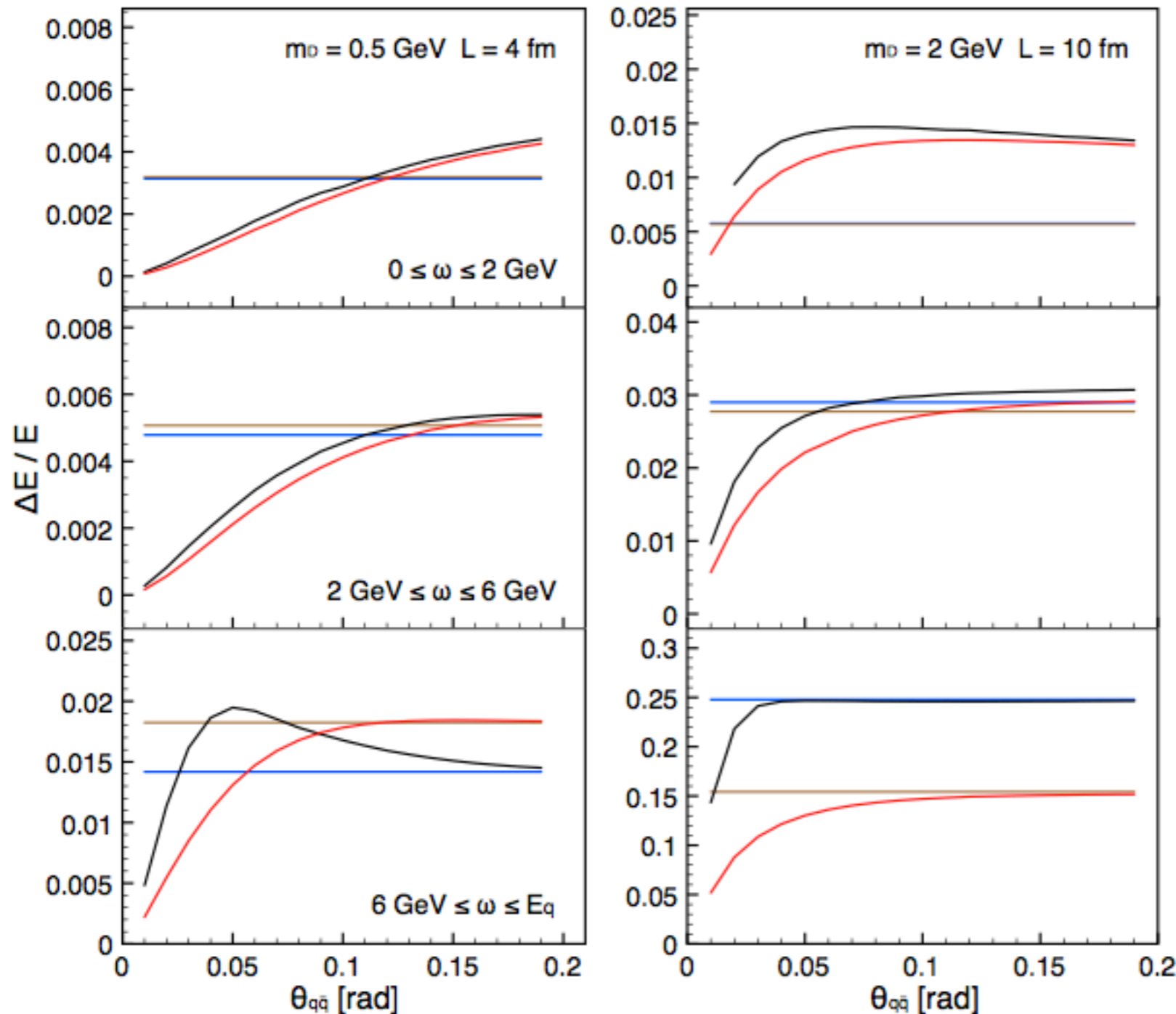


E = 100 GeV
 $\mu_D = 2 \text{ GeV}$
L = 10 fm
 $\frac{1}{2}$ Antenna (m = 0)
 $\frac{1}{2}$ Antenna (bottom)
indep (m = 0)
indep (bottom)

$$\Delta E(\theta) = \int_{\omega_{\min}}^{\omega_{\max}} d\omega \int_{\theta}^{\pi/2} d\theta' \omega \frac{dN}{d\omega d\theta'}$$

\Rightarrow No \mathbf{k}_{\perp} -broadening in Antenna
 as in indep for the soft sector.

Average energy loss



E = 100 GeV
1/2 Antenna (m = 0)
1/2 Antenna (bottom)
indep (m = 0)
indep (bottom)

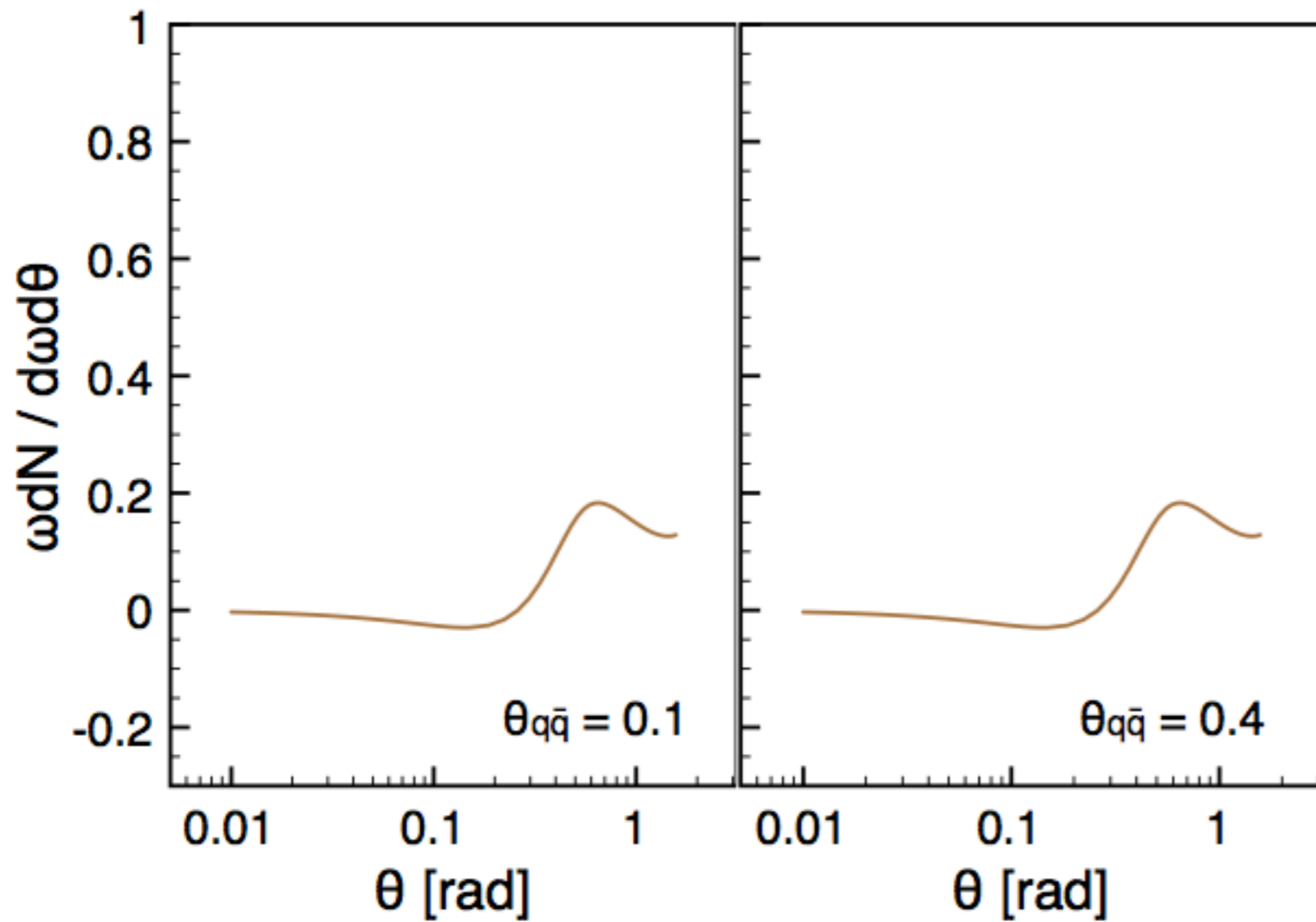
$$\Delta E = \int_{\omega_{\min}}^{\omega_{\max}} d\omega \int_0^{\pi/2} d\theta \omega \frac{dN}{d\omega d\theta}$$

⇒ More collimated jets lose less energy.

⇒ The size of the mass effect is similar to the indep.

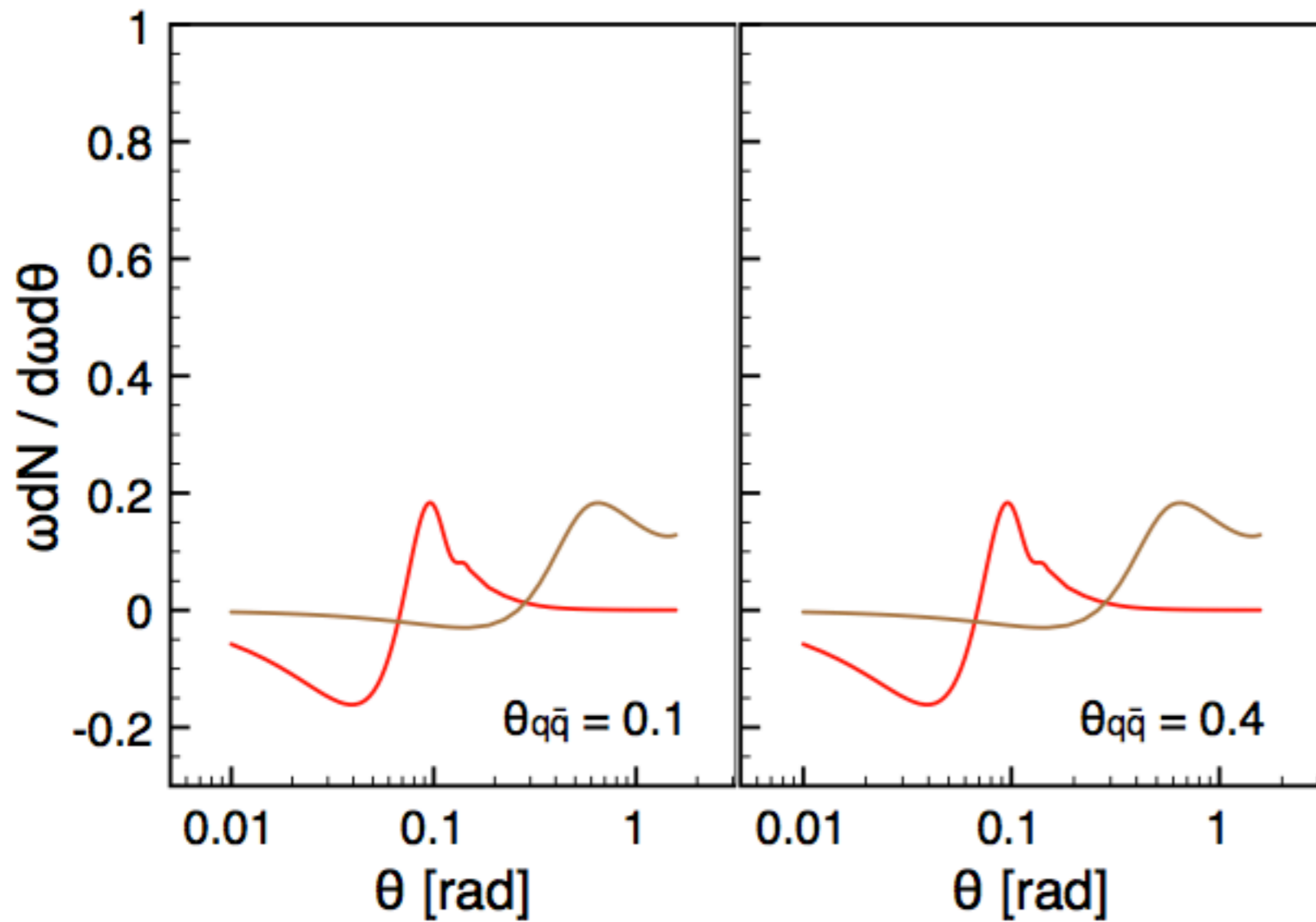
Armesto, Ma, Mehtar-Tani, Salgado and Tywoniuk, *arXiv:1107.0291v1 [hep-ph]*.

Transition between the massless antenna & indep



indep ($\omega = 2$ GeV)

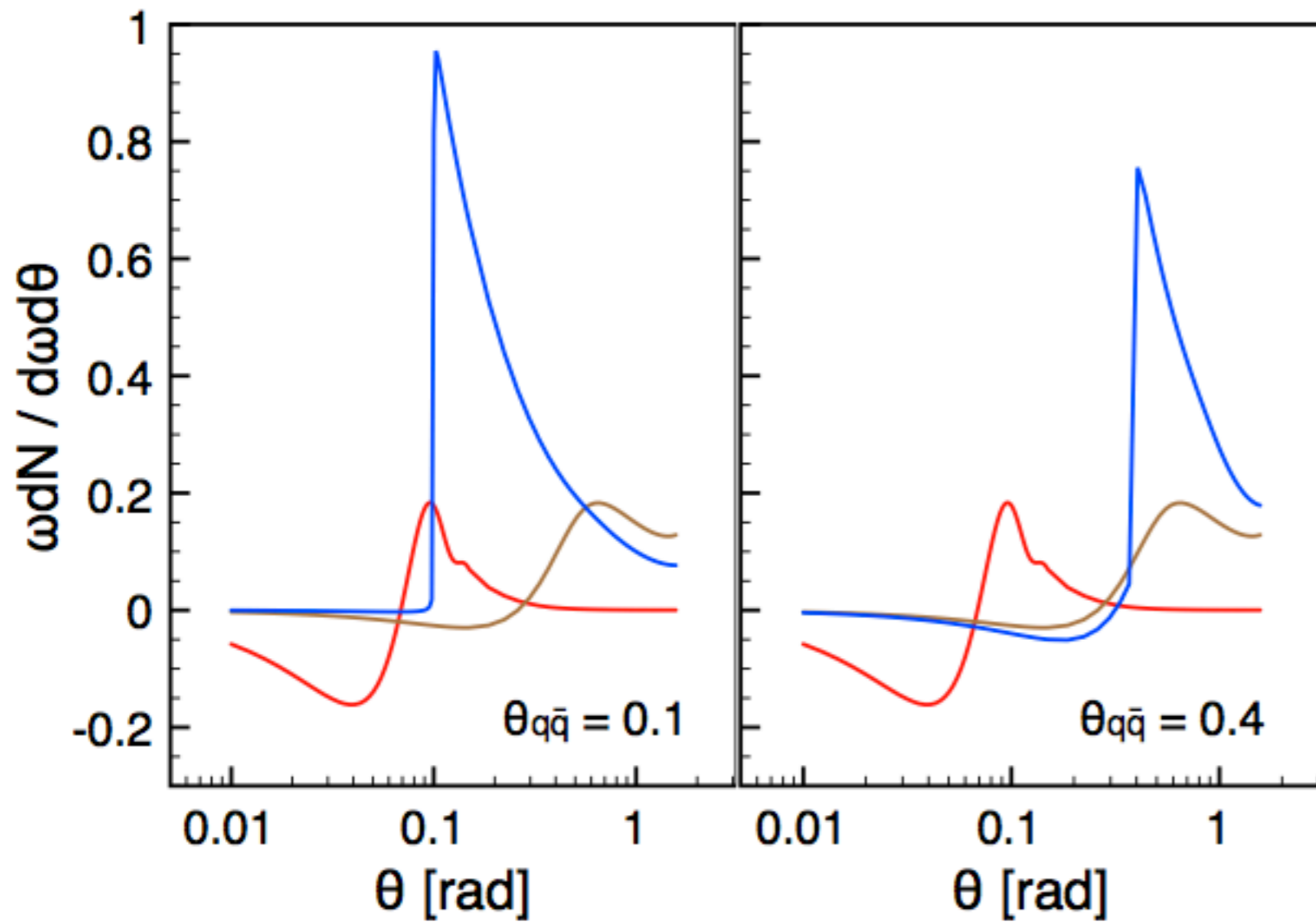
Transition between the massless antenna & indep



indep ($\omega = 2$ GeV)

indep ($\omega = 60$ GeV)

Transition between the massless antenna & indep

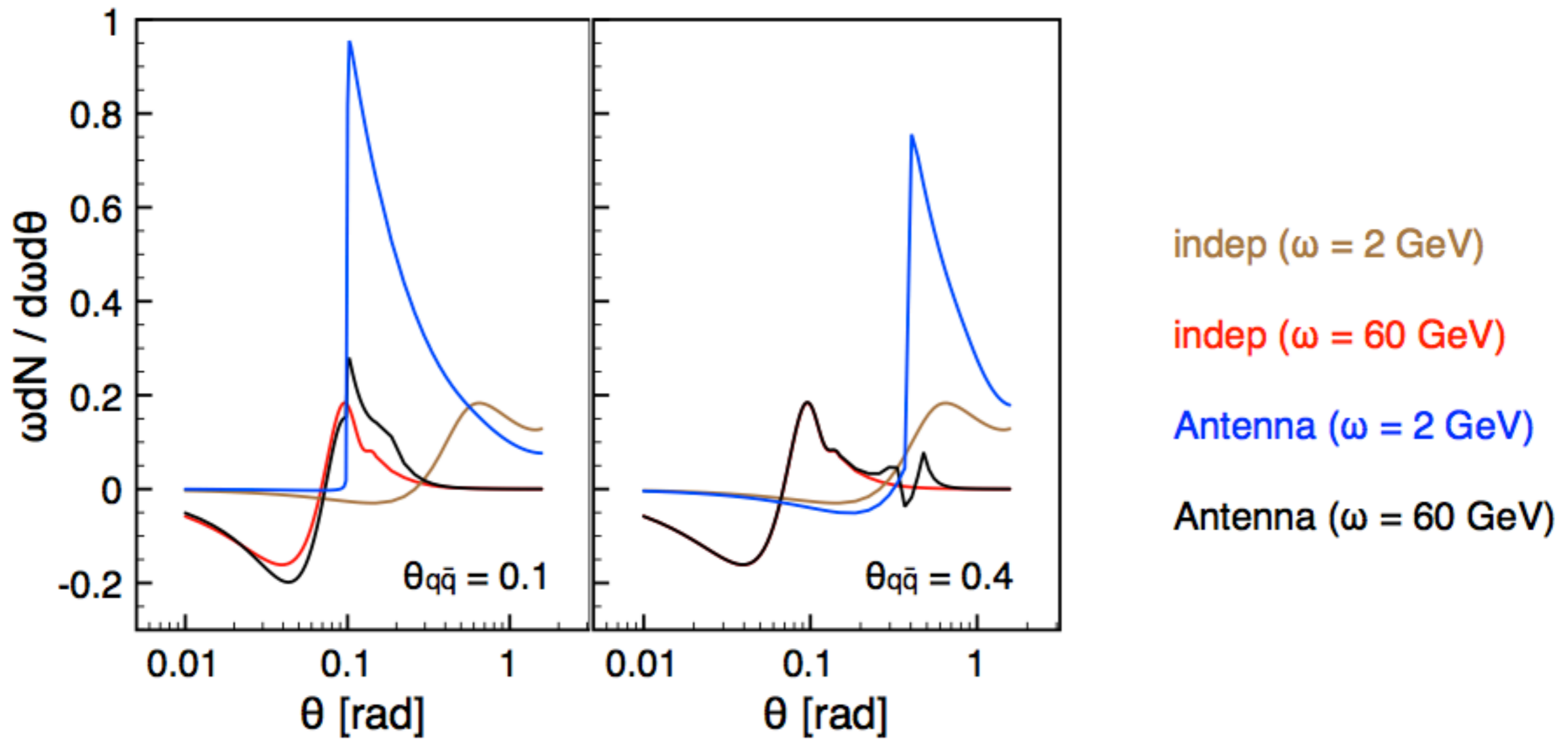


indep ($\omega = 2$ GeV)

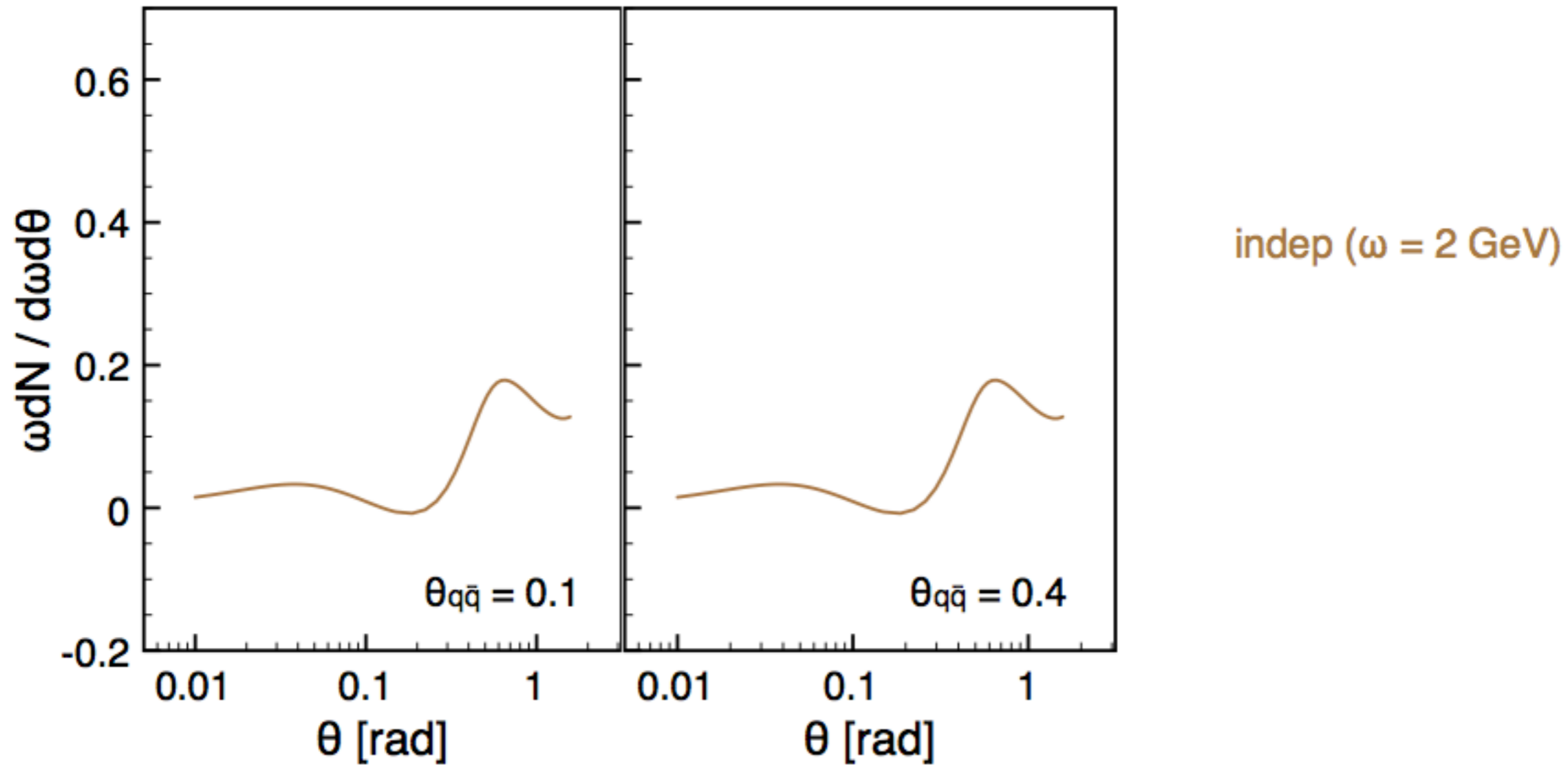
indep ($\omega = 60$ GeV)

Antenna ($\omega = 2$ GeV)

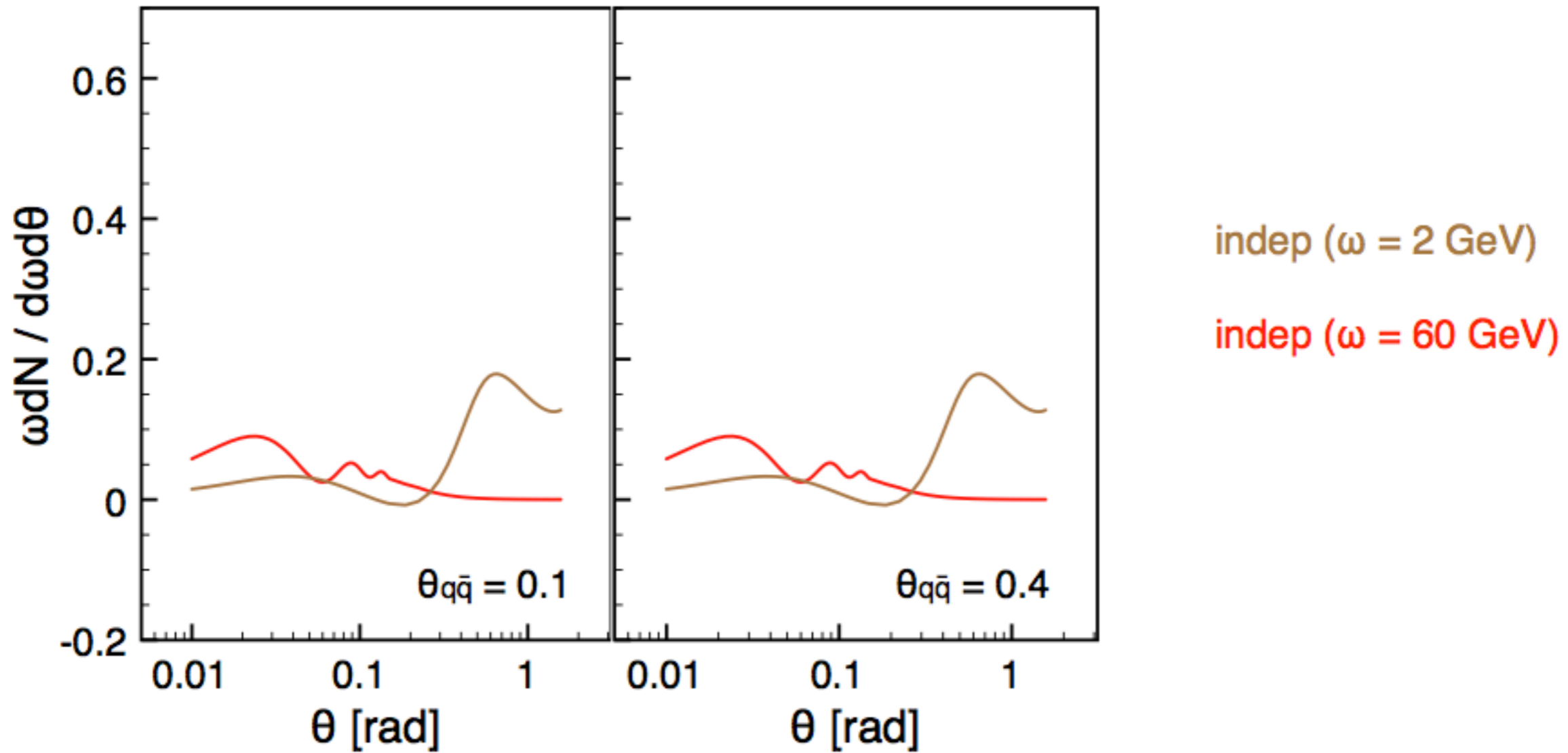
Transition between the massless antenna & indep



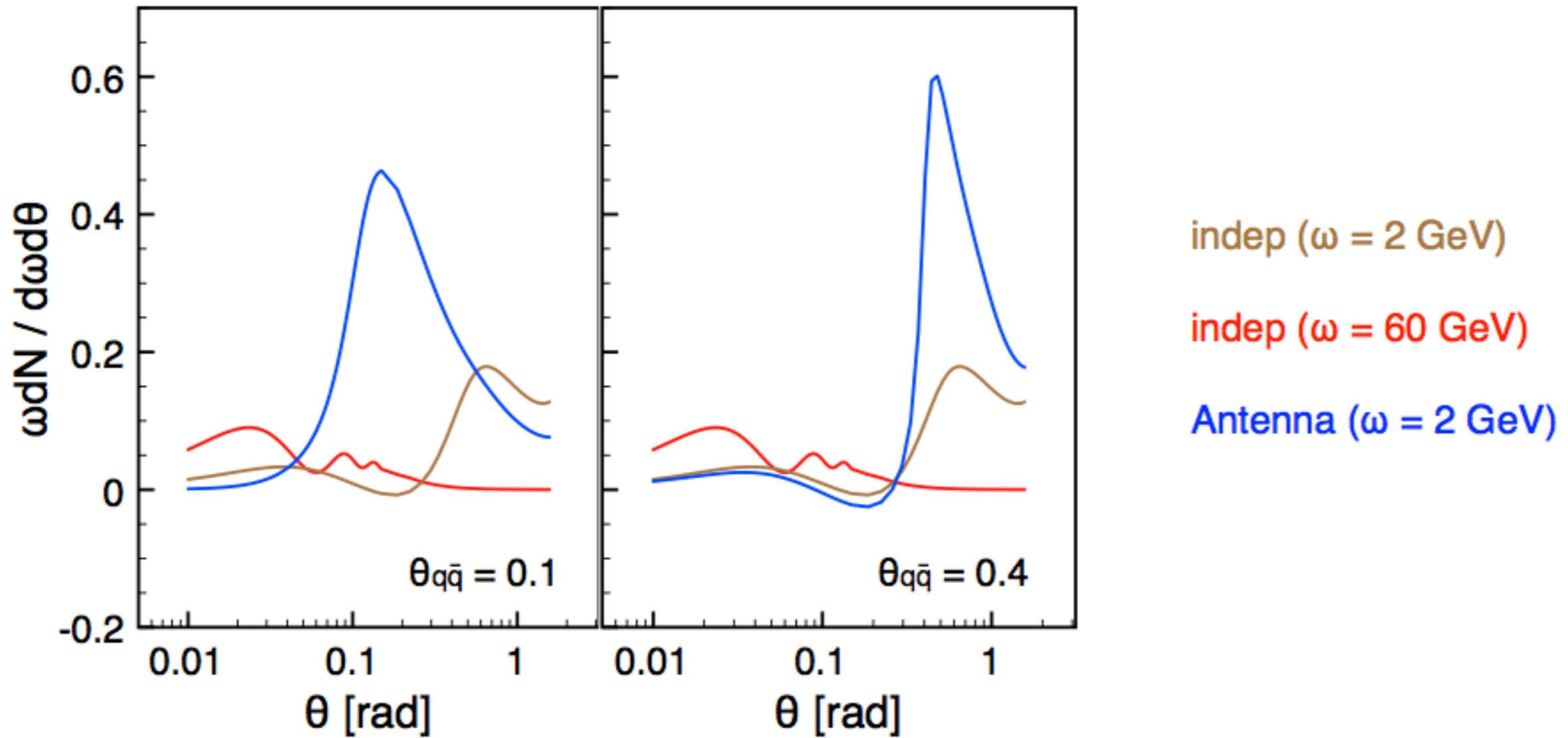
Transition between the bottom antenna & indep



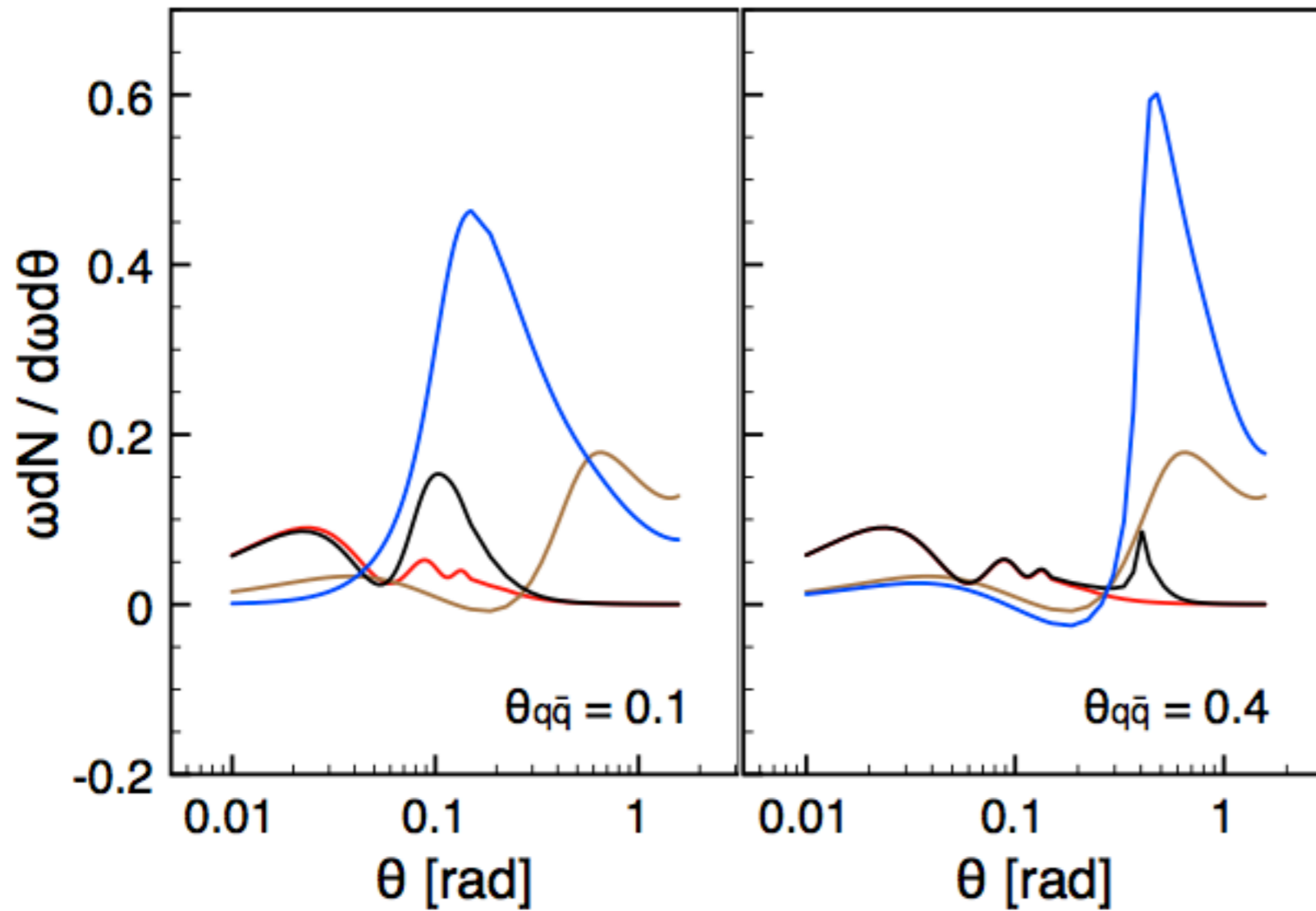
Transition between the bottom antenna & indep



Transition between the bottom antenna & indep



Transition between the bottom antenna & indep



indep ($\omega = 2$ GeV)

indep ($\omega = 60$ GeV)

Antenna ($\omega = 2$ GeV)

Antenna ($\omega = 60$ GeV)

Conclusion

- Medium-induced antiangular ordering in the soft limit is modified in the massive antenna case due to the dead cone effect.
- Decoherence in the soft limit is extended to the massive antenna case.
- Antenna opens phase space for soft gluon radiation at relatively large opening angles.
- Both dead cone effect and non-Abelian LPM effect appear in medium for the massive antenna case.
- More collimated jets lose less energy.
- The size of the mass effect of the antenna is similar to the independent emitter.

Back-up slides

Back-up

- Massless medium decoh. parameter

$$\Delta_{\text{med}}(\theta_{q\bar{q}}, L_+) = \frac{\hat{q}}{\mu_D^2} \int_0^{L_+} dx_+ \left(1 - \frac{|\mathbf{r}_\perp| \mu_D x_+}{L_+} K_1 \left(\frac{|\mathbf{r}_\perp| \mu_D x_+}{L_+} \right) \right)$$

- Massless dipole size $|\mathbf{r}_\perp| = \frac{\sqrt{2} \sin \theta_{q\bar{q}} L_+}{1 + \cos \theta_{q\bar{q}}}$

- Massive medium decoh. parameter

$$\Delta_{\text{med}}(\theta_{q\bar{q}}, \theta_0, L_+) = \frac{\hat{q}}{\mu_D^2} \int_0^{L_+} dx_+ \left(1 - \frac{|\mathbf{r}_\perp| \mu_D x_+}{L_+} K_1 \left(\frac{|\mathbf{r}_\perp| \mu_D x_+}{L_+} \right) \right)$$

- Massive dipole size $|\mathbf{r}_\perp| = \frac{\sqrt{2} \sin \theta_{q\bar{q}} L_+}{1 + \sqrt{1 - \theta_0^2} \cos \theta_{q\bar{q}}}$

- Gluon formation time $t_g^{\text{form}} \sim \frac{\omega}{|\mathbf{k}_\perp|^2 + \theta_0^2 \omega^2}$

Back-up

- independent emitter spectrum

$$\mathcal{I}_{quark}^{indep} = \int \frac{d^2\mathbf{q}}{(2\pi)^2} \int_0^{L^+} dx^+ n_0 \frac{m_D^2}{(\mathbf{q}^2 + m_D^2)^2} 8\alpha_s^2 (4\pi)^2 C_A C_F \frac{(\mathbf{k} - \mathbf{q})^2 \mathbf{k} \cdot \mathbf{q} - x^2 m_q^2 (\mathbf{k} - \mathbf{q}) \cdot \mathbf{q}}{[(\mathbf{k} - \mathbf{q})^2 + x^2 m_q^2]^2 (\mathbf{k}^2 + x^2 m_q^2)} \left[1 - \cos\left(\frac{(\mathbf{k} - \mathbf{q})^2 + x^2 m_q^2}{2k^+} x^+\right) \right]$$

Back-up

- Antenna spectrum in the soft limit

$$\mathcal{I}_{q\bar{q}}^{\text{interf II}} = \int \frac{d^2\mathbf{q}_\perp}{(2\pi)^2} \int_0^{L_+} dx_+ n_0 \frac{\mu_D^2}{(\mathbf{q}_\perp^2 + \mu_D^2)^2} (-2) \alpha_s^2 (4\pi)^2 C_A C_F \frac{1}{x\bar{x}} \frac{\boldsymbol{\kappa}_\perp \cdot \bar{\boldsymbol{\kappa}}_\perp}{p \cdot k \bar{p} \cdot k} \left(\cos\left(\frac{p \cdot v}{p_+} x_+ - \frac{\bar{p} \cdot v}{\bar{p}_+} x_+\right) - \cos\left(\frac{p \cdot k}{p_+} x_+ - \frac{\bar{p} \cdot k}{\bar{p}_+} x_+\right) \right)$$



$$\mathcal{I}_{q\bar{q}}^{\text{interf II}} = \int \frac{d^2\mathbf{q}_\perp}{(2\pi)^2} \int_0^{L_+} dx_+ n_0 \frac{\mu_D^2}{(\mathbf{q}_\perp^2 + \mu_D^2)^2} 2 \alpha_s^2 (4\pi)^2 C_A C_F \frac{1}{x\bar{x}} \frac{\boldsymbol{\kappa}_\perp \cdot \bar{\boldsymbol{\kappa}}_\perp}{p \cdot k \bar{p} \cdot k} \cos(\Omega_{q\bar{q}}^0 x_+) (1 - \cos(\delta\mathbf{n}_\perp \cdot \mathbf{q}_\perp x_+))$$