

Positronium Physics

D. J. Murtagh

For the AVA School on Precision Studies

24th March 2020

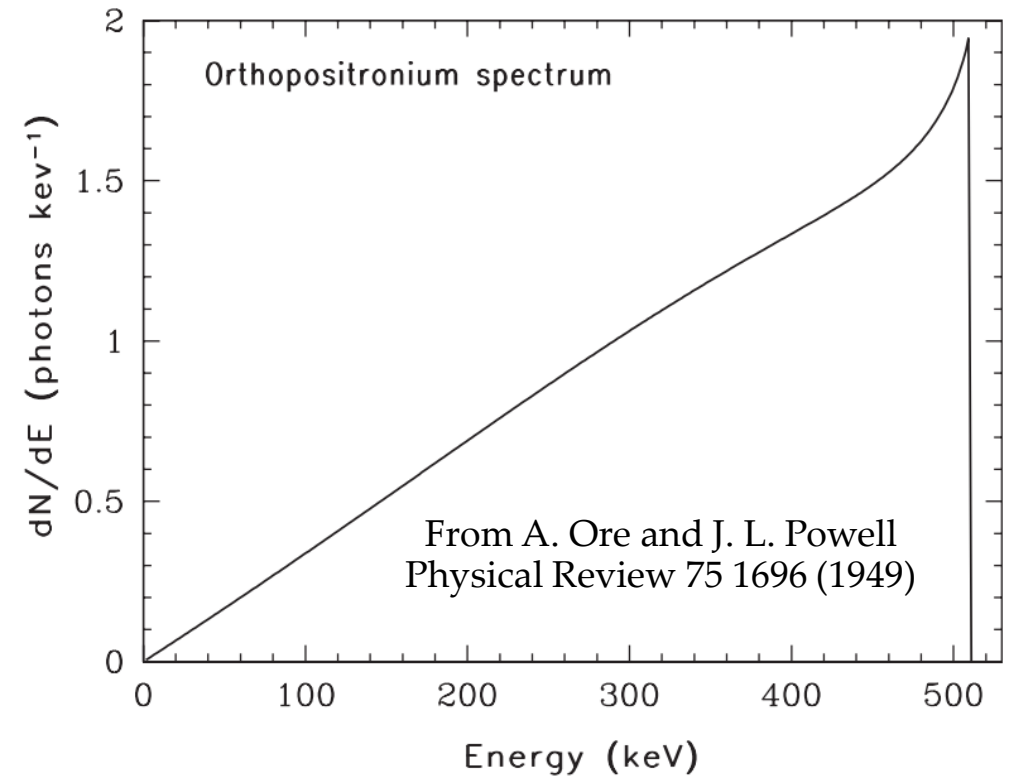


Overview

- In this lecture, I will introduce the field of Positronium physics from a historical perspective
- I will cover what you need to gain an understanding of positronium enabling you to go forward and explore this field for yourself
- At the end I will link to reviews which you can use as resources going forward

Start with the end

- Positronium is the end point of the life of most positrons in vacuua, the direct annihilation cross-section is generally considered to be negligible ¹
- Observations of gamma rays around positron annihilation peak (511 keV) in the interstellar medium find a positronium fraction of ~ 0.90 ²
- For some perspective approximately 10 billion kg of positrons annihilate in the Milky Way every second 90% of which is via Ps formation ³



1. Van Reeth, P. Laricchia, G. and Humberston., J. W *Physica Scripta*, **71** C9 (2005).
 2. T. Siegert, R. Diehl, G. Khachatryan, M. G. H. Krause, F. Guglielmetti, J. Greiner, A. W. Strong and X. Zhang *Astronomy and Astrophysics* **586** A84 (2016)
 3. Prantzos, N., Boehm, C., Bykov, A.M., Diehl, R., Ferrière, K., Guessoum, N., Jean, P., Knoedlseder, J., Marcowith, A., Moskalenko, I.V., Strong, A., Weidenspointner, G., *Rev. Mod. Phys.* **83**, 1001 (2011)

Prediction

- Dirac (1930) and Weyl (1931) predict that a particle with the same mass as the electron but with a positive charge should exist
- Anderson (1932a,b 1933) observes said positive electron in a cloud chamber
- Blackett and Occhialini (1933) confirm the charge to mass ratio
- Mohorovičić (1934) predicted the existence of a bound state of the positive and negative electron

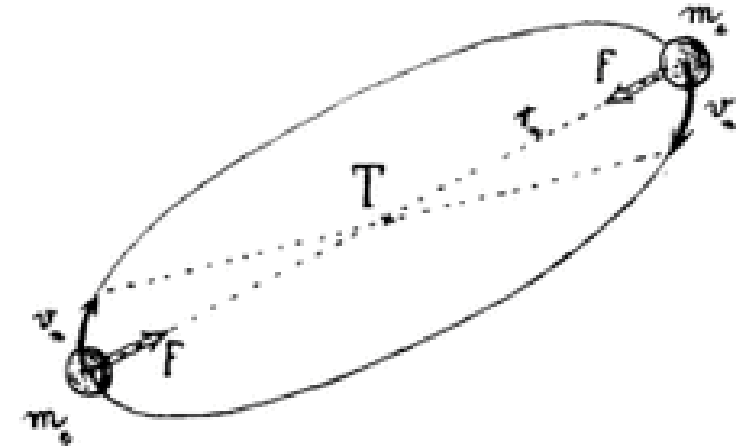


Fig. 1.

1. Dirac P. A. M. *Proc. Roy. Soc* **126** 360 (1930a)
2. Weyl R. N. *Gruppentheorie and Quantenmechanik* **2nd ed.** P234 (1931)
3. Anderson C. D. *Phys Rev* **41** 405 (1932a)
4. Anderson C. D. *Science* **76** 238 (1932b)
5. Anderson C. D. *Phys Rev* **43** 491 (1933)
6. Blackett P. M. S. and Occhialini G. P. *S Proc. Roy. Soc A* **139** 699 (1933)
7. Mohorovicic S. *Astrn. Nachr.* **235** 94 (1934)

Electrum ~~Positronium~~ Physics

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Properties of Electrum

- Rydberg 'constant' $\frac{M}{M+m} R_{\infty} = \frac{1}{2} R_{\infty}$
- Broad structure should follow $\frac{1}{\lambda} = \frac{1}{2} R_{\infty} \left(\frac{1}{n^2} - \frac{1}{k^2} \right)$
- Ground state binding energy or 6.8 eV
 - Half that of hydrogen which follows directly from the reduced mass

Prediction

- The spectroscopic structure was calculated by Ruark (1945)
- The binding energy and lifetime by Wheeler (1946) also considered were the positive and negative ions of Ps and Ps_2 - P^{+-} P^{++-} P^{--+} and P^{++--} in his own notation

Positronium

ARTHUR E. RUARK*

*Naval Research Laboratory, Office of Research and Inventions,
Anacostia, D. C.***

November 13, 1945

POLYELECTRONS*

BY

JOHN ARCHIBALD WHEELER†

1. Ruark A. E. *Phys. Rev* **68** 278 (1945)
2. Wheeler J. A. *Ann. N. Y. Acad. Sci.* **48** 291 (1946)

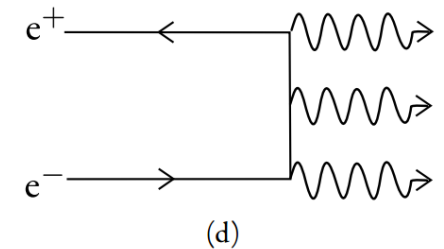
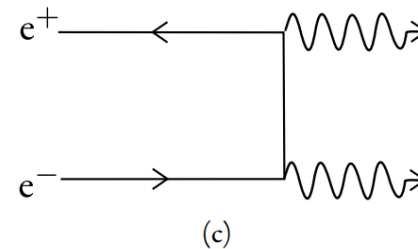
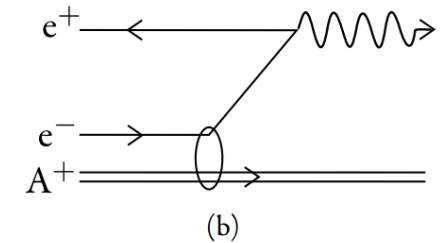
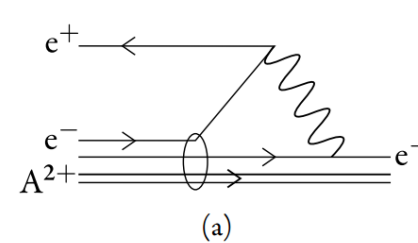


Properties of Positronium

- Lyman – α of 243 nm
- Blamer
 - α ($3 \rightarrow 2$) of 1312.6 nm
 - β ($4 \rightarrow 2$) of 972.2 nm
 - δ ($5 \rightarrow 2$) of 868.0 nm
 - Limit of 729 nm
- Bohr radius is twice that of hydrogen
- Due to the larger magnetic moment of the positron than the proton the structure has a triplet and singlet state
- It has a finite lifetime against annihilation which should be calculated
- The existence of annihilation channels will cause broadening of states

Properties of Positronium

- a^1 and b^2 show photonless and one photon annihilation, these require the presence of a third body in this case an atom
- c^3 and d^4 show two and three gamma annihilation
- Annihilation probability proportional to α^m where α is the fine structure constant and m the number of photons



1. Palathingal J. C., Asoka-Kumar P., Lynn K. G., Posada and Wu X. Y., Phys Rev. Lett. **67** 3491 (1991)
2. Klemperer O., Proc. Camb. Philos. Soc. **30** 347 (1934)
3. Chang T., Li Q., Wang Y and Li Y. in "Positron Annihilation" eds. P. G. Coleman, S. C. Sharman and L. M. Diana (North Holland: Amsterdam) p32 (1982)
4. Adachi S., Chiba M., Hirose T., Hagayama S., Nakamitsi Y., Sato T. and Yamada T. Phys. Rev. Lett **65** 2634 (1990)



Properties of Positronium

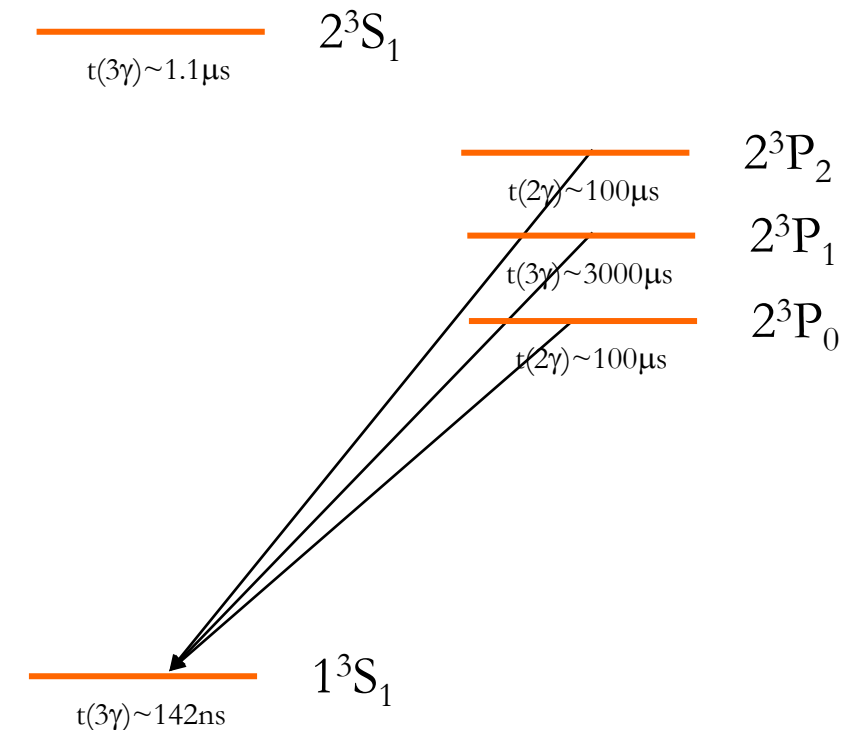
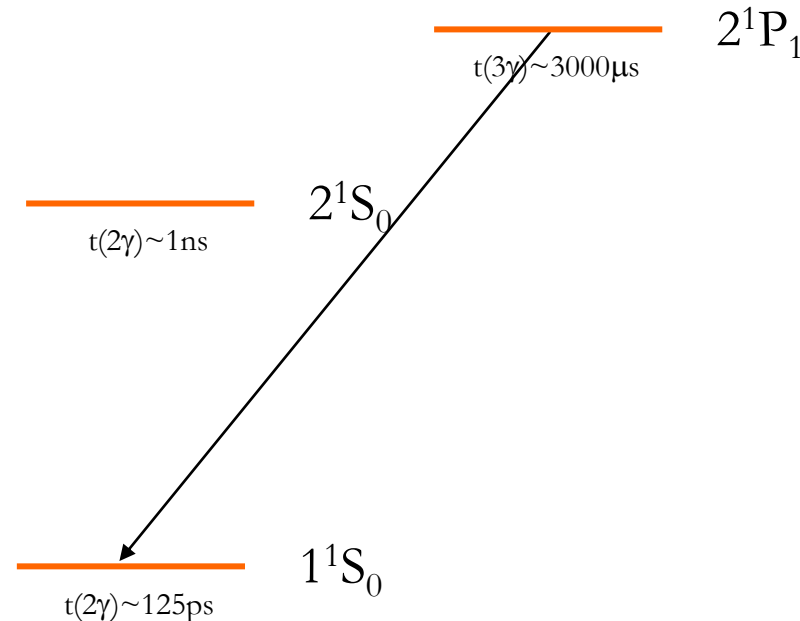
- Ground state lifetime
 - Singlet state : 125 ps
 - Triplet state 142 ns
- The number of annihilation quanta is determined by charge parity conservation
 - Photon $P_c = (-1)^n$
 - Ps system $P_c = (-1)^{L+S}$
 - $(-1)^n = (-1)^{L+S}$

$$\Gamma_{3\gamma}(n_{Ps}, {}^3S_1) = \frac{2}{9\pi} (\pi^2 - 9) \frac{mc^2}{\hbar} \frac{\alpha^6}{n_{Ps}^3}$$

$$\Gamma_{2\gamma}(n_{Ps}, {}^1S_0) = \frac{1}{2} \frac{mc^2}{\hbar} \frac{\alpha^5}{n_{Ps}^3}$$

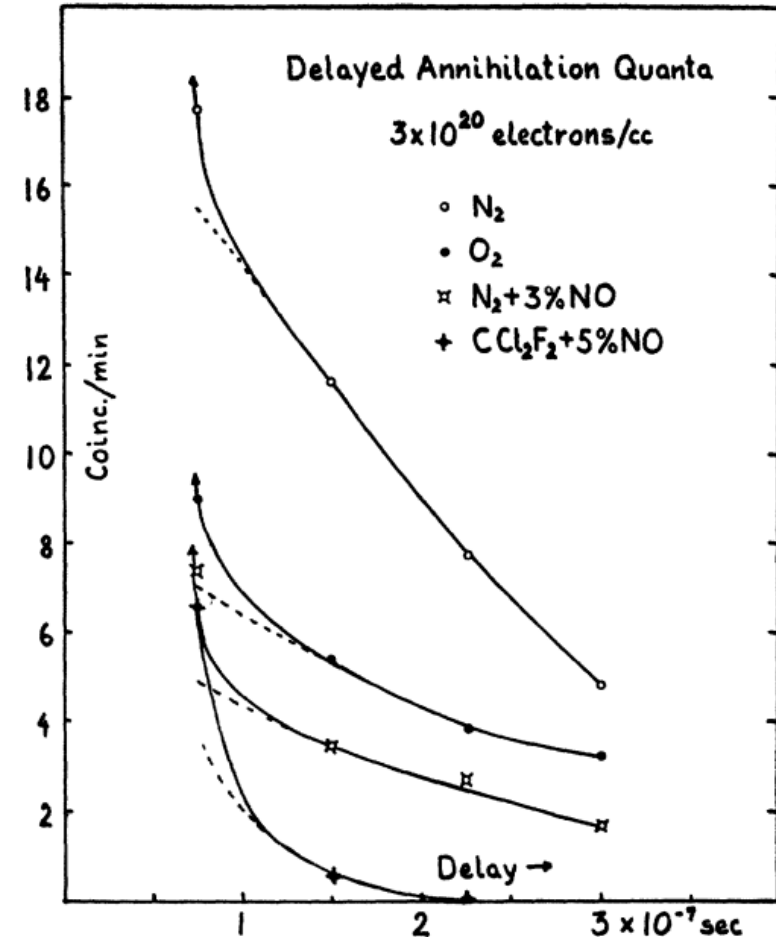
Properties of Positronium

- Energy level diagram for the $n < 3$ manifold of Ps
 - Lifetimes against annihilation and transition shown below each state

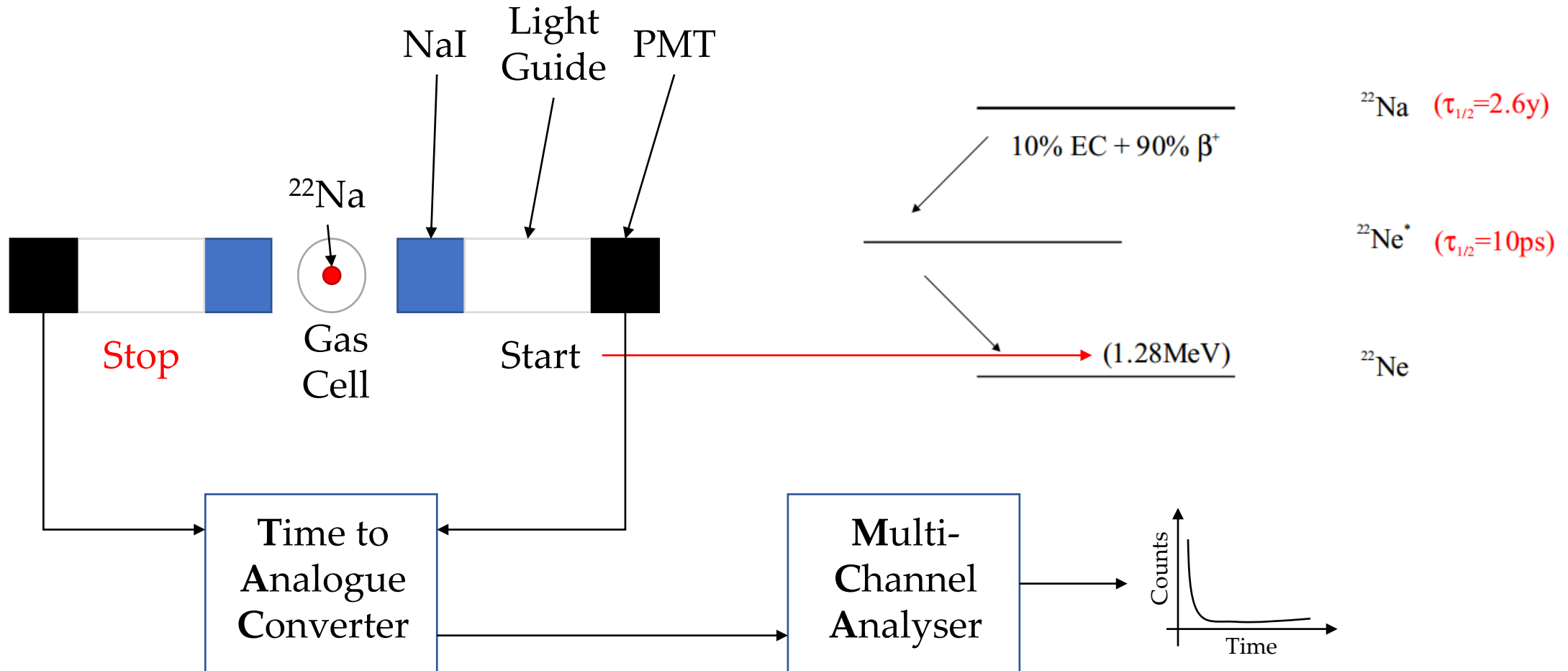


Observation

- The first experimental evidence for the existence of Ps came in 1951 in a paper titled: "Evidence for the formation of Positronium in Gases" by Deutsch



Positron annihilation lifetime



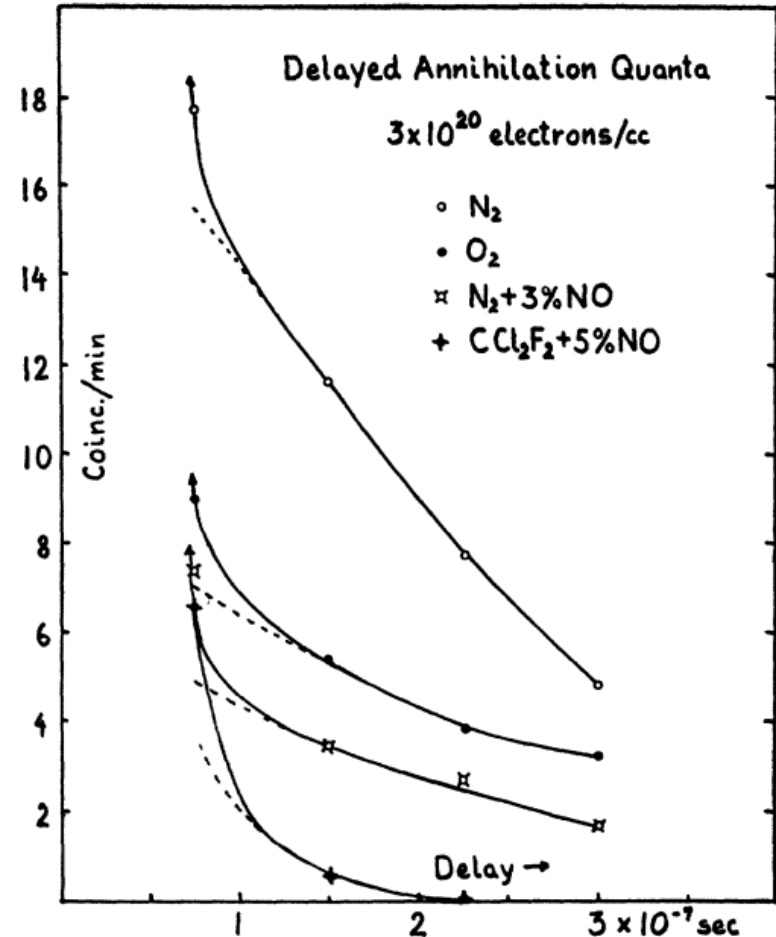


Positron annihilation lifetime

- Positron interactions with gas
 - Annihilation: $e^+ + A \rightarrow A^+ + 2\gamma$
 - Ps formation: $e^+ + A \rightarrow A^+ + Ps$
 - Direct ionization: $e^+ + A \rightarrow A^+ + e^+ + e^-$
- Positron interactions with a metal
 - The positron will lose energy, find a defect and annihilate.

Properties of Ps

- What we see is the time difference between the 1.28 MeV photon from the β^+ decay of ^{22}Na and the annihilation quanta from ortho-Ps
- Mixing NO with the gas being studied enhances the probability of spin flipping the Ps

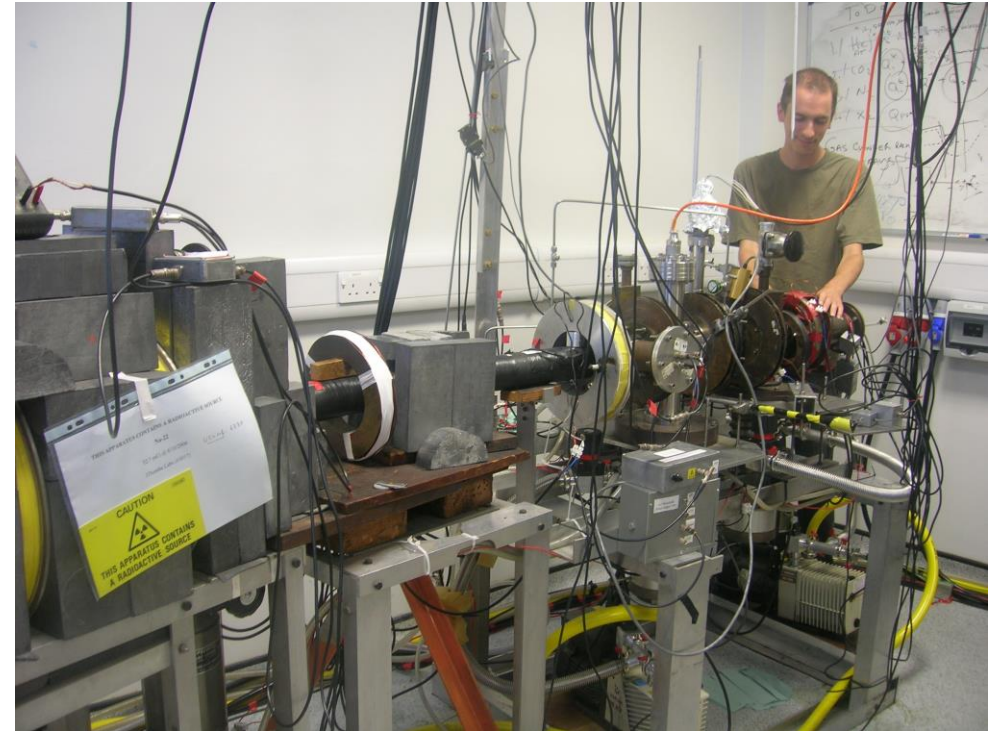


Summary: Properties of Ps

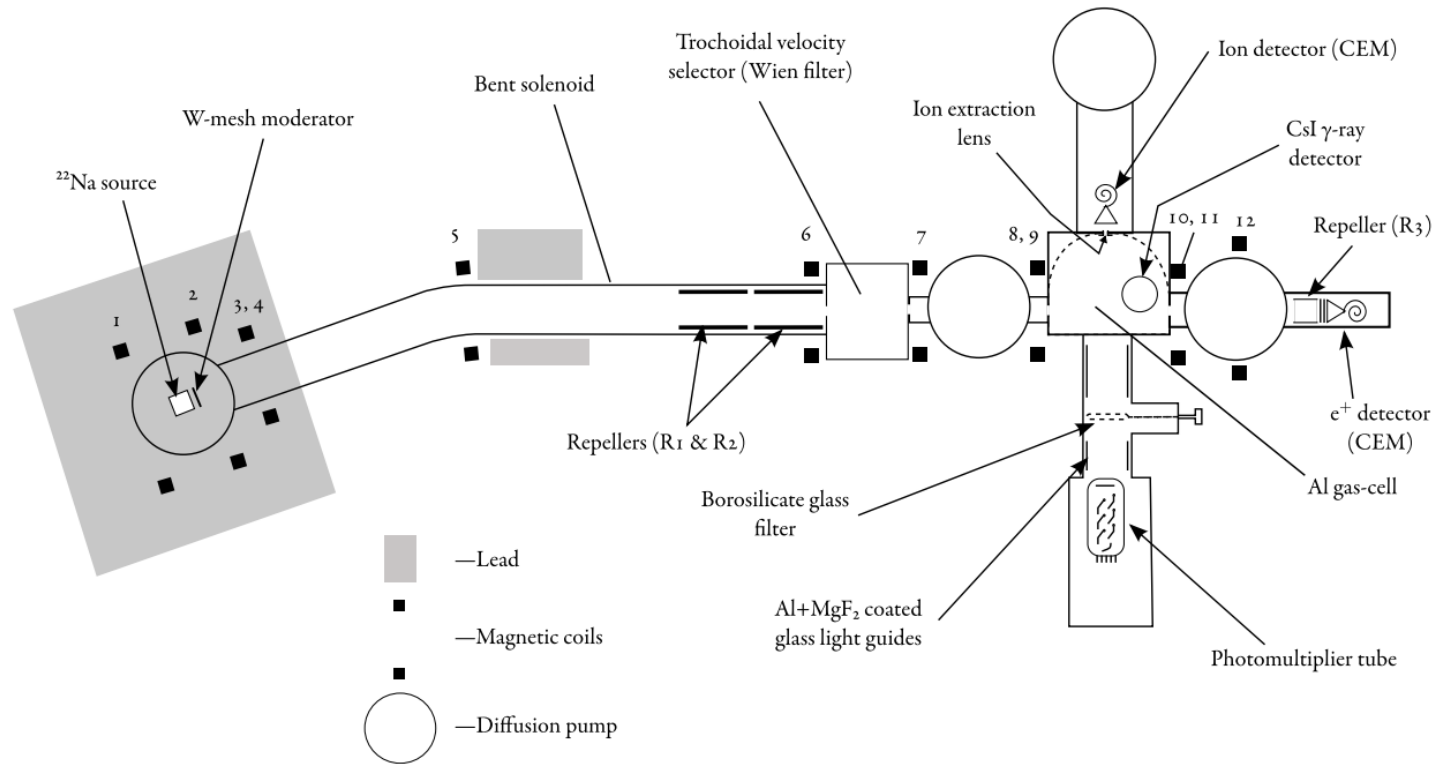
- Positronium is a hydrogenic quasi-stable atom
- Due to the larger magnetic moment of the positron than the proton the fine structure has a triplet and singlet state
- Ground state lifetime
 - Singlet state : 125 ps
 - Triplet state 142 ns
- The number of annihilation quanta is determined by charge parity conservation
 - Photon $P_c = (-1)^n$
 - Ps system $P_c = (-1)^{L+S}$
- Bohr radius is twice that of hydrogen
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Production of Ps

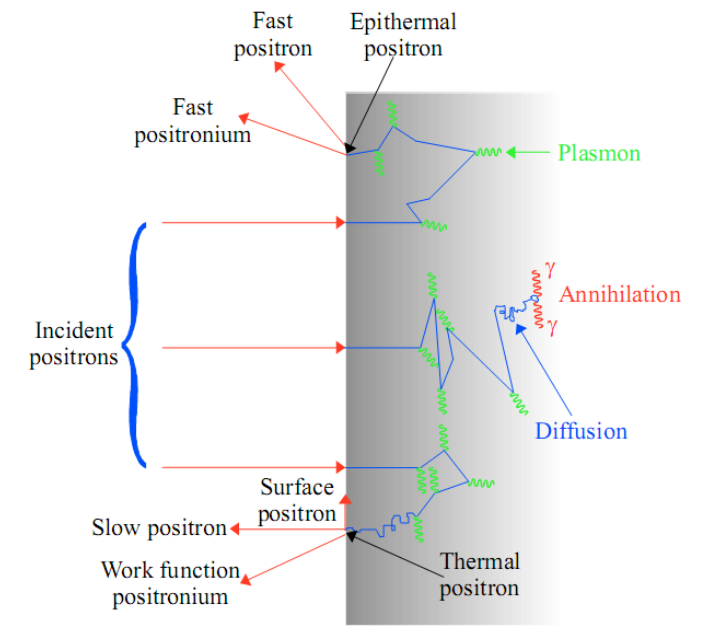
- The lifetime method used by Deutsch continues to be used to this day in the form of **Positron Annihilation Lifetime Spectroscopy (PALS¹)**
- To discuss the production of Ps we have to briefly discuss the production of positron beams



Positron beam production

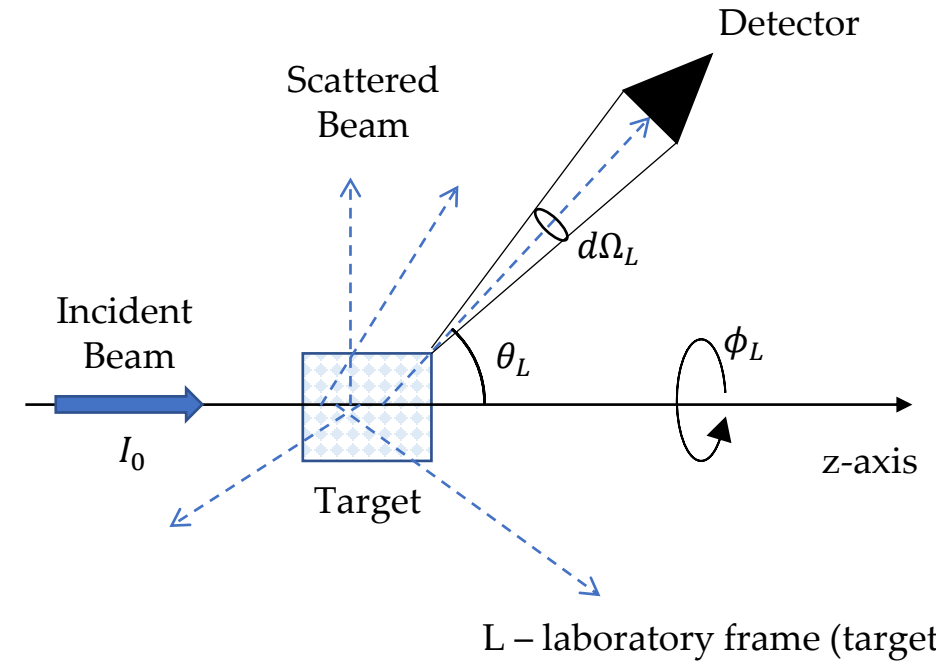


Moderation



Positronium Formation Cross-Section

- Once you can produce a beam of positrons and interact them with a gas of known pressure and cell length you can begin to measure cross-sections



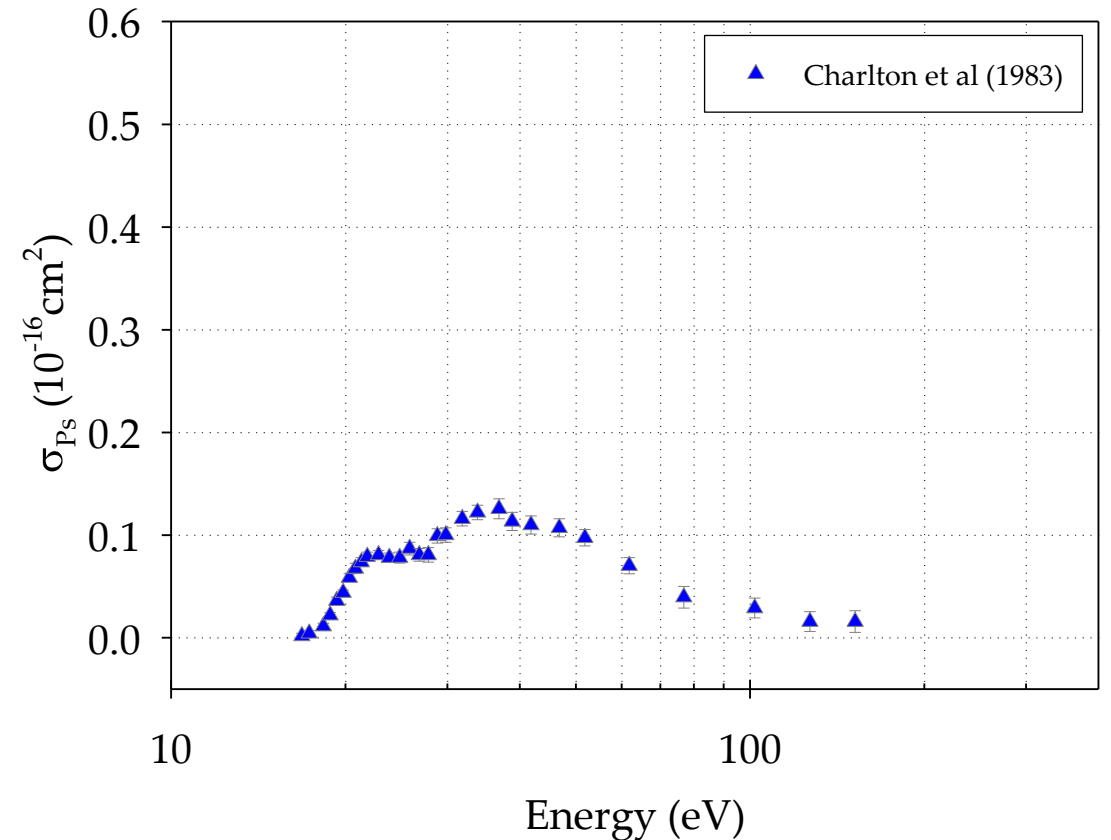
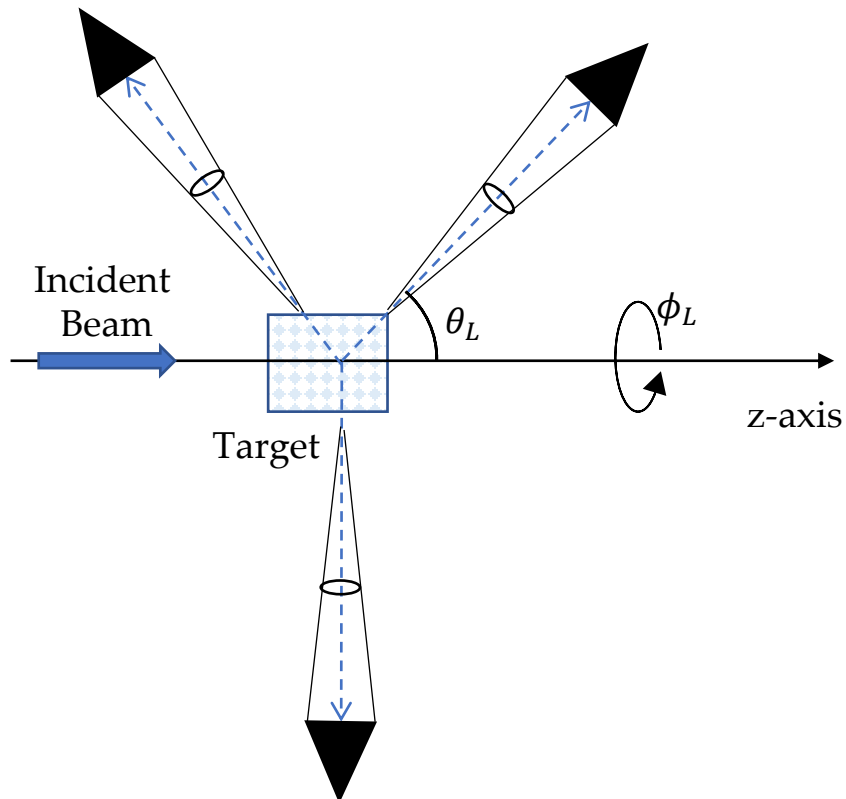
$$dN_{el} = I_0 \sigma_L^{el} (d\Omega_L) d\Omega_L$$

$$I_0 \sigma_L^{el} (d\Omega_L) d\Omega_L = I_0 \sigma_L^{el} (\theta_L, \phi_L) d\Omega_L$$

$$N_{el}^T = I_0 \sigma_L^{el}$$

Positronium Formation

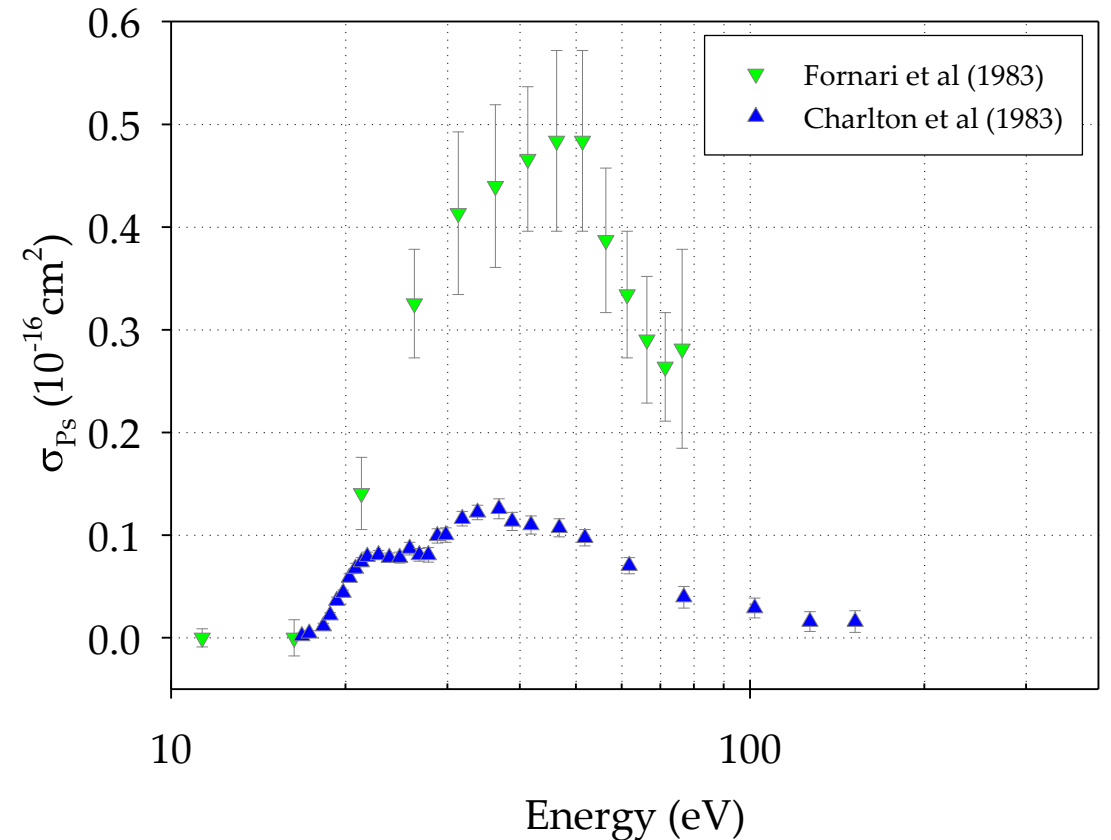
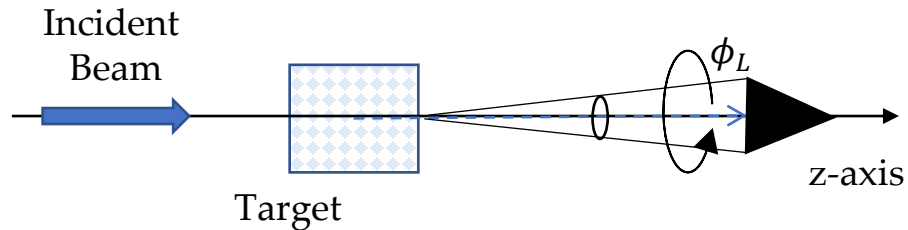
- First measurements of Ps formation cross-sections relied on triple coincidences from o-Ps annihilation



1. See D J Murtagh, M Szłuińska, J Moxom1, P Van Reeth and G Laricchia *J. Phys. B.* **38** 3857 (2005) and references there in

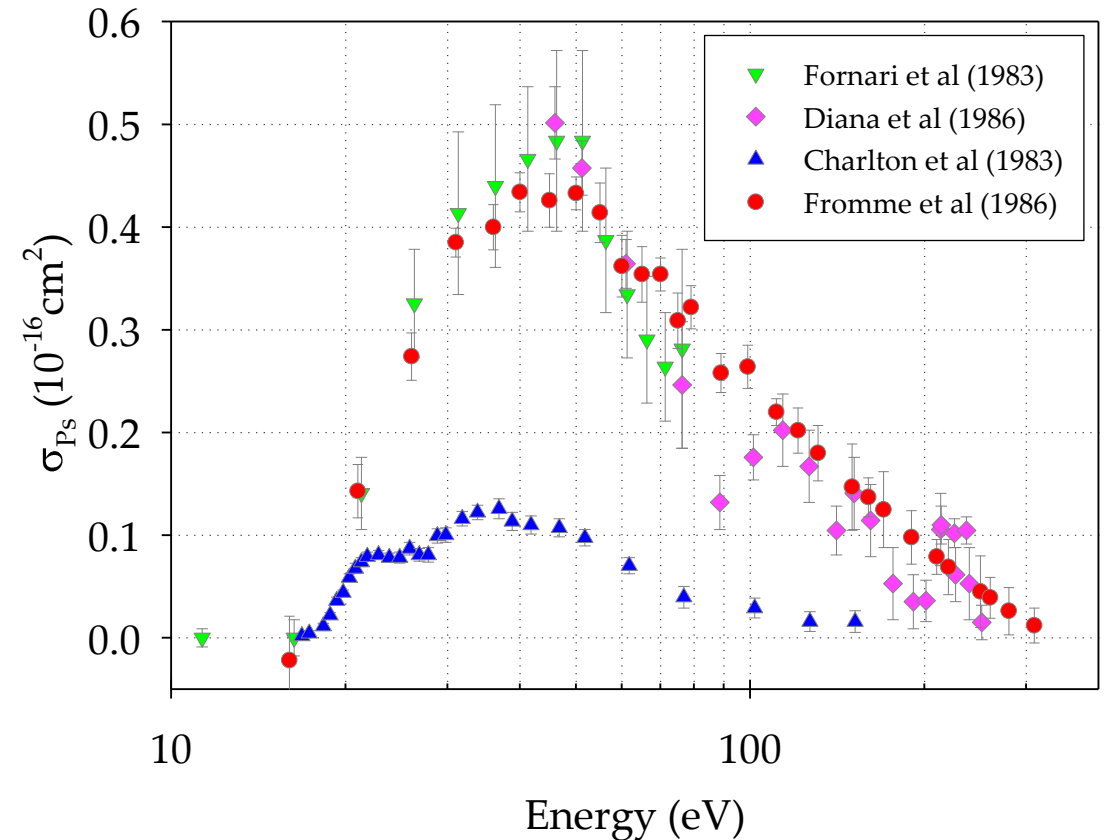
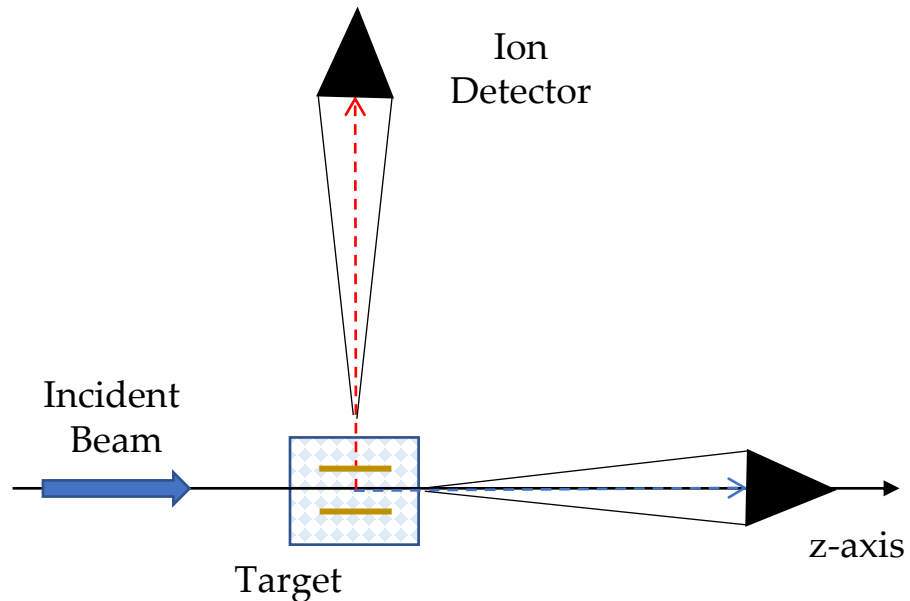
Positronium Formation

- The next method used was the loss of the positron in the final state
- The cross-sections are very different in magnitude - Charlton *et al* didn't know the Ps would exit the gas cell or hit the walls!



Positronium Formation

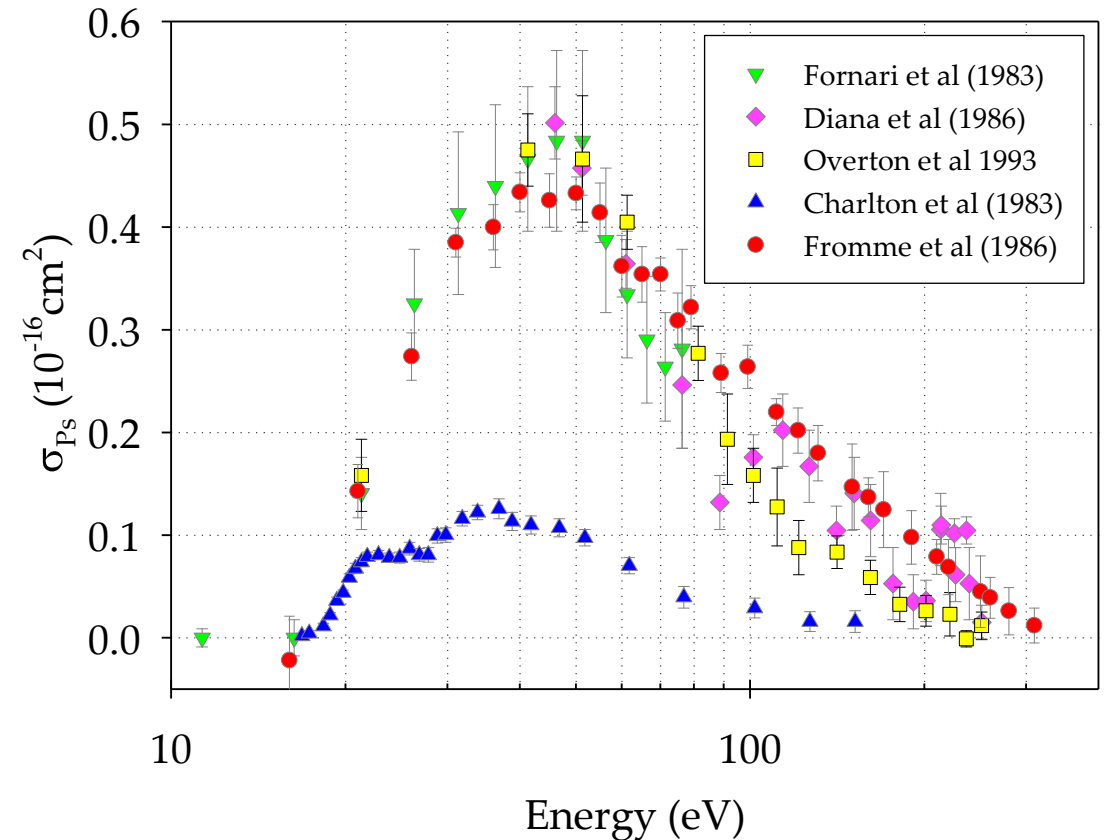
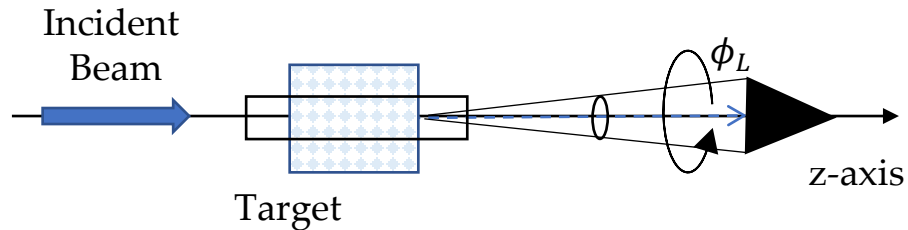
- Another idea was to measure the production of ions and positrons in coincidence. In the case of positronium an ion should be formed without an associated positron



1. See D J Murtagh, M Szłuińska, J Moxom1, P Van Reeth and G Laricchia *J. Phys. B.* **38** 3857 (2005) and references there in

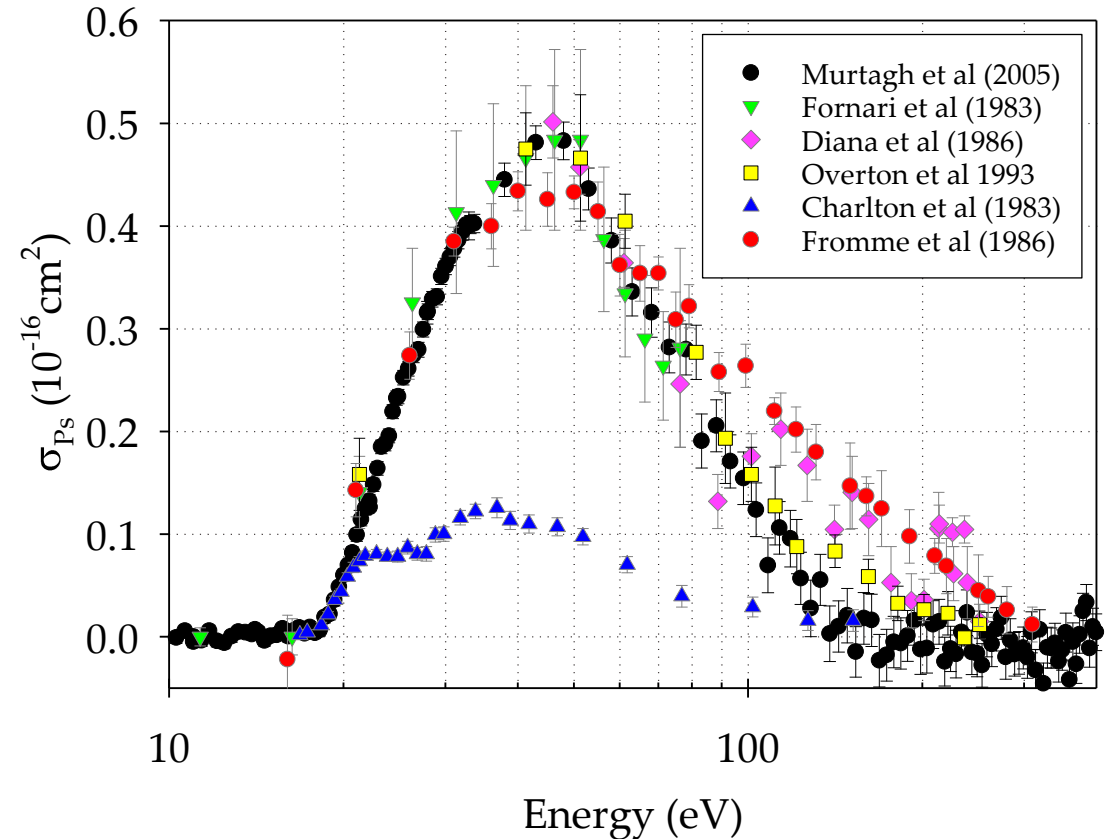
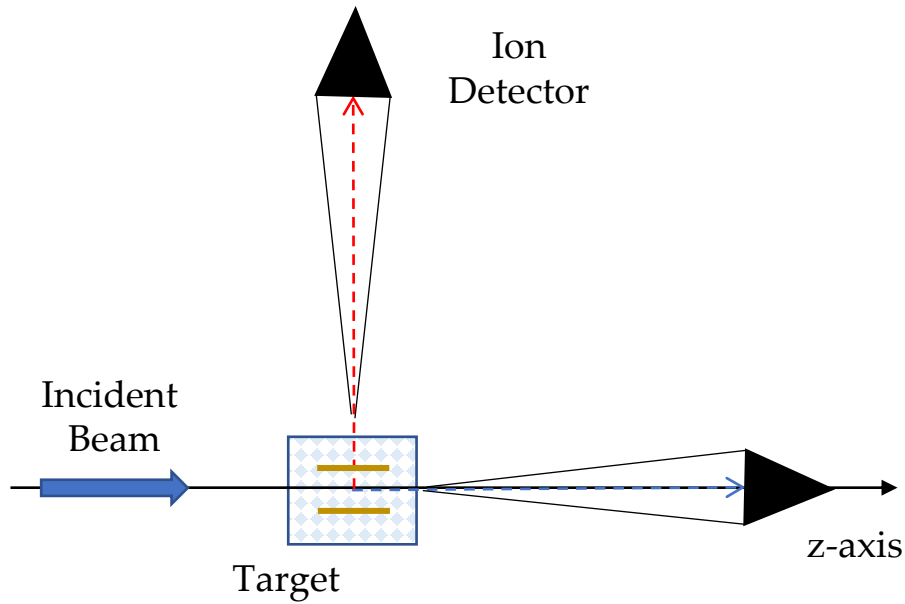
Positronium Formation

- Using the beam attenuation method special care has to be taken to confine scattered particles



Positronium Formation

- $\sigma_i^t = \sigma_{PS} + \sigma_i^+ + \Sigma\sigma_{HO}$
- $\sigma_{PS} = \sigma_i^t - \sigma_i^+$



Positronium in external fields

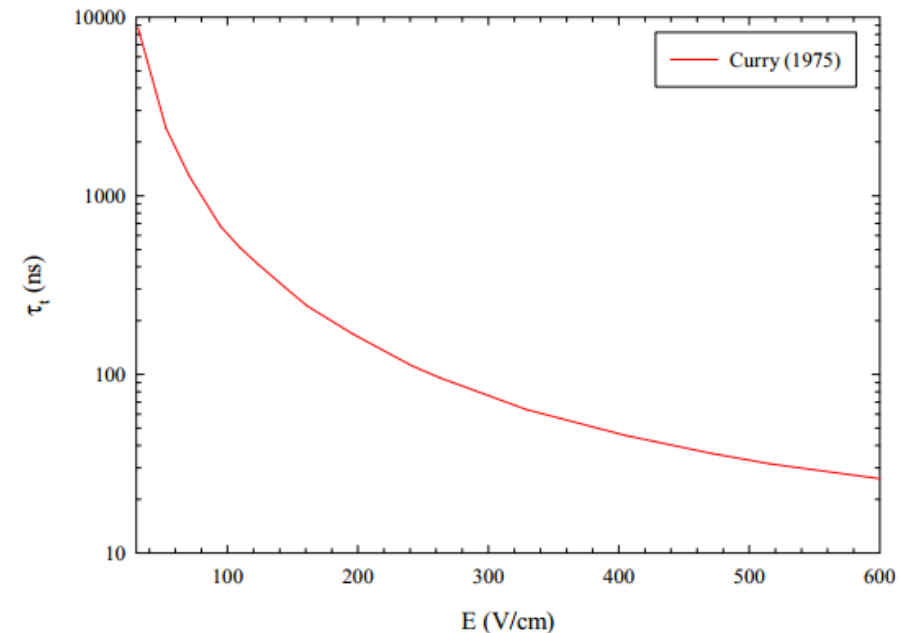
- Stark effect → Mixes l states
- Zeeman effect → Mixes triplet and singlet states
- Motional Stark effect → Mixes l states however unlike the conventional Stark effect is due to the Ps motion transverse to a **B** field :

$$\vec{E}_{mse} = \gamma(\vec{v} \times \vec{B})$$

$$\gamma = 1/\sqrt{1 - \beta^2}, \quad \beta = v/c$$

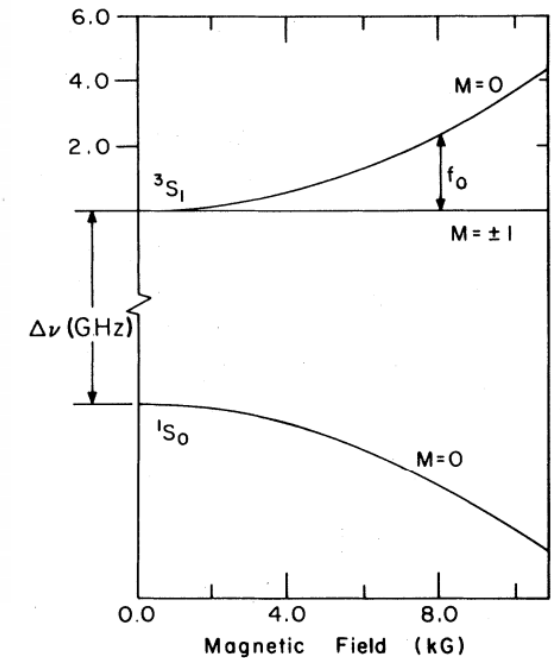
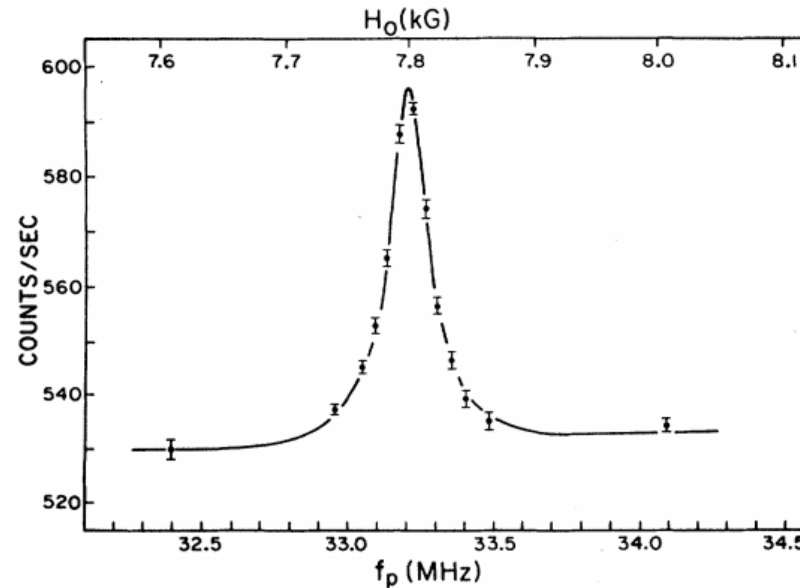
$$\lambda'_3 = \frac{\lambda_3 + y^2 \lambda_0}{1 + y^2} \quad y = \frac{\sqrt{1 + x^2} - 1}{x} \quad x = \frac{4\mu_B B}{\Delta W}$$

$$\Delta W = 25,423 \text{ MHz}$$



Hyperfine structure of Ps

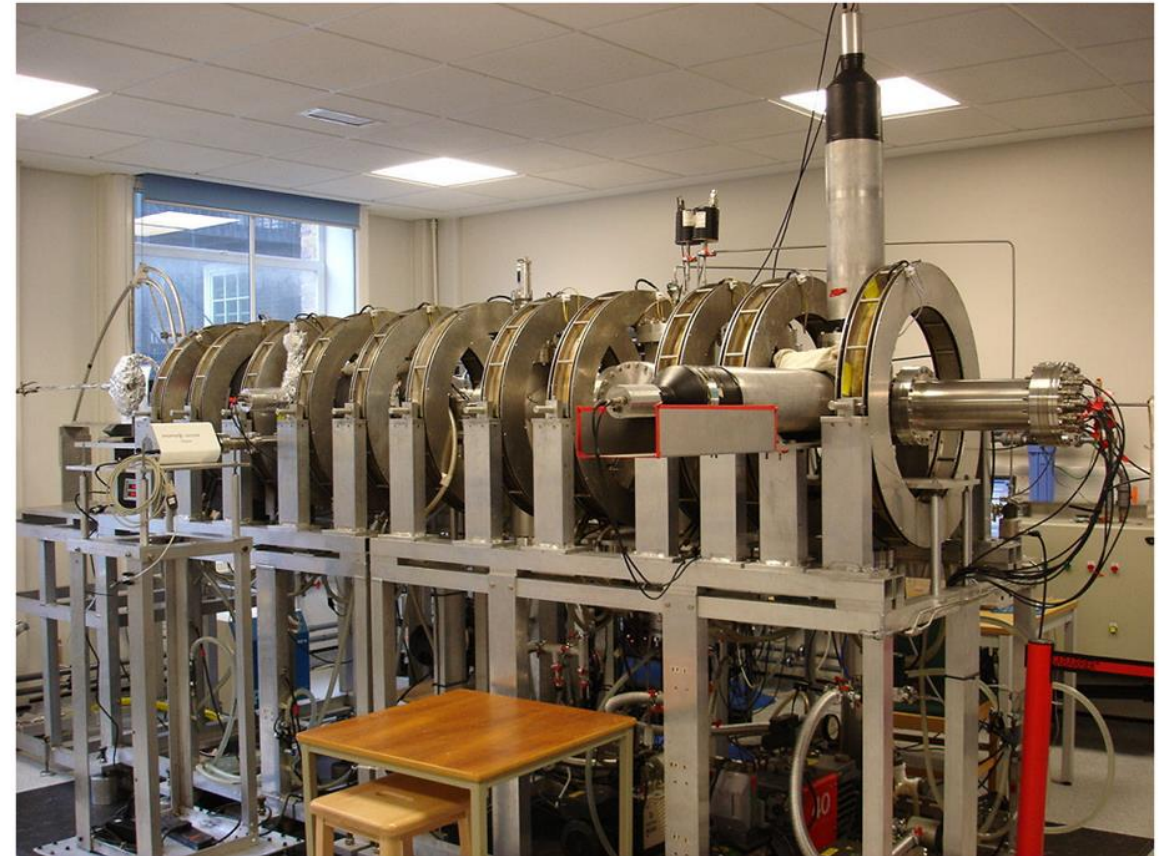
- Triplet and Singlet $M_J = 0$ states are mixed in a magnetic field
- By inducing a transition from the triplet $M_J = \pm 1$ states to $M_J = 0$ an increase in the 2γ annihilation rate can be observed



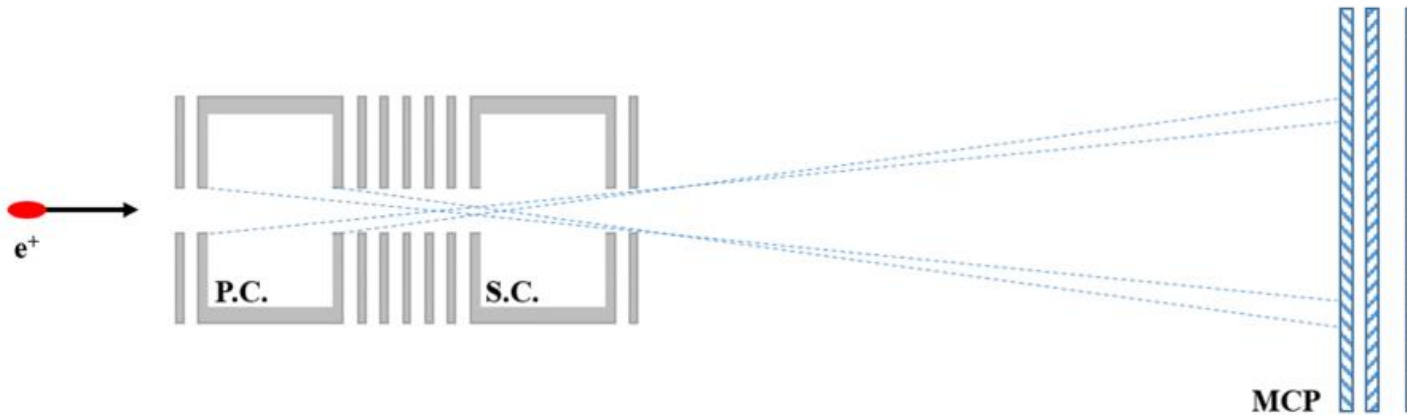
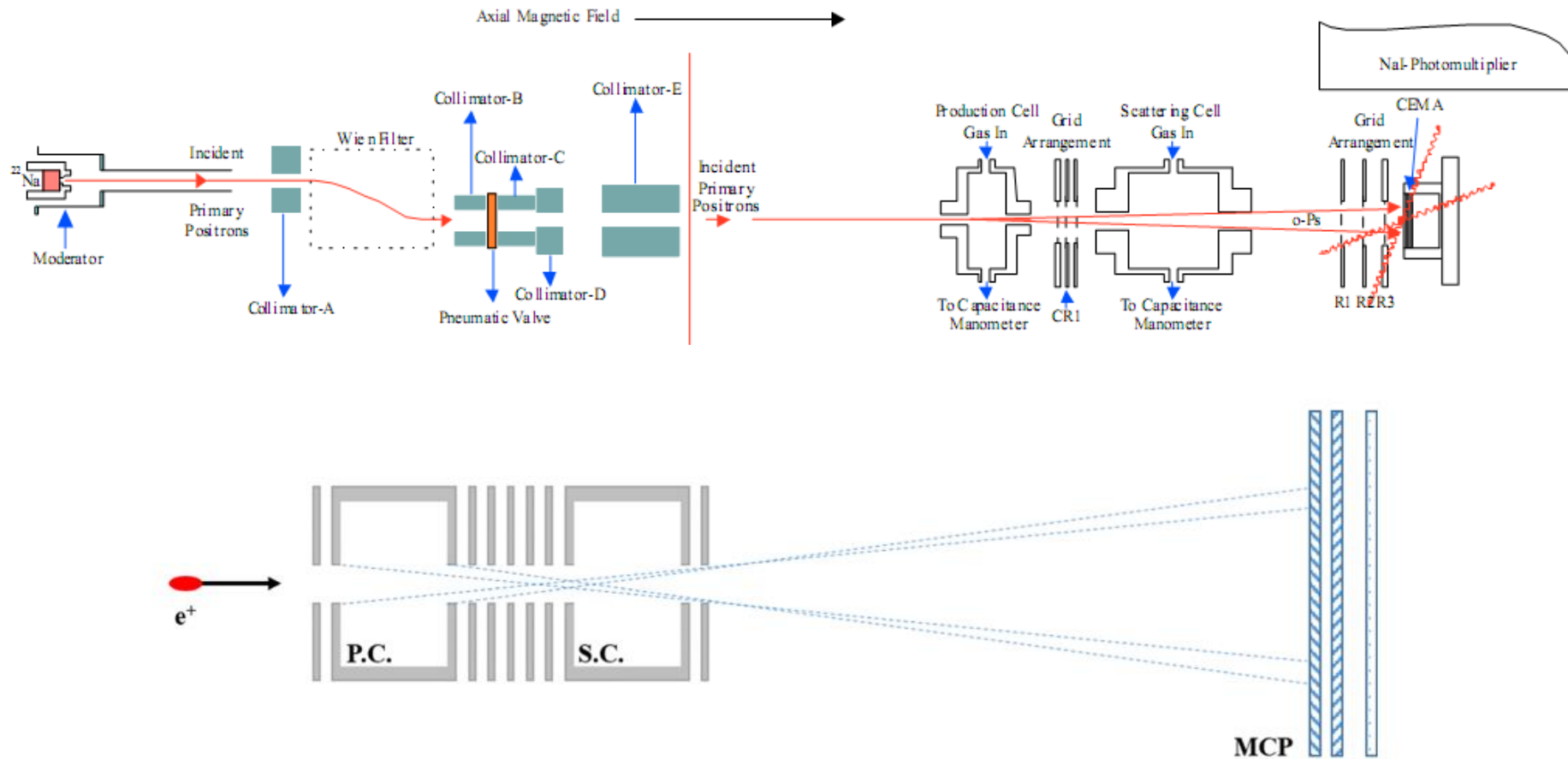
$$\Delta\nu = 203.38910(74) \text{ GHz}$$

Positronium Beam

- We know it's possible to collide positrons with a gas target and produce Ps
- We know from the early work of Charlton that the Ps has energy to leave the production region
- The next step was to create a beam of positronium



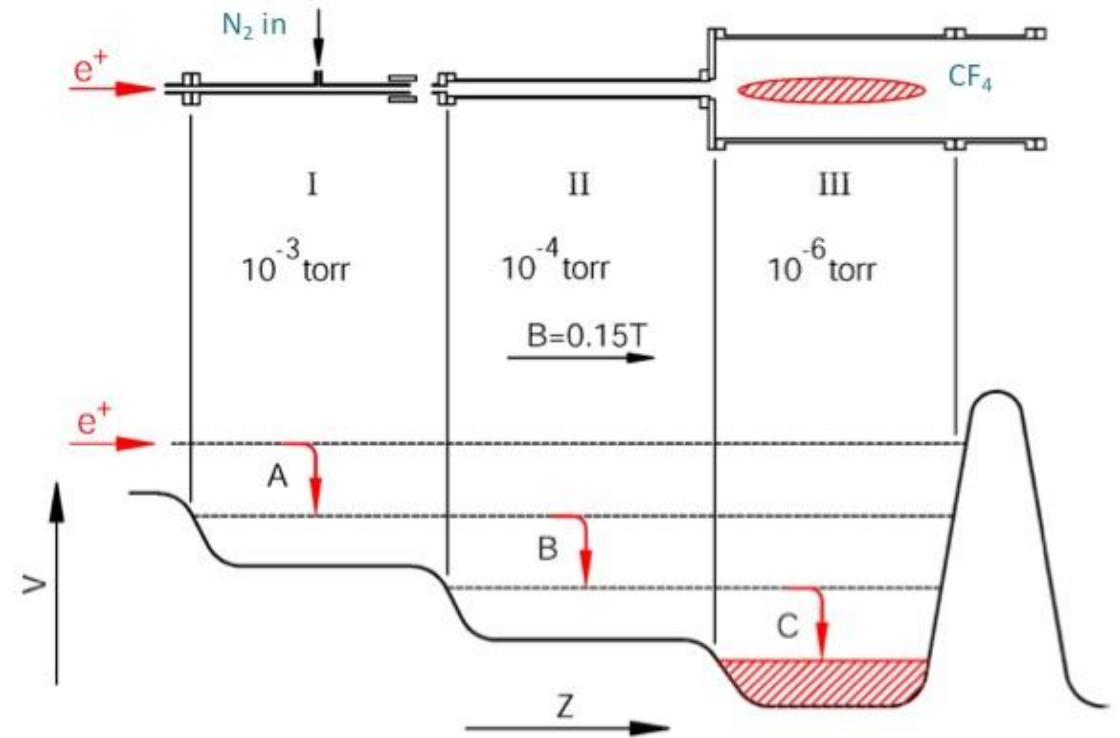
Positronium Beam



1. N. Zafar, G. Laricchia & M. Charlton *Hyperfine Interactions* **89** 243 (1994)
2. J. R. Machacek, S. J. Buckman, and J. P. Sullivan *Review of Scientific Instruments* **91**, 033311 (2020)

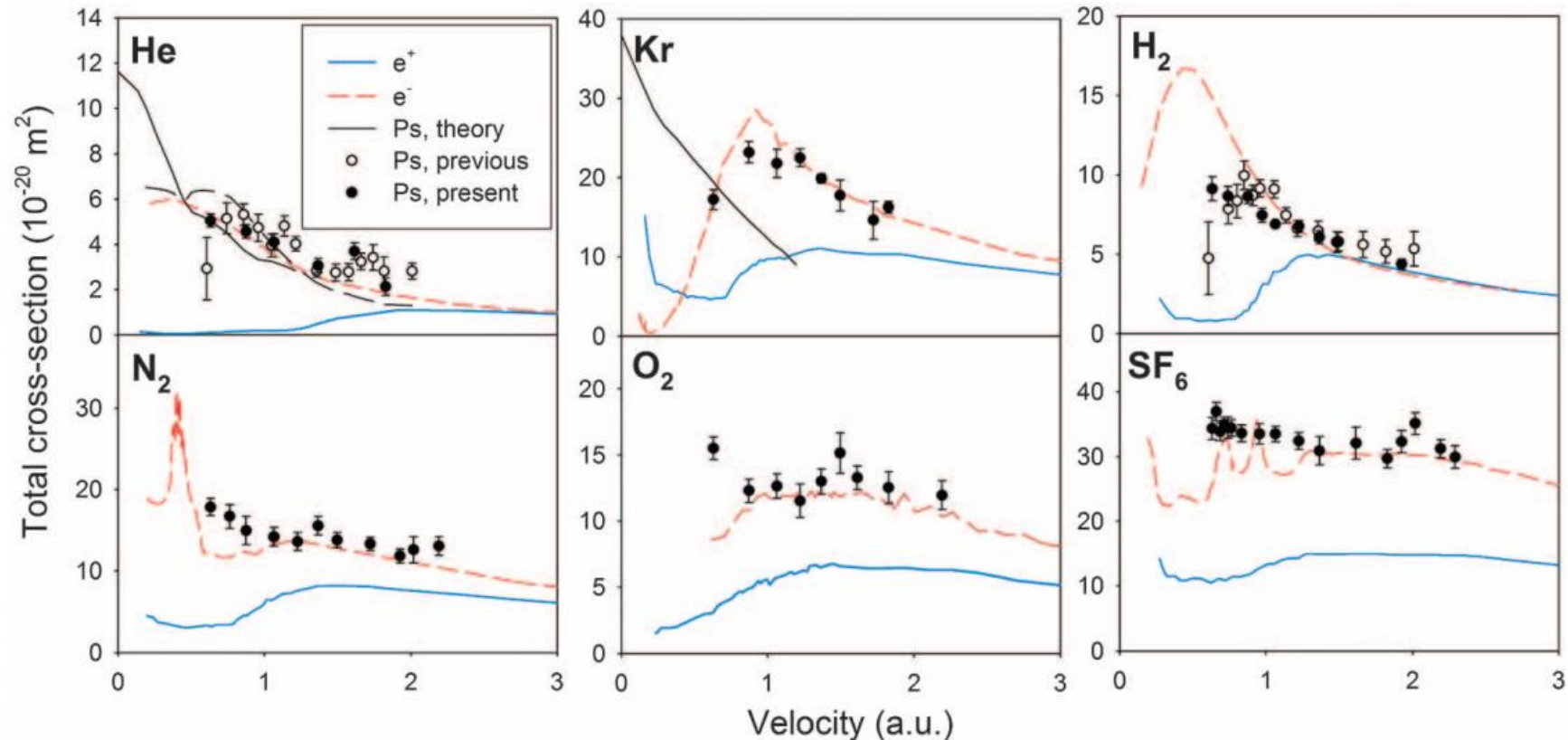
Buffer gas trap based beams

- The combination of a solid Ne moderated beam and N₂ buffer gas trap has enabled the production of high intensity, low energy spread, pulsed positrons beams in laboratories around the world



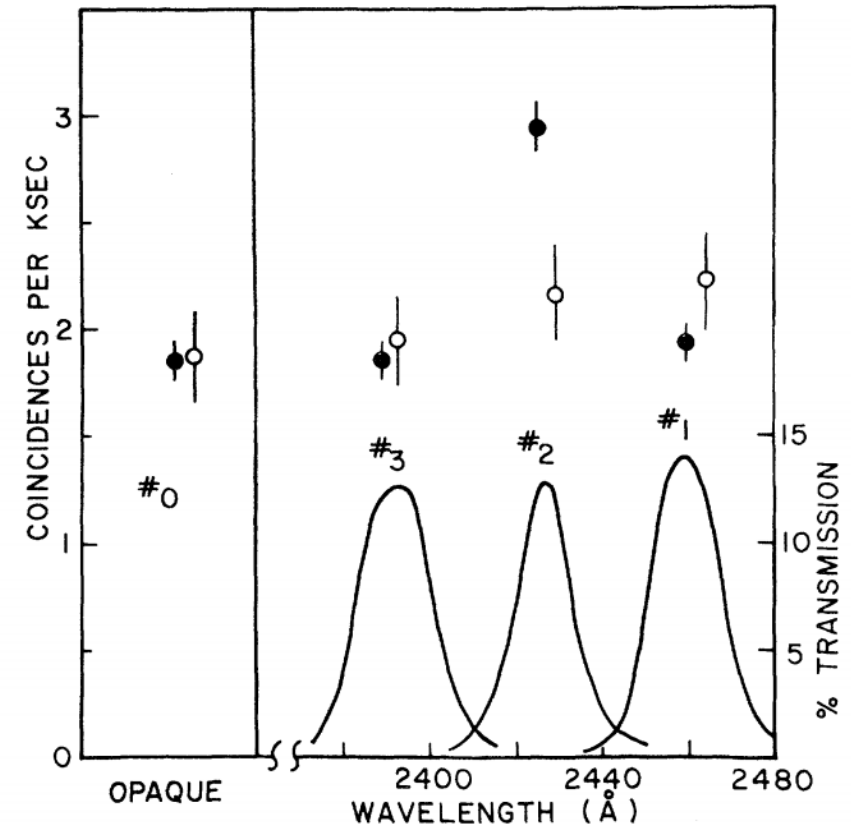
Positronium interactions with atoms and molecules

- Positronium scatters like an electron with the same velocity



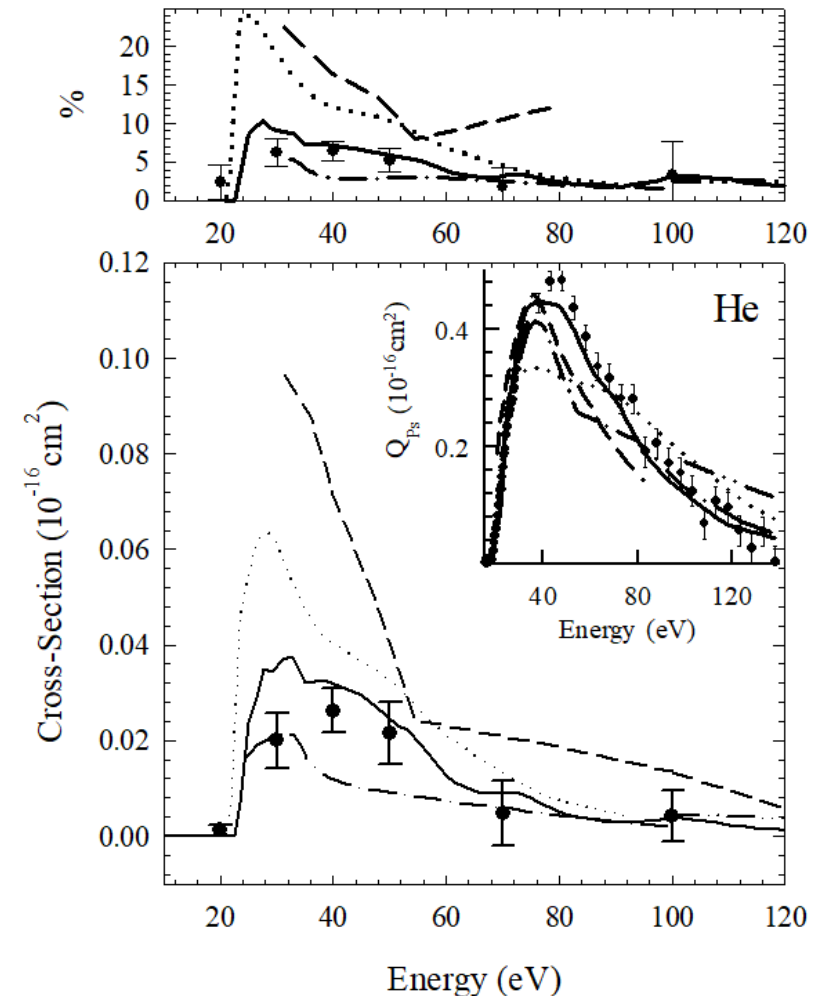
Excited states of Positronium

- The first observation of Ps formed with $n > 1$ was made by Canter et al in 1974
- Positrons were allowed to impinge on a Ge target
- A PMT was used to observe photons
- 3 different filters were placed in front of the PMT to isolate the 243nm light



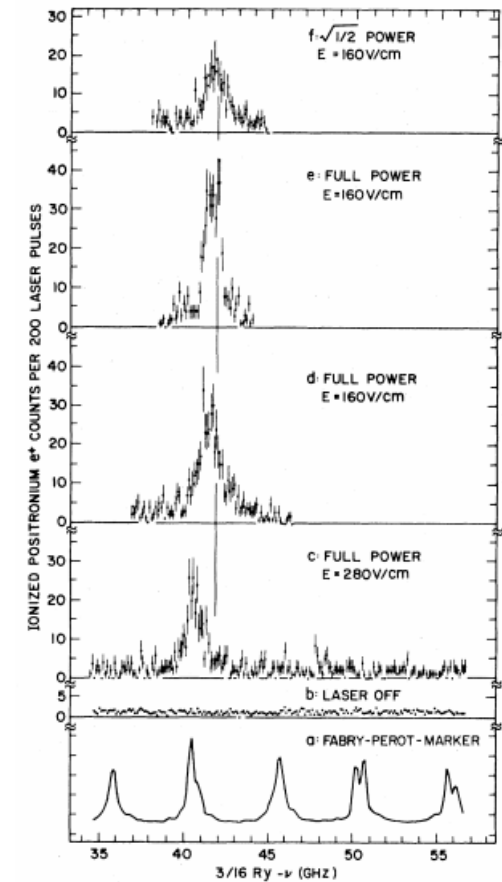
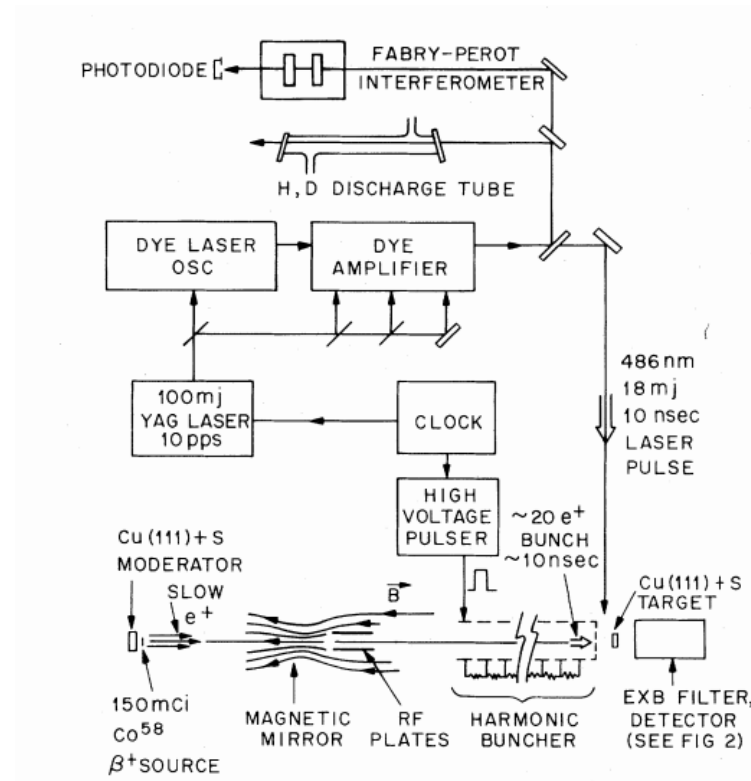
Excited states of Positronium

- First observation of Ps($n=2$) in a gas target was made in 1985 by Laricchia et al
- The Ps(2P) formation cross-section for He, Ar and Xe have been measured.
- The fractional contribution was found to be
 - He - $(6 \pm 1)\%$
 - Ar - $(12 \pm 4)\%$
 - Xe - $(26 \pm 9)\%$



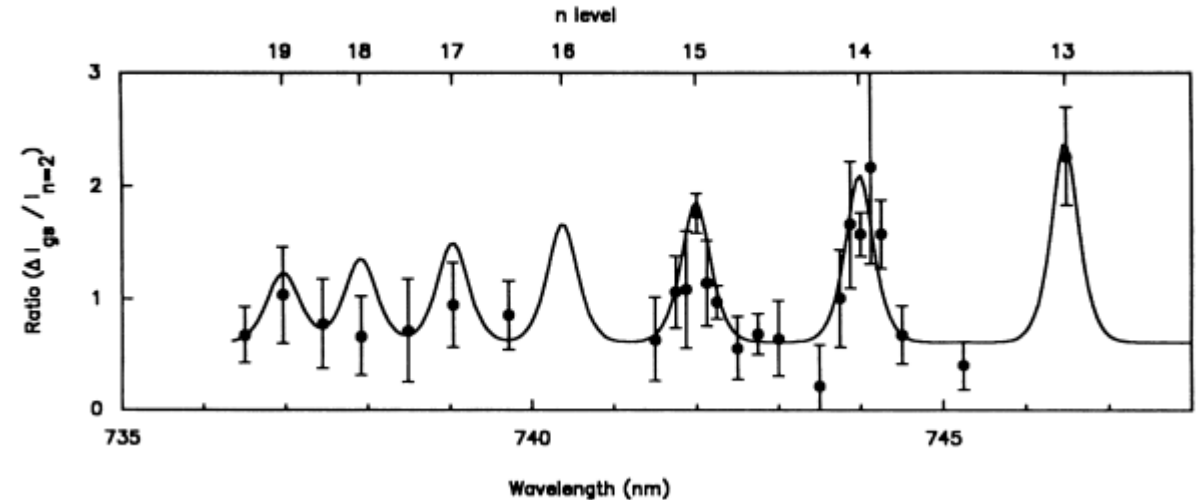
Laser Interactions with Ps

- The first observed laser induced transition was by Chu and Mills in 1982
- A resonant ionization process was used to detect positrons from the laser excited Ps atoms



Highly excited positronium

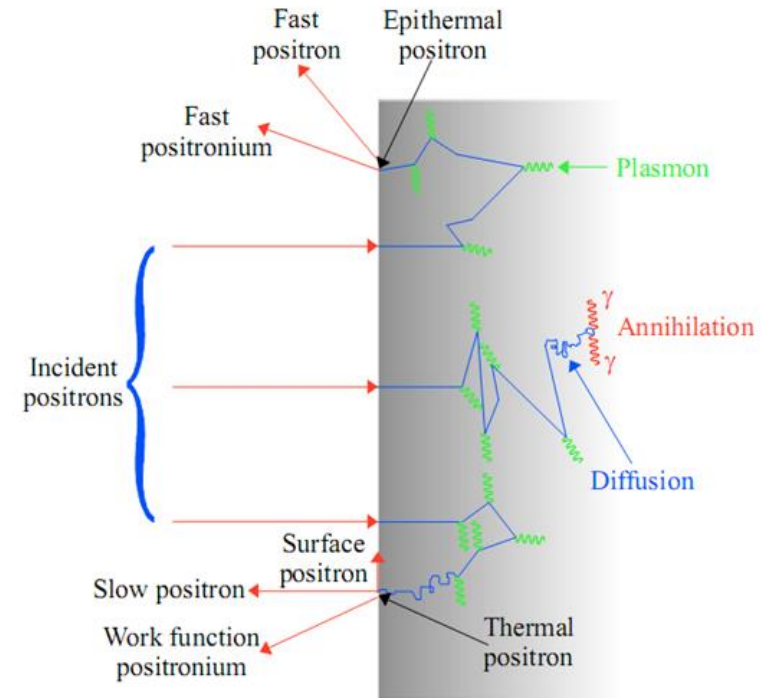
- Rydberg state Ps was first observed in 1990 by Ziock *et al*
- Pulses of positrons from the Lawrence Livermore 100 MeV linac $\rightarrow 10^5 e^+$ in 15 ns bunches
- $1s \rightarrow 2p \rightarrow 13l, 14l, 15l$ (highly excited states stark mixed)



For an excellent review of the topic of Ry Ps see D. B. Cassidy *Eur. Phys. J. D.* **72** 53 (2018)

Efficient production of Ps

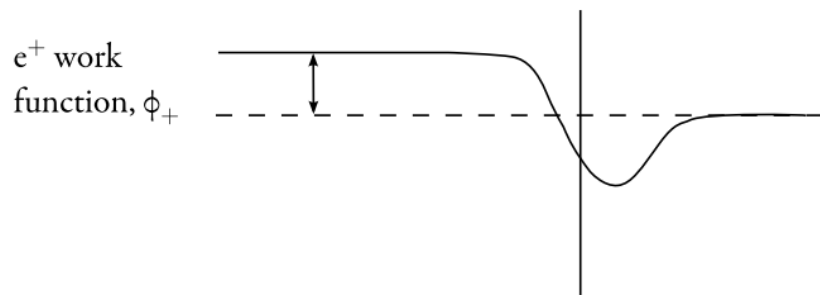
- Ps beam experiments make use of charge exchange collisions with gases
 - High energy Ps → okay for scattering not good for applications like optical excitation or charge exchange
- The previous optical excitation experiments made use of a single crystal Cu foils heated to 1000 K



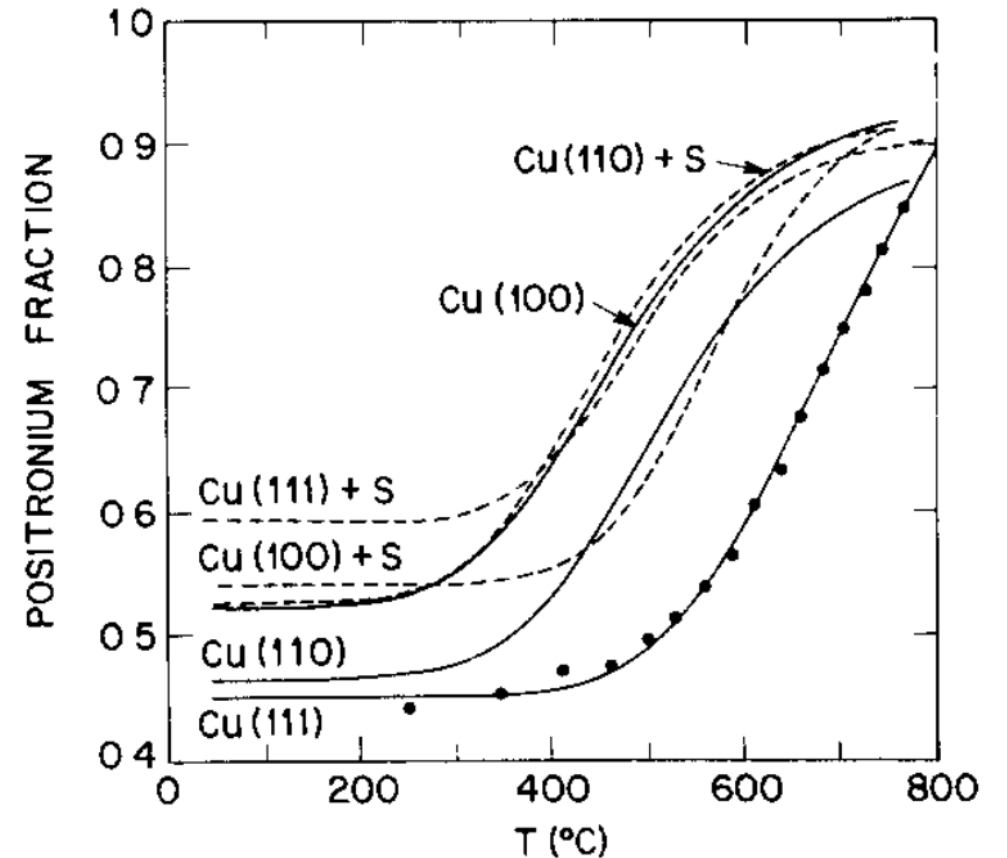
$$\phi_{PS} = \phi_+ + \phi_- - E_b$$

Efficient production of Ps

- Positrons can become trapped at the surface of metals
- Positronium production fraction
 - $f_{Ps} = f_{dir} + f_t(T)$



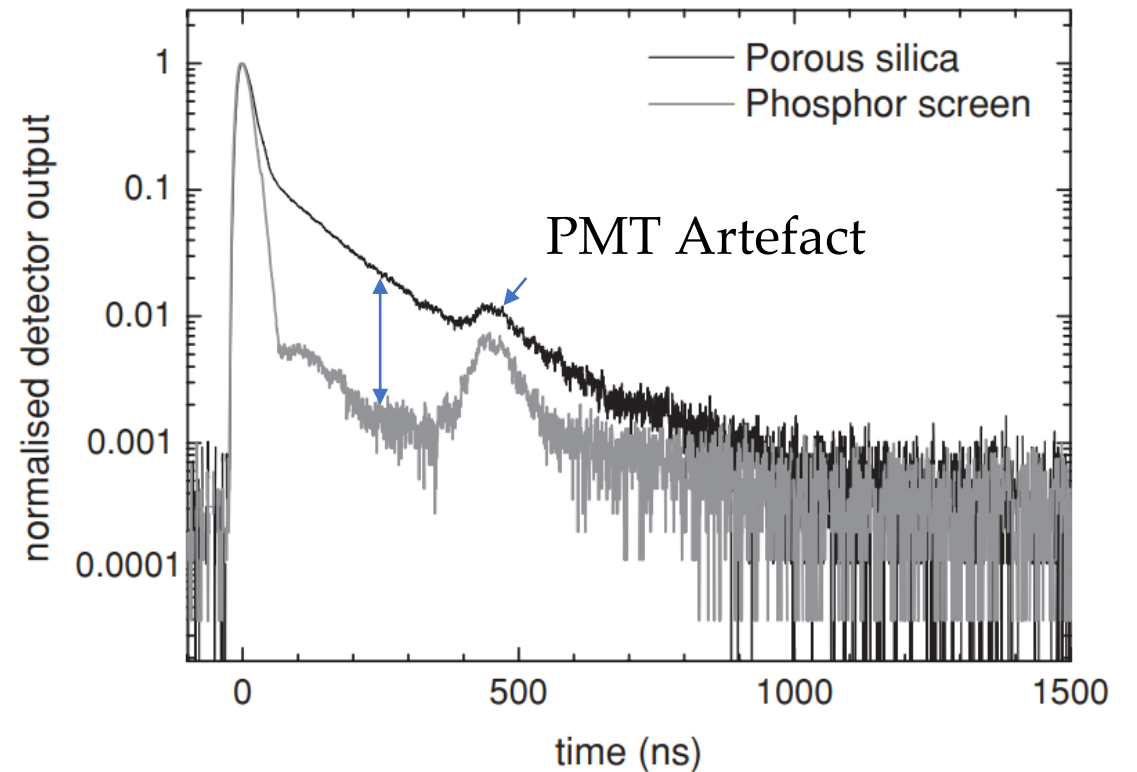
$$E_a = E_t + \phi_- - E_B$$



Efficient Ps production

- Thin mesoporous films have been used effectively as efficient Ps converters
- **Single Shot Positron Annihilation Lifetime Spectroscopy (SSPALS)**

Ps - 1.5 keV
Background - 5 keV

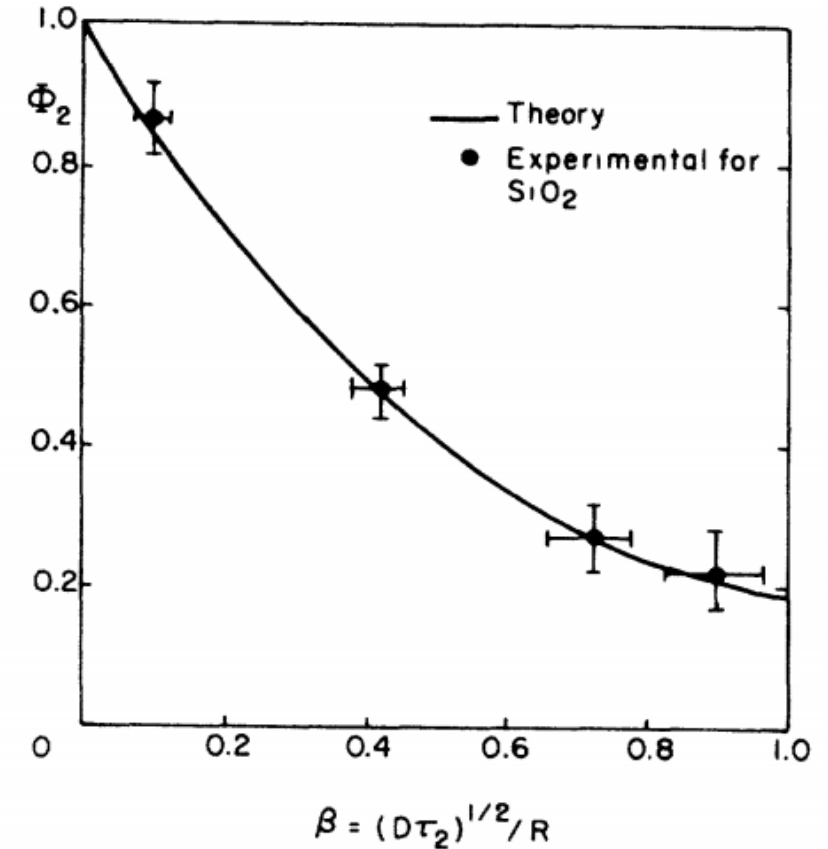


1. E.g. D. B. Cassidy, P. Crivelli, T. H. Hisakado, L. Liskay, V. E. Meline, P. Perez, H. W. K. Tom, and A. P. Mills, Jr. *Phys. Rev. A* **81**, 012715 (2010)

2. See references in D. B. Cassidy *Eur. Phys. J. D.* **72** 53 (2018)

Efficient production of Ps

- Powders (e.g. MgO, SiO₂, Al₂O₃) can also efficiently form Ps
- Ps can also diffuse from the bulk in these materials and reach the surface to be emitted into vacuum



Diffusion Coefficient $D = (5.8 \pm 1.9) \times 10^{-5} \text{ cm}^2\text{s}^{-1}$

Radius $R = 273, 76, 45, 36 \text{ \AA}$

O-Ps in bulk $\tau_2 = (1.8 \pm 0.2) \text{ ns}$

Powders and Pours

- Ps can be formed in the bulk and on surfaces

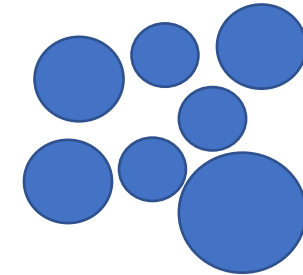
- $T_{bulk} = -\phi_+ - \phi_- + E_G - E_B + 6.8 eV$

- $T_S^v = -\phi_+ - \phi_- + 6.8 eV$

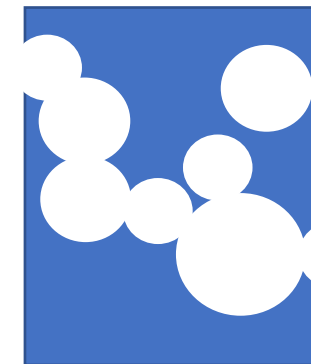
- $T_S^c = -\phi_+ - \phi_- + E_G + 6.8 eV$

- Ps atoms formed on surfaces with thermal conduction electrons can penetrate into the bulk
- Ps formed in the bulk are confined within the voids

Indecent Beam



$r \sim 2-8 \text{ nm}$



Pours and trapping

- Positronium trapped in a cavity will cool via multiple collisions with the walls
 - Rate is dependent on pour size
 - $\Delta E = \frac{2m_{Ps}}{M} (E(t) - \frac{3}{2} K_b T)$
 - Classical Model \rightarrow okay for large pour size
- Quantum mechanical effects will begin to dominate when the de Broglie wavelength is close to the size of pours
 - $\lambda = h \sqrt{\frac{1}{2E_{Ps}m_{Ps}}}$
- As the wavelength approaches the pour size cooling stops
- Quantum tunnelling then dominates the diffusion of the Ps through the material
- Ps emitted to vacuum will convert this potential energy into KE

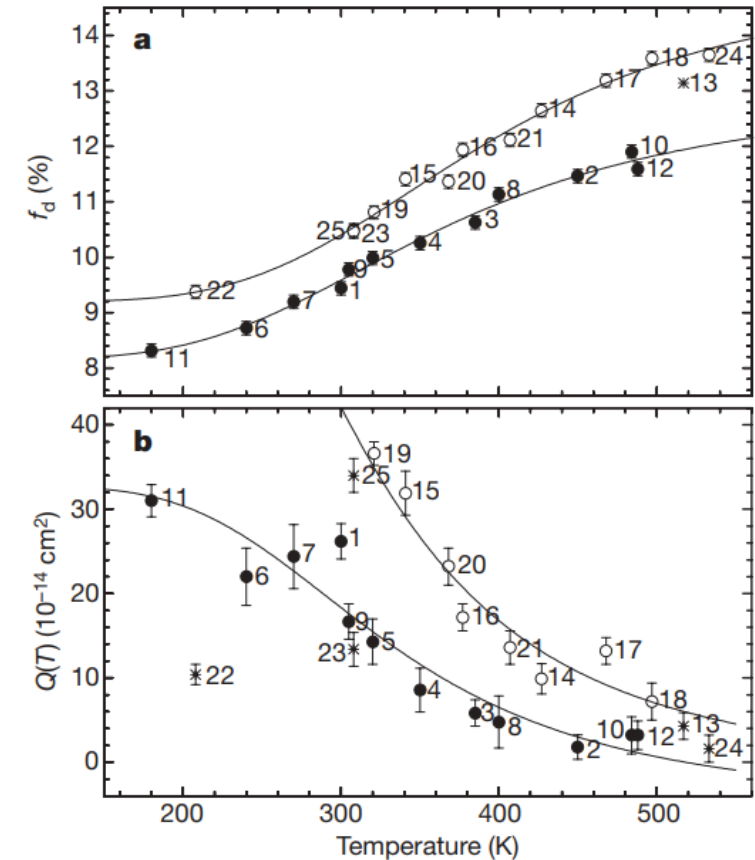


Dense Positronium Clouds

- If you can produce a dense enough positron plasma and have a good enough Ps converter you may be able to interact more than one Ps atom
 - Sub ns bunches
 - Areal density $\sim 1 \times 10^{10} \text{ cm}^{-2}$
- Elastic Scattering
 - $oPs + oPs \rightarrow oPs + oPs$
- Spin Exchange
 - $oPs + oPs \rightarrow 2pPs + 2E_{hfs}$
- Ps molecule formation
 - $X + oPs + oPs \rightarrow X + Ps_2 + E_b$
 - Three body recombination

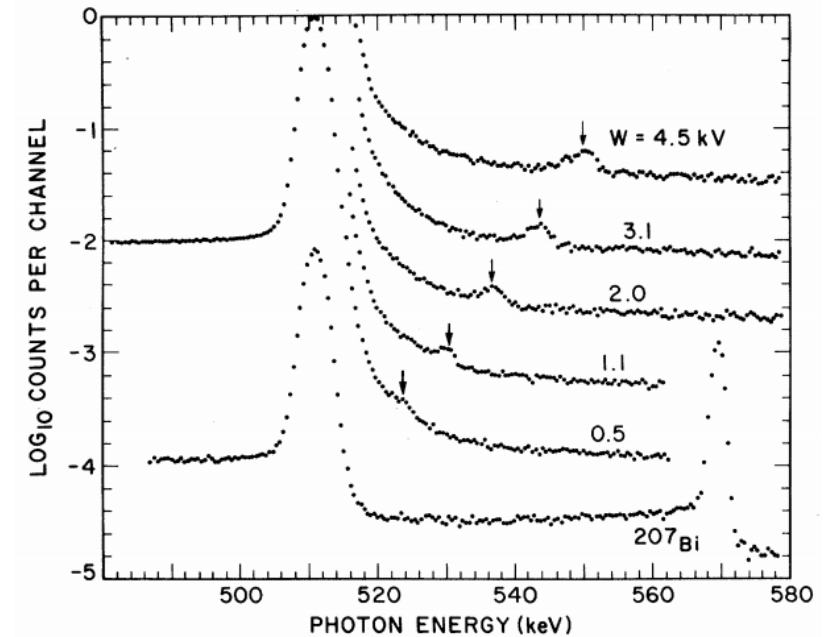
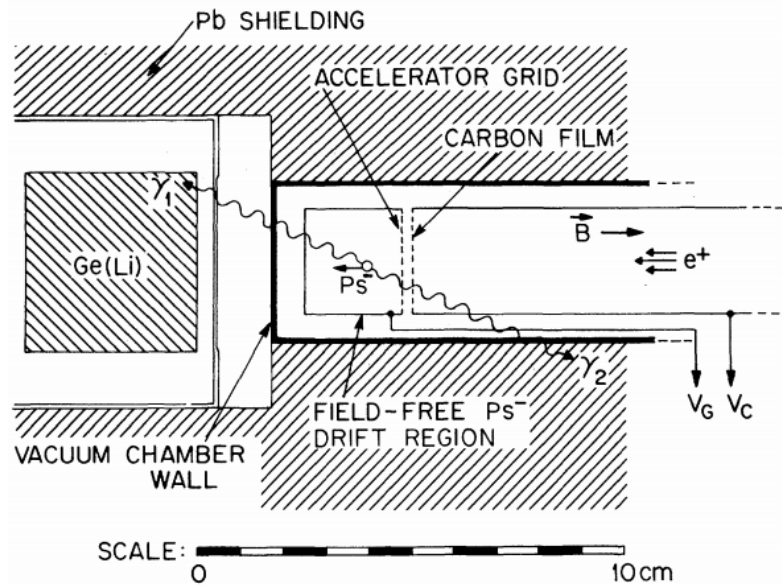
Molecular Ps

- Using SSPALS Cassidy and Mills looked for a signal from Ps_2 formation
- Both spin exchange and Ps_2 formation will cause oPs to annihilate more quickly
- This reduces the so called 'delayed fraction' (f_d)
- By increasing the temperature of the sample Ps atoms cannot bind to surfaces limiting the availability of a third body and hence suppressing Ps_2



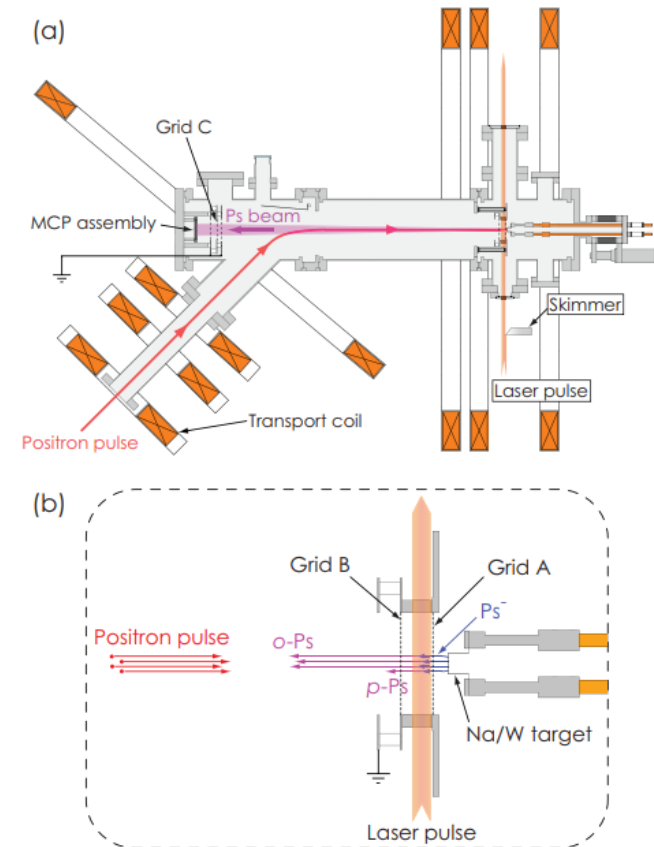
Observation of the Positronium Ion

- Mills observed a blue shifted annihilation line from Ps^- formed by passing a positron beam through a carbon foil $\sim 50 \text{ \AA}$



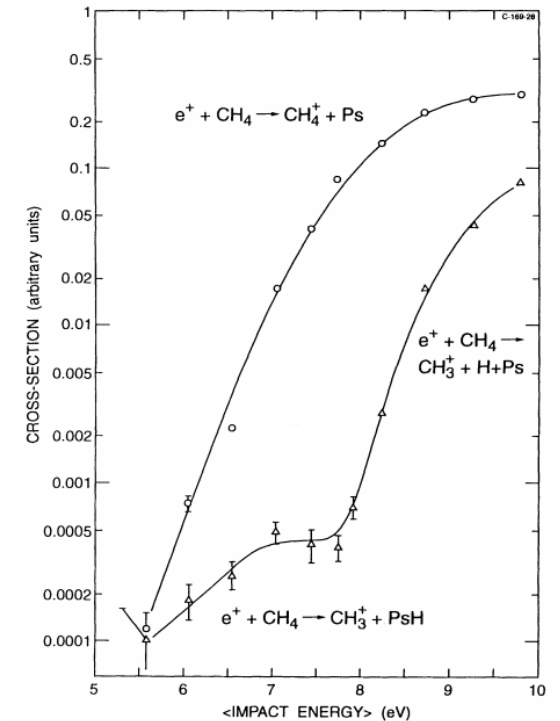
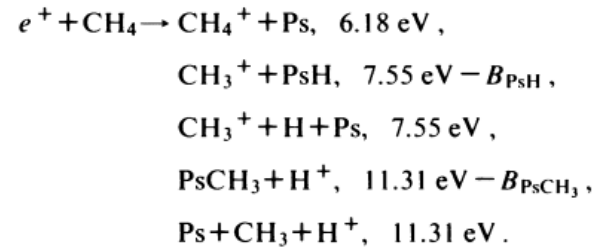
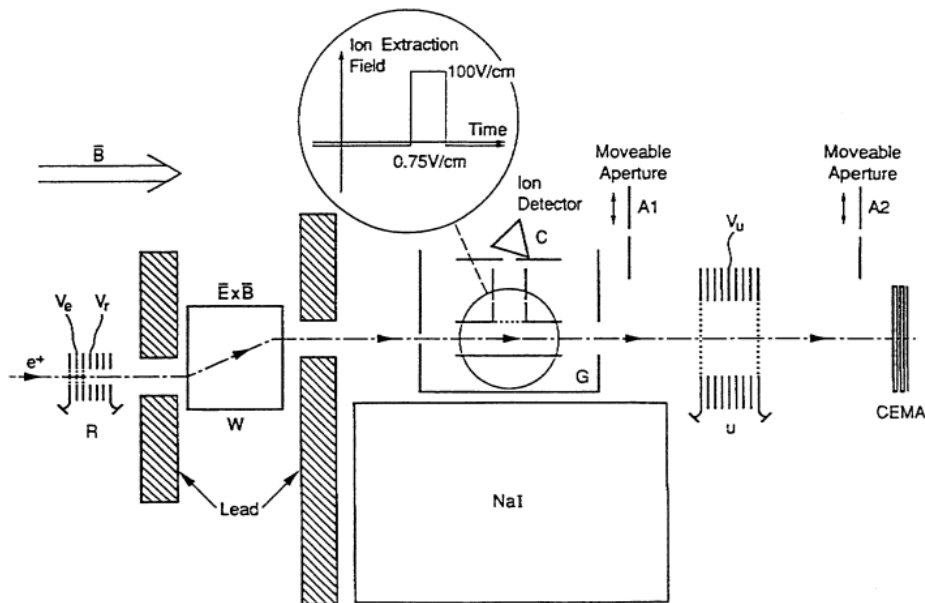
Positronium Ion (negative)

- Properties
 - Analog to the H^- ion
 - Binding energy ($Ps+e^-$) $0.33eV$
 - Lifetime 479 ps
 - 2γ annihilation
- Conversion efficiency from C foil 0.028%
- Recently Nagashima and co-workers have produced efficient Ps^- -converters using Na coated W (1.5%)



Positronium Hydride

- We have seen that Ps atoms behave much like H, forming Ps^- and Ps_2 what about binding to other atoms and molecules?



$$\text{BE} = (1.1 \pm 0.2) \text{ eV}$$



Conclusion

- From the first prediction of Mohorovičić in 1934 the field of positronium physics has been an active and interesting.
- Positronium is the most common pathway for positron annihilation in vacuum and some solids.
- Of the so called 'polyelectrons' Wheeler predicted Ps , Ps^- and Ps_2 have been produced (Ps^+ is more difficult)
- Positronium bound to atoms is an active field of research but at present only theory has been presented since the measurement of Schrader *et al*

Reviews

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- M. Charlton and J. W. Humberston *'Positron Physics'* Cambridge University Press (2000)
- P. Coleman *'Positron Beams and their applications'* World Scientific (2000)
- G. Laricchia, S. Armitage, A. Kover, D. J. Murtagh *Ad. At. Mol. Opt. Phys.* **56** 1 (2008)
- D. B. Cassidy *Euro J. Phys. D.* 71 **72** (2018)