

More on Galaxy Classification

Trends within the Hubble Sequence

- E0 --> S0 --> Sb
 - Decreasing bulge to disk ratio
 - Decreasing stellar age
 - Increasing gas content
 - Increasing star formation rate

Problems

- Constructed to classify massive galaxies
- Spiral parameters not well defined in the sequence
- Bars are yes/no, observations show that bars for a continuum of strengths
- Works best for isolated systems, in clusters classifications can be difficult

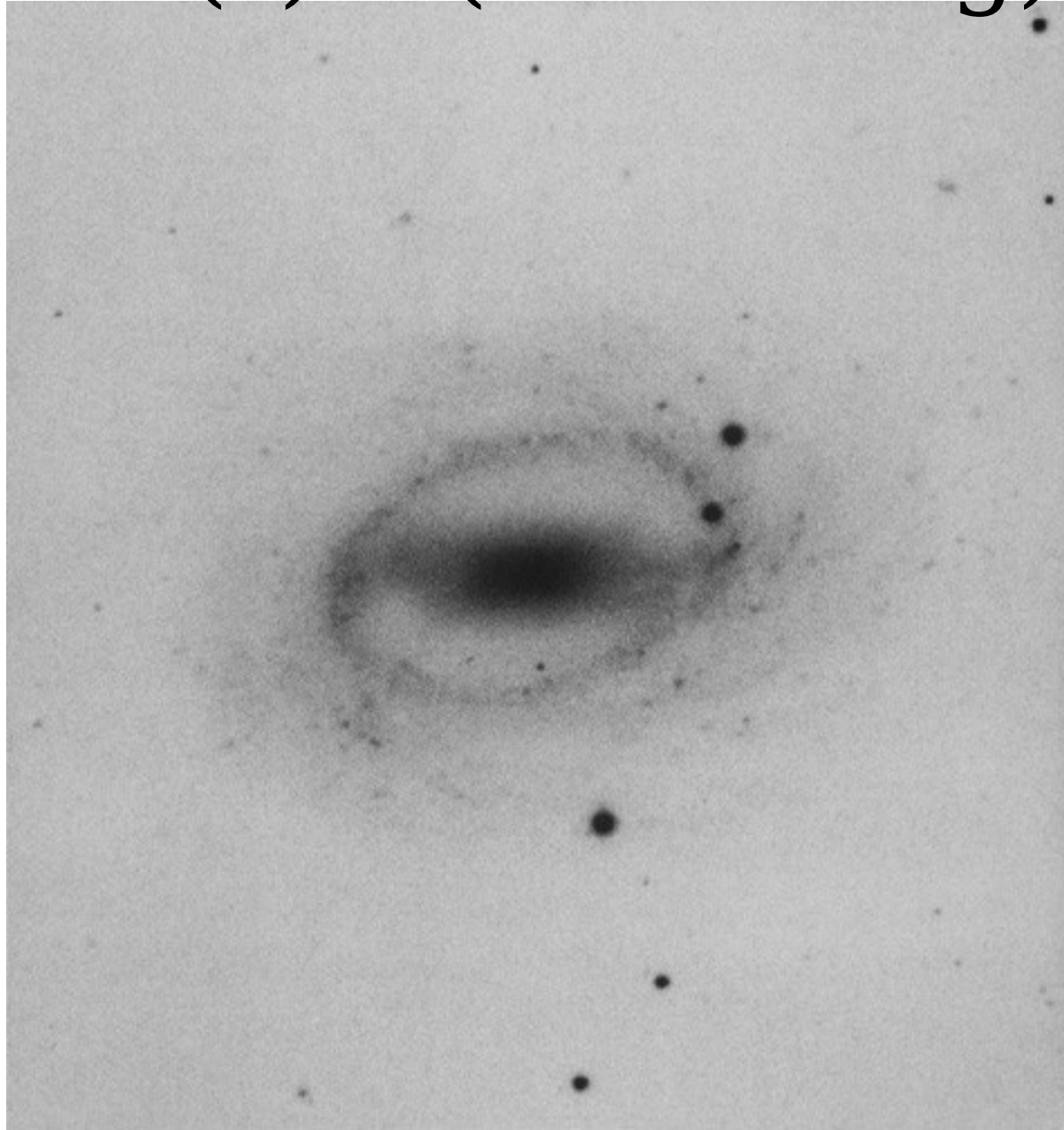
de Vaucouleurs' revision

- Mixed types
 - Sab , Scd, S0/a, E/S0 for intermediate types
 - S no bar, SB strong bar, SAB intermediate bar
- Inner rings
 - Arms in ring (r)
 - Arms out of ring (s)
 - And (rs)

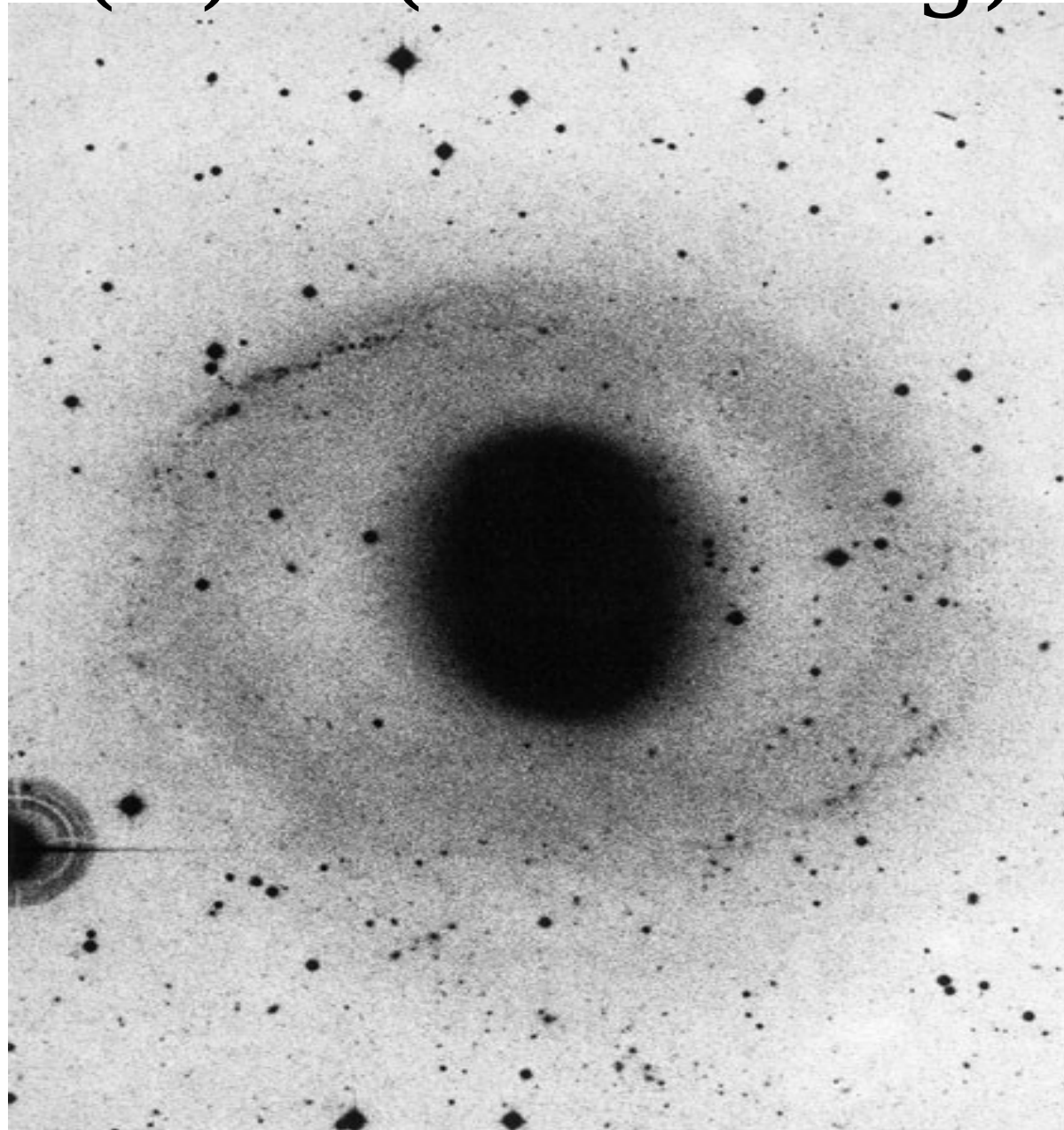
de Vaucouleurs' revision cont.

- Outer ring (R)S
- Extend the Spiral/Irr type
 - Sm (LMC type spiral)
 - Sd (very small bulge)
 - Sdm (intermediate between Sd and Im)
 - Im (Irr class)
 - Sm (Magallanic type spirals)
 - LMC is SB(s)m

IC 5240 example of a SB(r)ab(inner ring)



NGC 1543 example of a (R)SB(Outer ring)



NGC 6783 example of a (R)SB(r)

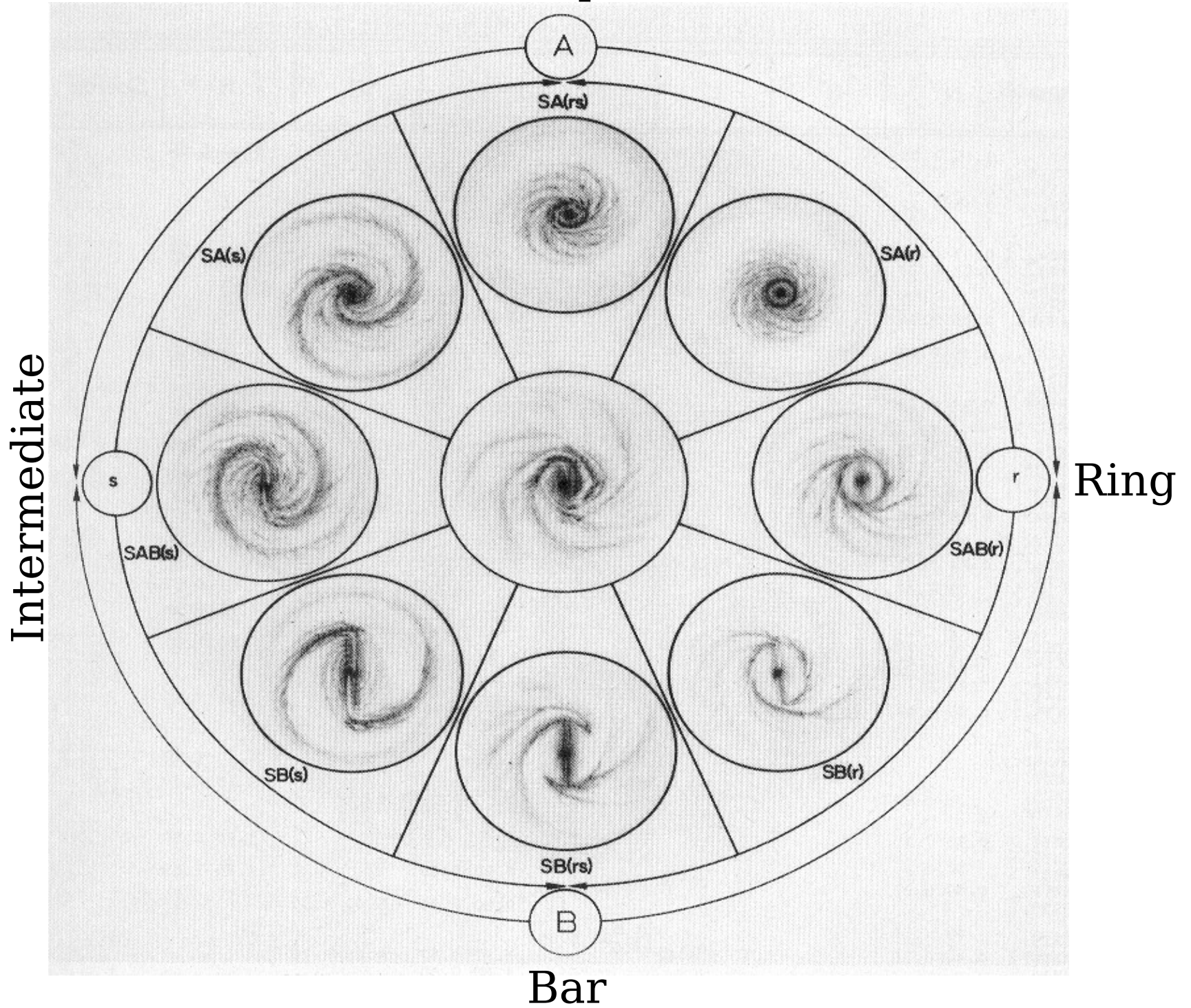


De Vaucouleurs also introduced the T type

E0	-->	S0	-->	Sa	-->	Sb	-->	Sc	-->	Im
-5		-1		1		3		5		10

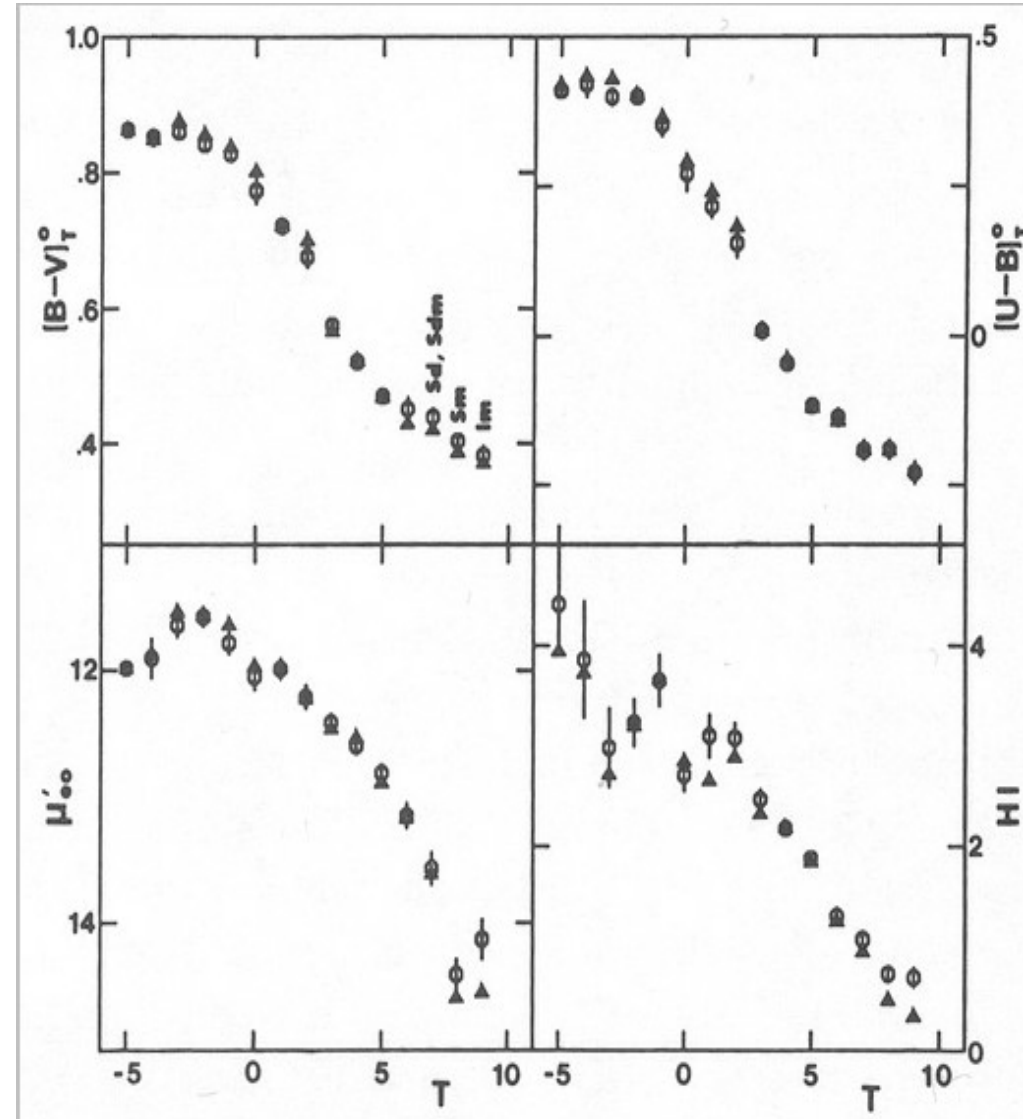
so that computers could be used with his
catalogs (RC1 1964, RC2 1974, RC3 1991)

Pure Spiral



Advantages of the new system

- Classes are more continuous
- Can classify up to 97% of all galaxies without special bins
- Describes features that are clues to the dynamics
- In wide use



Limitations of the new system

- $E \rightarrow I_m$ is not a linear sequence of 1 parameter
- Terms are not universally applicable
 - What do ring or bars mean for ellipticals?
- Visual system
 - Not based of a physical parameter (mass, luminosity etc)
- Parameters are not always independent
 - Rings and bars

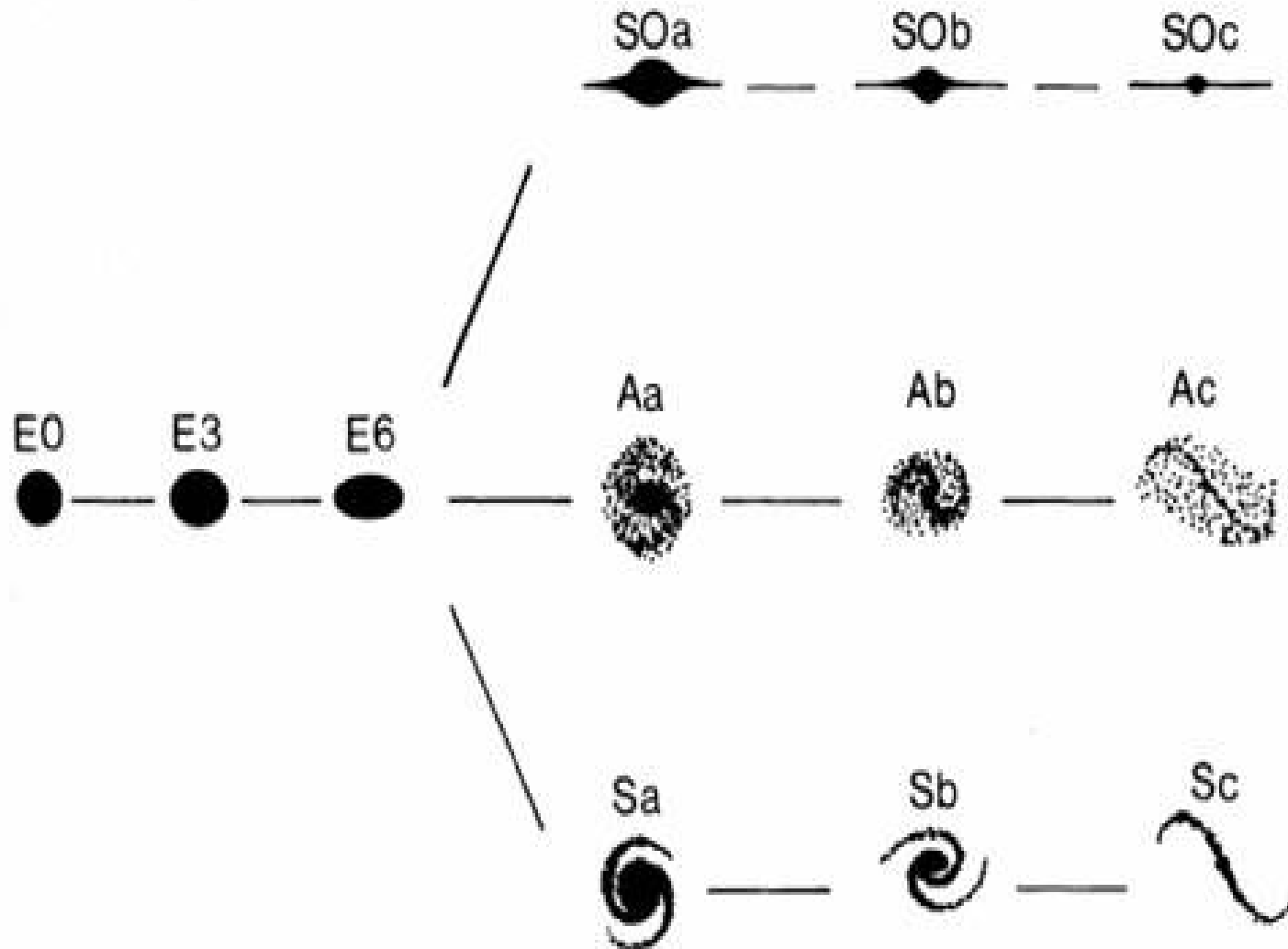
Modifications

- DDO system (Van den Bergh 1960)
 - Added luminosity as a parameter
 - Star formation, gas content, spiral arm “development”
 - Mostly added subclasses to spirals
 - Sc I - well developed arms
 - Sc III - short stubby arms
 - Sc IV - faint spiral structure (LMC)

Modifications cont.

- DDO system (Van den Bergh 1976)
 - Added Anemic as a class
 - Lower gas content than galaxies with similar luminosity
 - Lower star formation rates
 - Spiral arms less well developed

DDO classification of Galaxies

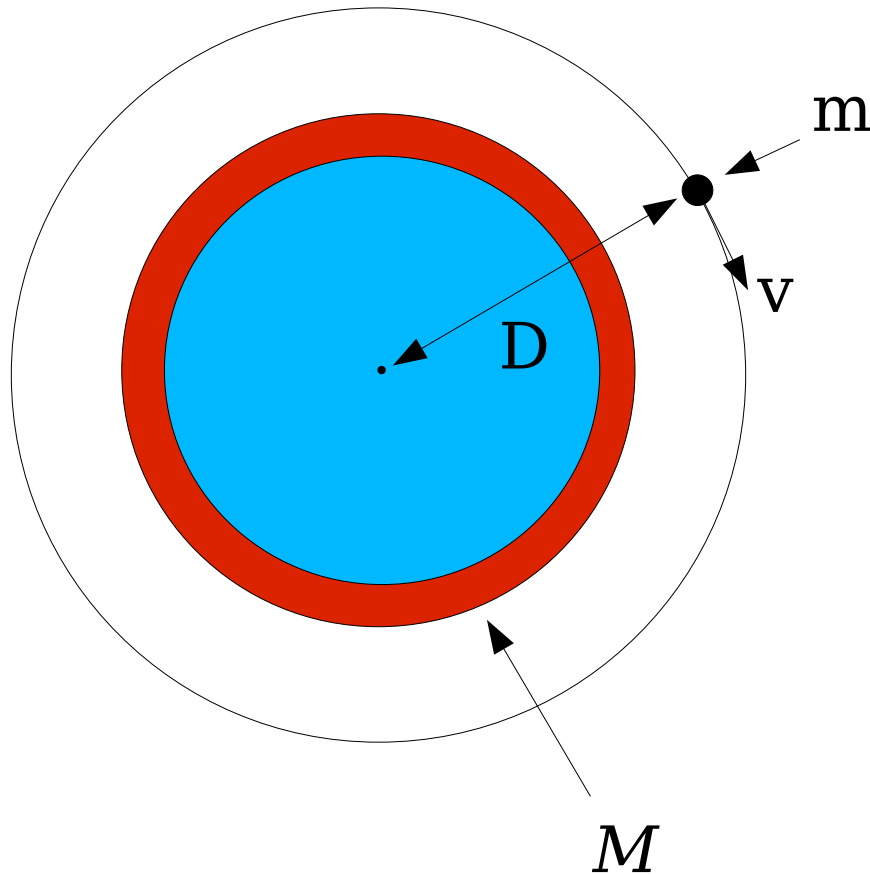


Problems with the DDO system

- Complicated
- Anemic is not necessarily an intrinsic parameter
 - Ram pressure stripping
 - Merger activity

As a result it is not as commonly used as the de Vaucouleurs system

Tidal Stripping

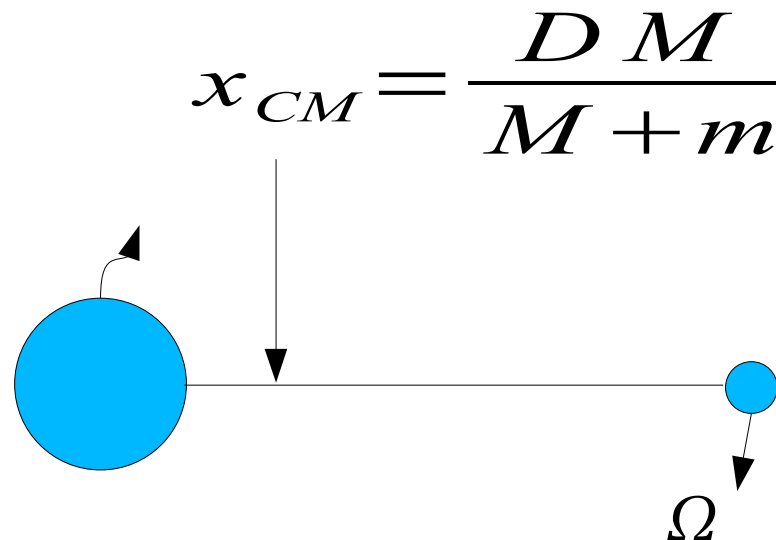


If the mass m is close enough to the particles in M then the particles closest to M are at risk of being removed or stripped from the larger body.

Tidal Stripping cont.

If we select a frame of reference that rotates at the same rate as the satellite m ($\Omega = v/(2\pi D)$). This is so that we have a stationary problem.

Lets also look at this in Center of Mass coordinates.



We can write the effective potential in the form

$$\Phi_{eff}(x) = -\frac{GM}{|D-x|} - \frac{Gm}{|x|} - \frac{\Omega^2}{2} \left(x - \frac{DM}{M+m} \right)^2$$

Normal gravitational force and angular momentum

This potential will have 3 maxima and we can find these by

$$\frac{\partial \Phi_{eff}}{\partial x} = -\frac{GM}{(D-x)^2} \pm \frac{Gm}{x^2} - \Omega^2 \left(x - \frac{DM}{M+m} \right) = 0$$

But remember for a circular orbit

$$V^2 = \frac{GM}{r} \quad \text{and the acceleration is} \quad a = \frac{V^2}{r} = \frac{GM}{r^2}$$

Since $\Omega = v/(2\pi r)$ then $a = \frac{V^2}{r} = \Omega^2 r$

But since we are in the CM coordinate system

$$a = \Omega^2 r = \Omega^2 \frac{DM}{(M+m)} = \frac{GM}{D^2}$$

Solving this for Ω^2

$$\Omega^2 = \frac{G(M+m)}{D^3}$$

Substituting this for Ω^2 below

$$\frac{\partial \Phi_{eff}}{\partial x} = -\frac{GM}{(D-x)^2} \pm \frac{Gm}{x^2} - \Omega^2 \left(x - \frac{DM}{M+m} \right) = 0$$

We get

$$-\frac{GM}{(D-x)^2} \pm \frac{Gm}{x^2} - \frac{G(M+m)}{D^3} \left(x - \frac{DM}{M+m} \right) = 0$$

If $m \ll M$ then $x \ll D$ we can rewrite our equation as

$$\frac{GM}{D \left(1 - \frac{x}{D}\right)^2} \pm \frac{Gm}{x^2} - \frac{GM}{D^2} + \frac{M+m}{D^3} = 0$$

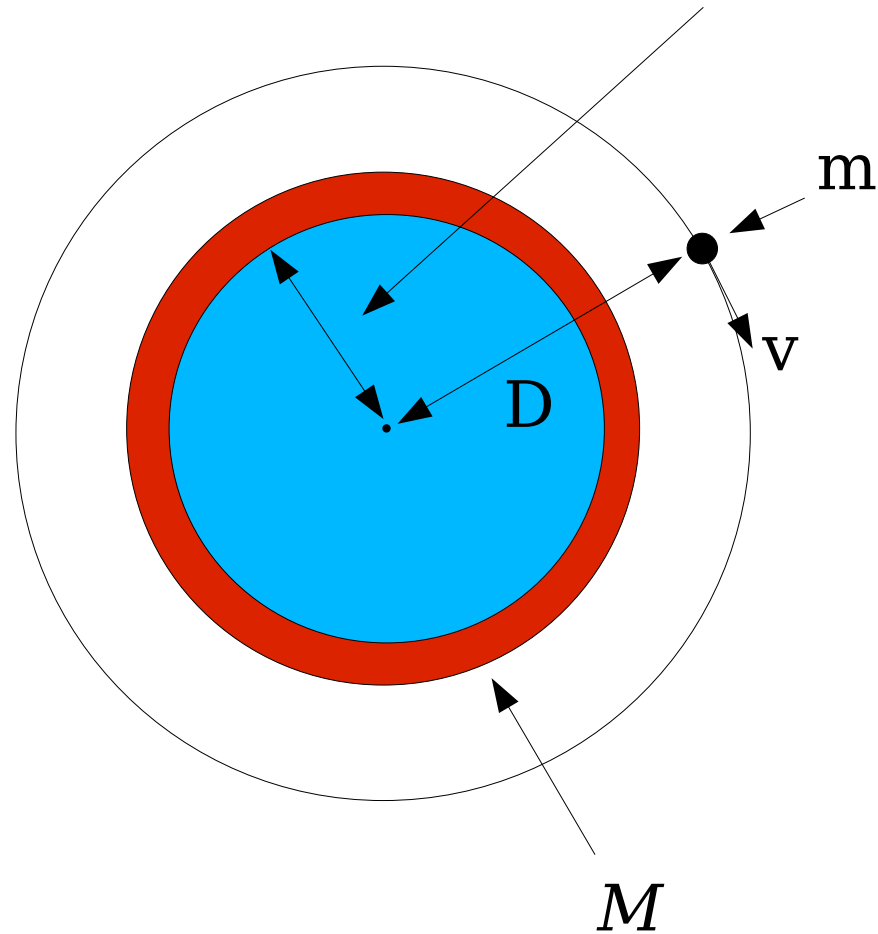
Expanding $(1 - x/D)^{-2}$ in a Taylor series

$$\frac{GM}{D^2} \left(1 + \frac{2x}{D} + \dots\right) \pm \frac{Gm}{x^2} - \frac{GM}{D^2} + \frac{M+m}{D^3} = 0$$

Solving for x we get

$$x = \pm D \left[\frac{m}{M(3 + m/M)} \right]^{\frac{1}{3}} \approx \pm \left(\frac{m}{3M} \right)^{\frac{1}{3}} D$$

x is called the Jacobi limit or the Roche limit and is written as r_J



- This provides a crude estimate of the true tidal radius
 - In general the system will not have a circular orbit
 - We derived this for point masses and most systems are extended
 - If done in 3-d (so we get a surface) this is not a spherical surface
 - If $r_j > x$ a particle will not necessarily escape. Numerical studies show that there are some stable orbits up to $r = 2r_j$.

Cartesian Coordinates

$$\nabla^2 \psi(x, y, z) = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2}$$

Polar Coordinates

$$\nabla^2 \psi(\rho, \phi, z) = \frac{\partial^2 \psi}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial \psi}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{\partial^2 \psi}{\partial z^2}$$

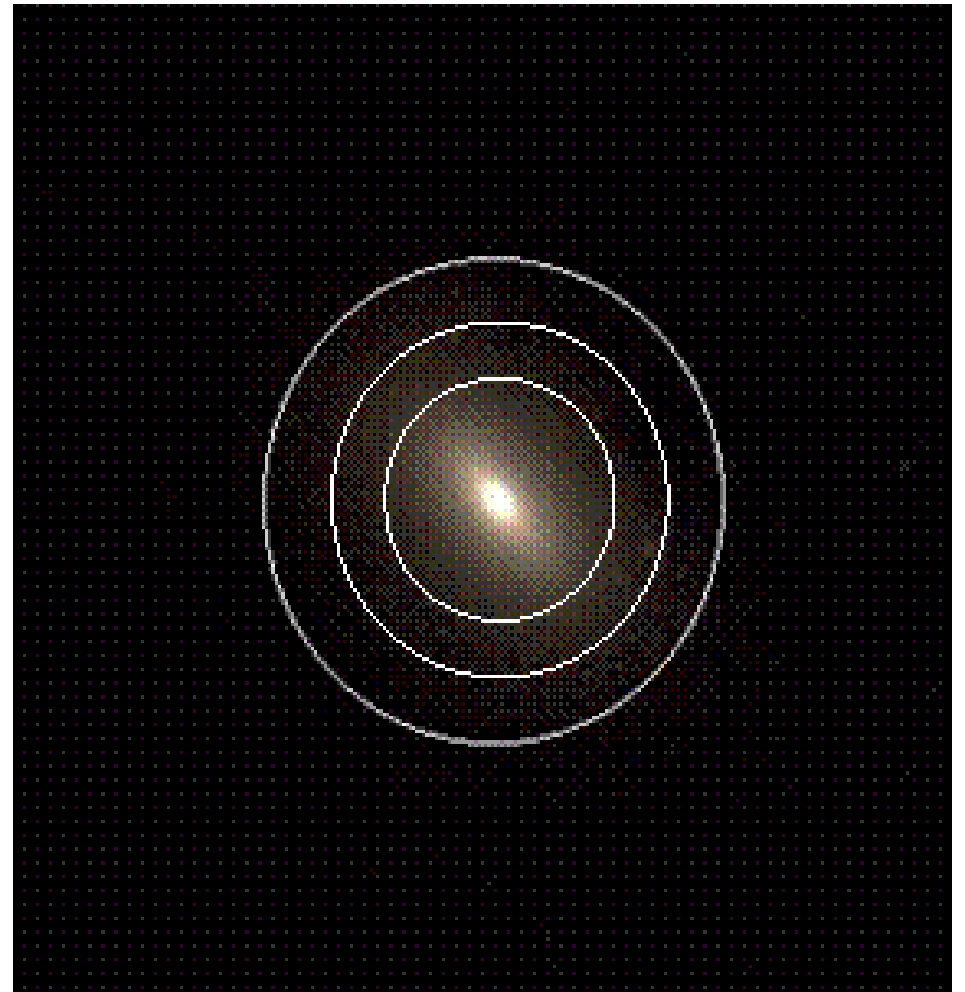
Spherical Coordinates

$$\nabla^2 \psi(r, \theta, \phi) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2}$$

Dust and Gas in External Galaxies

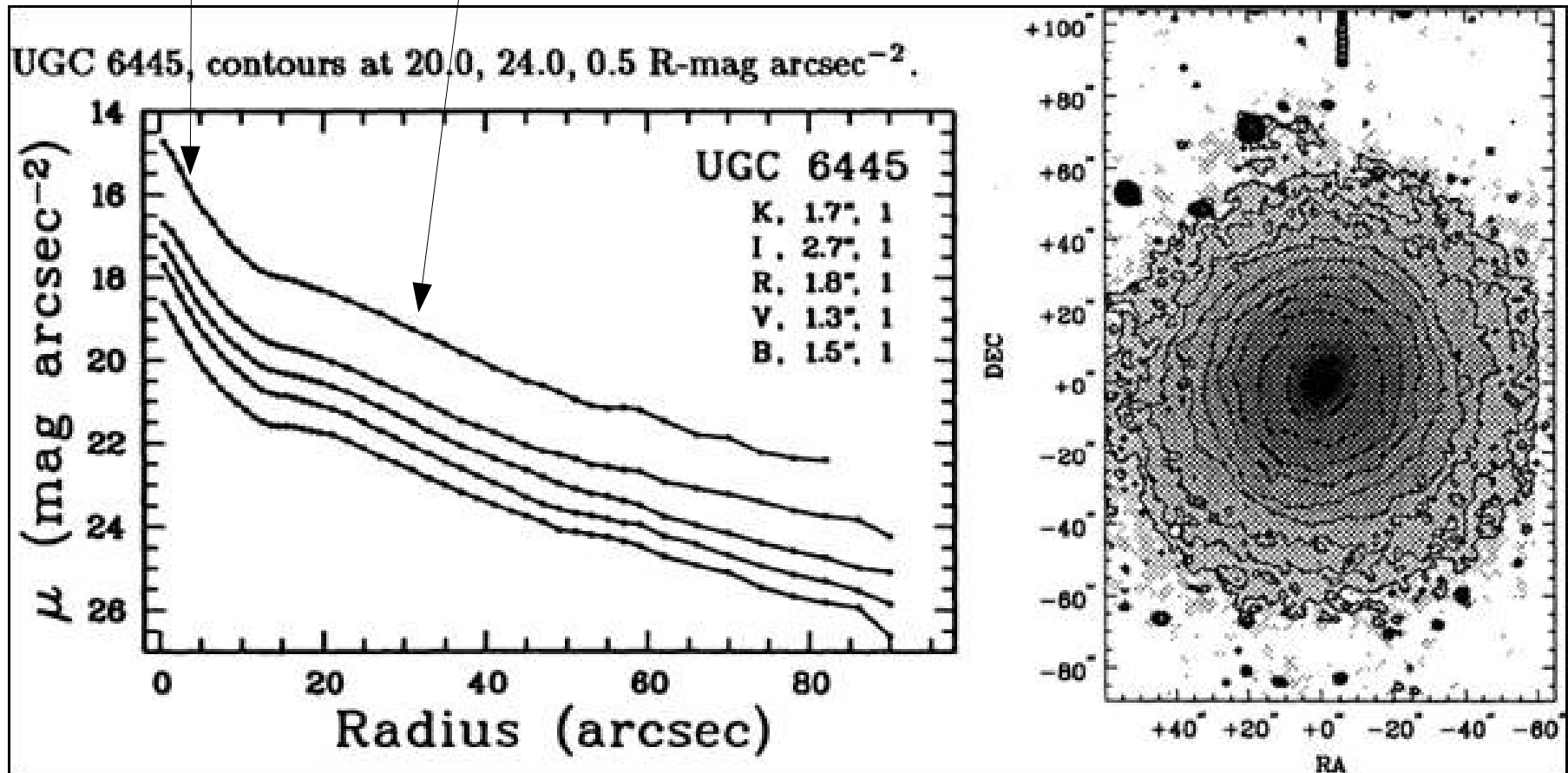
Galaxy Photometry

- Surface Photometry is often used to measure stellar distribution
- Measured in concentric radii (mag/sq arcsec)
- Use a fit function to measure SB
- Compare different anulii e.g. Concentration index, assymetry



Spiral galaxy profiles

Two components
Bulge Disk



Elliptical galaxy profiles

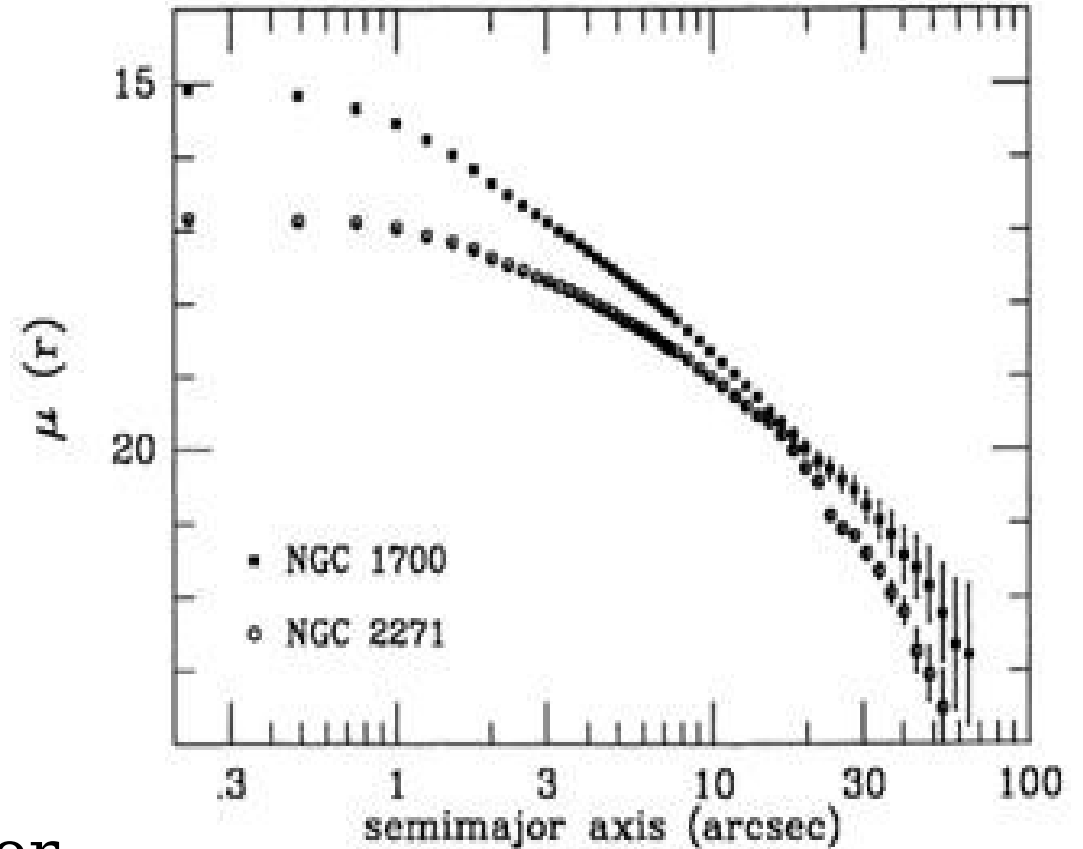
Often well fit with an exponential profile

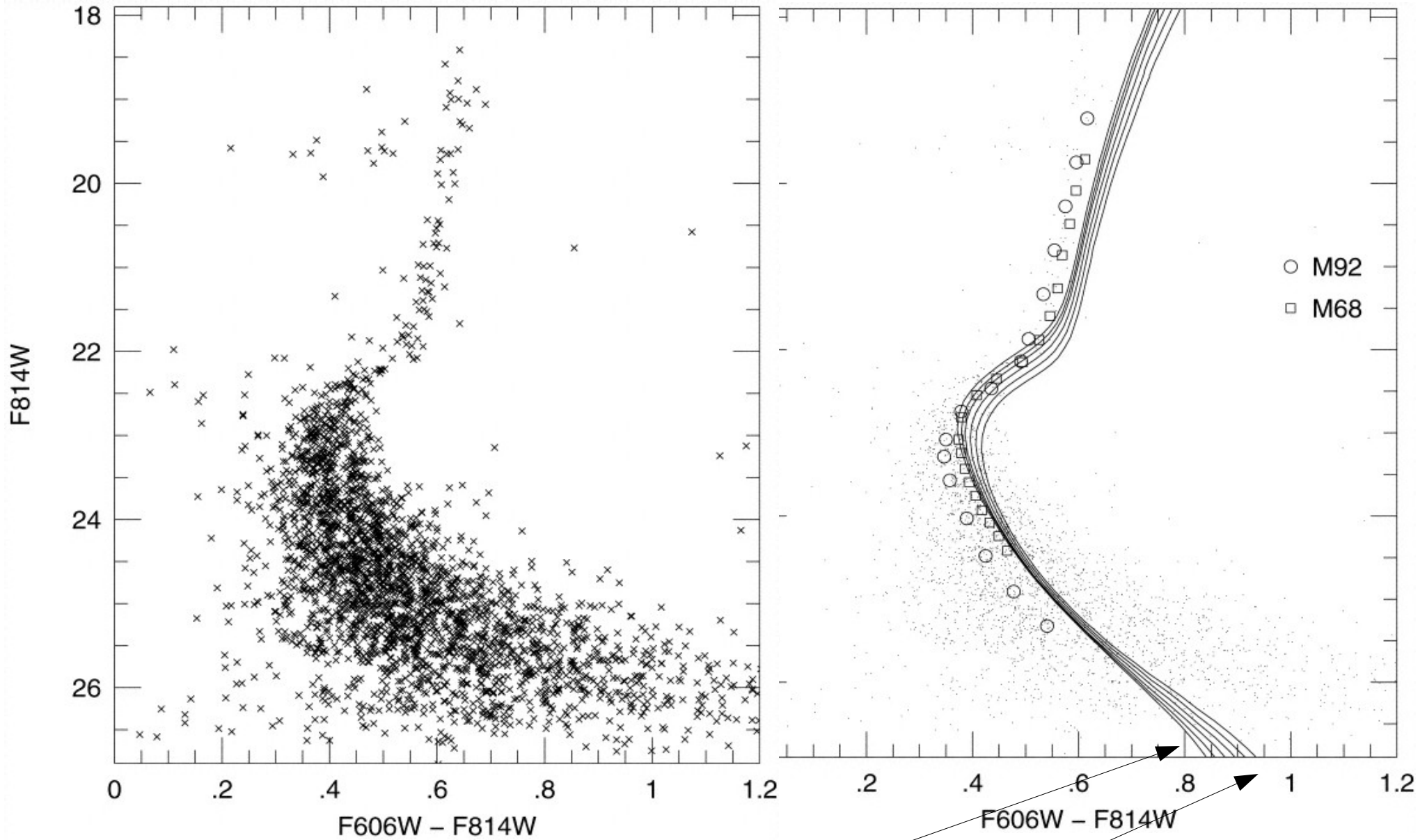
$$I(r) = I(r_e) \exp\left[-b \left(\frac{r}{r_e}\right)^{\frac{1}{n}} - 1\right]$$

If $n = 4$ this is a de Vaucouleur's law for general n Sersic's law.

Often flattens in the center

Fit spiral bulges as well





Measured and color magnitude diagram for the draco dSph the fit is for an age of 16 Gyr with metallicities $2.2 > [Fe/H] > 1.2$ (Grillmair et al. 1998)

