# Temperature dependence of the acentric factor of normal hydrogen, orthohydrogen and parahydrogen 

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## Research Article

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# Temperature dependence of the acentric factor of normal hydrogen, orthohydrogen and parahydrogen 

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#### Abstract

Temperature-dependence correlations of vapor pressure and acentric factor for normalhydrogen $\left(n-H_{2}\right)$, orthohydrogen $\left(o-H_{2}\right)$ and parahydrogen $\left(p-H_{2}\right)$, have been formulated. The obtained correlations are statistically excellent. The characteristic parameters such as the Pitzer's acentric factor, Riedel's parameter, Filippov's parameter have been determined for $n-H_{2}, o-H_{2}$ and $p-H_{2}$. And, the curvatures of vapor pressure curve for $n-H_{2}$, $o-H_{2}$ and $p-H_{2}$ have been determined in a wide range of temperature. It is found that the curvatures of vapor pressure curve for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ have a maximum at about $17.11 \mathrm{~K}, 17.12 \mathrm{~K}$ and 17.00 K , respectively.


## Keywords

Acentric factor, Filippov's parameter, normal hydrogen, orthohydrogen, parahydrogen, Riedel parameter

## Introduction

Hydrogen $\left(H_{2}\right)$ is a contender for alternative energy. $H_{2}$ fuel cell vehicles and $H_{2}$ based lowcarbon fuels will contribute to the decarbonisation of the mobility sector, shipping and aviation. $H_{2}$ is used as a rocket fuel. And, petroleum refining, semiconductor manufacturing, aerospace industry, fertilizer production, metal treatment, pharmaceutical, power plant generator, methanol production, commercial fixation of nitrogen from air, reduction of metallic ores use $\mathrm{H}_{2}$.
$o-H_{2}$ and $p-H_{2}$ are nuclear spin isomers of $H_{2}$. The nuclear spins of $o-H_{2}$ are parallel and that of $p-\mathrm{H}_{2}$ are antiparallel. They are chemically identical. They have the same atomic and isotopic structure. They differ in the nuclear spin states of their atoms. The energy difference associated with nuclear spin transitions of $H_{2}$ is about $0.1 \mathrm{~J} . \mathrm{mol}^{-1}$. However, this tiny change leads to different thermodynamic and spectroscopic properties of $H_{2}$ molecules. The $o-H_{2}$ and $p-H_{2}$ are characterized by different values of specific heat, boiling point and heat of vapor formation. It follows from the Pauli's principle that nuclear spin state and rotational state of the $H_{2}$ molecule are correlated. This is attributed to the fact that the molecules of these gases are rotating differently. Conversion between two nuclear spin states of $\mathrm{H}_{2}$ molecule occurs extremely slowly as the transitions between symmetric and antisymmetric nuclear spin states are forbidden by the selection rule of quantum mechanics. Hence, $p-H_{2}$ can be stored as a individual gas for longer periods. However, the use of paramagnetic catalysis promote the establishment of Boltzman's thermodynamic equilibrium between $o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$ states for a given temperature at accelerated rate. The paramagnetic materials create a strong inhomogeneous magnetic field at the atomic scale. In such fields, the two $\mathrm{H}_{2}$ isomers are no longer equivalent. Hence, the spin-flip transitions between o $-\mathrm{H}_{2}$ and $\mathrm{p}-\mathrm{H}_{2}$ are no longer forbidden. At room temperature, the $n-H_{2}$ at thermal equilibrium consists of $75 \% o-H_{2}$ and $25 \% p-H_{2}$. Knowledge about $o-H_{2}$ to $p-H_{2}$ conversion is important for the storage of liquid
$H_{2}$. Due to the energy difference associated with different rotational level, energy is released when $o-H_{2}$ converts to $p-H_{2}$ and energy is absorbed in the reverse process.

The scientific and technical significance has led to numerous experimental and theoretical studies on the thermodynamic properties of $H_{2}[1-14]$. The effect of $o-H_{2}$ and $p-\mathrm{H}_{2}$ composition on the performance of a proton exchange membrane fuel cell is calculated and experimentally studied[9]. Equation of state of $o-H_{2}$ and $p-H_{2}$ has been derived[10]. The influence of $o-\mathrm{H}_{2}$ and $\mathrm{p}-\mathrm{H}_{2}$ conversion is considered to recommend the parameters for $\mathrm{H}_{2}$ storage[11]. Sound velocity in liquid $p-H_{2}$, dielectric constant of liquid $p-\mathrm{H}_{2}$ along the saturation line, surface tension of $p-H_{2}$ in the temperature range from triple point to critical point and density of liquid $p-\mathrm{H}_{2}$ along saturation line have been determined[12]. In contrast to bulk metals, the nanoparticles of copper, silver and gold catalyze the low temperature $o-H_{2}$ to $p-H_{2}$ conversion[14]. $p-H_{2}$ is employed in the NMR and MRI signal enhancement[15].

The technological applications of hydrogen require the knowledge of its thermodynamic properties including its acentric factor. Acentric factor is a characteristic thermodynamic parameter of substances. It is a measure of nonsphericity of molecules[16,17]. It accounts[18] for the deviations in the thermodynamic properties of fluids consisting of non-spherical molecules from that of fluids made of spherical molecules. Acentric factor is a used as a parameter in the corresponding state principle[19]. The acentric factor is widely used in determining the thermodynamic properties of substances such as the compressibility factor[20-25], fluid phase equilibrium[26,27], virial coefficients[28,29], vapour pressure[30-32] and enthalpy of vaporization[17,33,34]. Hence, the knowledge of acentric factor of substances acquires significance. In recent years, several studies have been made on the acentric factors of various substances. Acentric factors are correlated[35] to molecular energies of $n$-alkanes. Artificial neural network group contribution method is used[36] to calculate the acentric factors of pure compounds. The values of acentric factors of limonene and linalool were optimized[37] to improve the performance of the SRE EoS. A new second order group contribution method has been developed[38] to predict the acentric factors of organic compounds. Using the normal boiling temperature, molecular weight and the number of atoms and bonds, empirical correlations are developed[39] to estimate the acentric factors of s-containing compounds. Two different intelligent systems are used to estimate[40] the acentric factors of binary and ternary mixtures of ionic liquids. New generalized models are introduced[41] to estimate the acentric factors of pure compounds. Acentric factors of fluoroalkylsilane compounds have estimated[42] by a group contribution method. Acentric factor of carbon dioxide is required[43] by the cubic equations of state to predict the solid solute solubility in supercritical carbon dioxide. However, the studies on the acentric factors of $H_{2}$ and their isomers, particularly their temperature dependence, are scarce.

This work formulates the temperature-dependence correlations for the acentric factor, vapour pressure and the curvature of vapour pressure curve for $n-H_{2}, o-H_{2}, p-H_{2}$ over a wide range of temperature upto the vicinity of their liquid-vapour critical point.

## Vapor pressure

Analysis of the vapour pressure $P_{v p}$ data[44,45] for $n-H_{2}$, o- $H_{2}$, and $p-H_{2}$ shows that its temperature-dependence may be represented by a second degree polynomial. That is,

$$
\begin{equation*}
P_{v p}=A_{1}+A_{2} T+A_{3} T^{2} \tag{1}
\end{equation*}
$$

where $P_{v p}$ is in MPa and $T$ is in K .
The coefficients in Eq.(1) are $A_{1}=1.3422, A_{2}=-0.1560$ and $A_{3}=0.0046$. The correlation given by Eq.(1) is characterized by the correlation coefficient $(R)$ of 0.9976 and the
coefficient of determination $\left(R^{2}\right)$ of 0.9952 , respectively. The vapor pressure correlation given by Eq.(1) is depicted in Fig. 1 along with the literature data[44,45].


Fig. 1 - Vapour pressure of $n-H_{2}, o-H_{2}$ and $p-H_{2}$.
Eq.(1) gives the temperature derivatives of vapour pressure as
$\frac{d P}{d T}=A_{2}+2 A_{3} T$
$\frac{d^{2} P}{d T^{2}}=2 A_{3}$
For, $n-H_{2}, o-H_{2}$ and $p-H_{2}$, the derivatives $d P / d T$ and $d^{2} P / d T^{2}$ are determined by Eqs.(2) and (3), respectively. The results are presented in the reduced coordinates $T^{*}=T / T_{c}$ and $P^{*}=P / P_{c}$, in Tables 1-3 and depicted in Fig. 2


Fig. $2-d P^{*} / d T^{*}$ versus $T^{*}$ for $n-H_{2}, o-H_{2}$ and $p-H_{2}$.

## Curvature of vapor pressure curve

The vapour pressure curve's curvature at a particular temperature is a measure of how the surface bends away from its tangent plane at this point. The curvature of the vapour pressure curve is examined for the proposed $P_{v p}$ correlation. The curvature (inverse of the radius of curvature) of the vapour pressure curve is defined[46] as
$k=\left(d^{2} P^{*} / d T^{* 2}\right)\left[1+\left(\frac{d P^{*}}{d T^{*}}\right)^{2}\right]^{-3 / 2}$
$\frac{P_{c}}{T_{c}} \frac{d P^{*}}{d T^{*}}=\frac{d P}{d T}=A_{2}+2 A_{3} T$
$\frac{P_{c}}{T_{c}} \frac{d^{2} P^{*}}{d T^{2^{*}}}=\frac{d^{2} P}{d T^{2}}=2 A_{3}$
For $n-H_{2}, o-H_{2}$ and $p-H_{2}$, the curvature $k$ is determined by Eq.(4) using the values of $d P / d T$ and $d^{2} P / d T^{2}$ tabulated in Tables 1-3. The results are also given in Tables 1-3 and depicted in Fig. 3. The curvature of vapour pressure for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ is found have a maximum at temperatures $0.5162 T_{c}, 0.5153 T_{c}$, and $0.5161 T_{c}$, respectively.


Fig. 3 - Curvature of the vapour pressure curve for $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$.
Table 1 - Curvature of vapour pressure curve $\boldsymbol{k}$ and acentric factor $\boldsymbol{\omega}$ of $\boldsymbol{n}-\boldsymbol{H}_{\mathbf{2}}$.

| $T^{*}$ |  | $d P^{*} / d T^{*}$ | $d^{2} P^{*} / d T^{2 *}$ | $k$ | $\omega$ | $\omega$ | $d \omega / d T^{*}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Eq.(5) | Eq.(6) | Eq.(4) | $d^{2} \omega / d T^{* 2}$ <br> Eq.(8) | $\omega(9)$ <br> Eq.(9) | Eq.(10) | Eq.(11) |  |
| 0.4555 | -0.4369 | 6.0138 | 4.6272 | 0.9798 | 0.941 | -5.7761 | 8.4398 |
| 0.4707 | -0.3187 | 6.0138 | 5.2015 | 0.8744 | 0.8544 | -5.6481 | 8.4398 |
| 0.4859 | -0.2004 | 6.0138 | 5.6688 | 0.7749 | 0.7697 | -5.5201 | 8.4398 |
| 0.5010 | -0.0822 | 6.0138 | 5.9533 | 0.6807 | 0.6869 | -5.3921 | 8.4398 |
| 0.5162 | 0.0361 | 6.0138 | 6.0020 | 0.5915 | 0.6061 | -5.2641 | 8.4398 |
| 0.5314 | 0.1543 | 6.0138 | 5.8052 | 0.5067 | 0.5272 | -5.1361 | 8.4398 |
| 0.5466 | 0.2726 | 6.0138 | 5.4008 | 0.4260 | 0.4503 | -5.0080 | 8.4398 |


| 0.5617 | 0.3908 | 6.0138 | 4.8590 | 0.3492 | 0.3753 | -4.8800 | 8.4398 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0.5769 | 0.5091 | 6.0138 | 4.2561 | 0.2758 | 0.3023 | -4.7520 | 8.4398 |
| 0.5921 | 0.6274 | 6.0138 | 3.6555 | 0.2056 | 0.2311 | -4.6240 | 8.4398 |
| 0.6072 | 0.7456 | 6.0138 | 3.0986 | 0.1384 | 0.162 | -4.4960 | 8.4398 |
| 0.6224 | 0.8639 | 6.0138 | 2.6060 | 0.0740 | 0.0948 | -4.3679 | 8.4398 |
| 0.6376 | 0.9821 | 6.0138 | 2.1840 | 0.0122 | 0.0295 | -4.2399 | 8.4398 |
| 0.6527 | 1.1004 | 6.0138 | 1.8294 | -0.0472 | -0.0339 | -4.1119 | 8.4398 |
| 0.6679 | 1.2186 | 6.0138 | 1.5351 | -0.1044 | -0.0953 | -3.9839 | 8.4398 |
| 0.6831 | 1.3369 | 6.0138 | 1.2923 | -0.1595 | -0.1547 | -3.8559 | 8.4398 |
| 0.6982 | 1.4551 | 6.0138 | 1.0926 | -0.2126 | -0.2122 | -3.7279 | 8.4398 |
| 0.7134 | 1.5734 | 6.0138 | 0.9281 | -0.2639 | -0.2678 | -3.5998 | 8.4398 |
| 0.7286 | 1.6916 | 6.0138 | 0.7925 | -0.3135 | -0.3214 | -3.4718 | 8.4398 |
| 0.7437 | 1.8099 | 6.0138 | 0.6802 | -0.3615 | -0.3731 | -3.3438 | 8.4398 |
| 0.7589 | 1.9282 | 6.0138 | 0.5869 | -0.4079 | -0.4229 | -3.2158 | 8.4398 |
| 0.7741 | 2.0464 | 6.0138 | 0.5089 | -0.4529 | -0.4707 | -3.0878 | 8.4398 |
| 0.7892 | 2.1647 | 6.0138 | 0.4436 | -0.4966 | -0.5165 | -2.9598 | 8.4398 |
| 0.8044 | 2.2829 | 6.0138 | 0.3884 | -0.5390 | -0.5605 | -2.8317 | 8.4398 |
| 0.8196 | 2.4012 | 6.0138 | 0.3417 | -0.5802 | -0.6025 | -2.7037 | 8.4398 |
| 0.8347 | 2.5194 | 6.0138 | 0.3019 | -0.6203 | -0.6425 | -2.5757 | 8.4398 |
| 0.8499 | 2.6377 | 6.0138 | 0.2679 | -0.6593 | -0.6806 | -2.4477 | 8.4398 |
| 0.8651 | 2.7559 | 6.0138 | 0.2386 | -0.6973 | -0.7168 | -2.3197 | 8.4398 |
| 0.8803 | 2.8742 | 6.0138 | 0.2134 | -0.7343 | -0.751 | -2.1916 | 8.4398 |
| 0.8954 | 2.9925 | 6.0138 | 0.1915 | -0.7705 | -0.7832 | -2.0636 | 8.4398 |
| 0.9106 | 3.1107 | 6.0138 | 0.1724 | -0.8058 | -0.8136 | -1.9356 | 8.4398 |
| 0.9258 | 3.2290 | 6.0138 | 0.1557 | -0.8404 | -0.842 | -1.8076 | 8.4398 |
| 0.9409 | 3.3472 | 6.0138 | 0.1411 | -0.8742 | -0.8684 | -1.6796 | 8.4398 |
| 0.9561 | 3.4655 | 6.0138 | 0.1282 | -0.9074 | -0.8929 | -1.5516 | 8.4398 |
| 0.9713 | 3.5837 | 6.0138 | 0.1168 | -0.9400 | -0.9155 | -1.4235 | 8.4398 |
| 0.9864 | 3.7020 | 6.0138 | 0.1066 | -0.9721 | -0.9361 | -1.2955 | 8.4398 |

Table 2 - Curvature of vapour pressure curve $\boldsymbol{k}$ and acentric factor $\omega$ of $\boldsymbol{o}-\boldsymbol{H}_{2}$.

| $T^{*}$ | $d P^{*} / d T^{*}$ <br> Eq.(5) | $d^{2} P^{*} / d T^{2 *}$ <br> Eq.(6) | $k$ <br> Eq.(4) | $\omega$ <br> Eq.(8) | $\omega$ <br> Eq.(9) | $d \omega / d T^{*}$ <br> Eq.(10) | $d^{2} \omega / d T^{* 2}$ <br> Eq.(11) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.4548 | -0.4311 | 5.9104 | 4.5770 | 0.9845 | 0.9454 | -5.7825 | 8.4398 |
| 0.4699 | -0.3138 | 5.9104 | 5.1337 | 0.8792 | 0.8588 | -5.6548 | 8.4398 |
| 0.4851 | -0.1965 | 5.9104 | 5.5837 | 0.7797 | 0.7742 | -5.5270 | 8.4398 |
| 0.5002 | -0.0793 | 5.9104 | 5.8551 | 0.6856 | 0.6915 | -5.3992 | 8.4398 |
| 0.5153 | 0.0380 | 5.9104 | 5.8976 | 0.5964 | 0.6107 | -5.2715 | 8.4398 |
| 0.5305 | 0.1553 | 5.9104 | 5.7029 | 0.5116 | 0.5319 | -5.1437 | 8.4398 |
| 0.5456 | 0.2725 | 5.9104 | 5.3081 | 0.4310 | 0.455 | -5.0159 | 8.4398 |
| 0.5608 | 0.3898 | 5.9104 | 4.7804 | 0.3396 | 0.38 | -4.8882 | 8.4398 |
| 0.5759 | 0.5071 | 5.9104 | 4.1932 | 0.2807 | 0.307 | -4.7604 | 8.4398 |
| 0.5910 | 0.6243 | 5.9104 | 3.6073 | 0.2105 | 0.2359 | -4.6326 | 8.4398 |
| 0.6062 | 0.7416 | 5.9104 | 3.0628 | 0.1433 | 0.1667 | -4.5049 | 8.4398 |
| 0.6213 | 0.8589 | 5.9104 | 2.5803 | 0.0789 | 0.0995 | -4.3771 | 8.4398 |
| 0.6364 | 0.9761 | 5.9104 | 2.1658 | 0.0170 | 0.0342 | -4.2494 | 8.4398 |
| 0.6516 | 1.0934 | 5.9104 | 1.8168 | -0.0425 | -0.0291 | -4.1216 | 8.4398 |
| 0.6667 | 1.2107 | 5.9104 | 1.5265 | -0.0997 | -0.0906 | -3.9938 | 8.4398 |
| 0.6819 | 1.3279 | 5.9104 | 1.2866 | -0.1549 | -0.1501 | -3.8661 | 8.4398 |
| 0.6970 | 1.4452 | 5.9104 | 1.0888 | -0.2081 | -0.2076 | -3.7383 | 8.4398 |


| 0.7121 | 1.5625 | 5.9104 | 0.9258 | -0.2594 | -0.2632 | -3.6105 | 8.4398 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0.7273 | 1.6797 | 5.9104 | 0.7911 | -0.3091 | -0.3169 | -3.4828 | 8.4398 |
| 0.7424 | 1.7970 | 5.9104 | 0.6795 | -0.3571 | -0.3687 | -3.3550 | 8.4398 |
| 0.7575 | 1.9143 | 5.9104 | 0.5867 | -0.4036 | -0.4185 | -3.2272 | 8.4398 |
| 0.7727 | 2.0315 | 5.9104 | 0.5091 | -0.4487 | -0.4664 | -3.0995 | 8.4398 |
| 0.7878 | 2.1488 | 5.9104 | 0.4439 | -0.4924 | -0.5124 | -2.9717 | 8.4398 |
| 0.8030 | 2.2661 | 5.9104 | 0.3889 | -0.5348 | -0.5564 | -2.8439 | 8.4398 |
| 0.8181 | 2.3834 | 5.9104 | 0.3423 | -0.5760 | -0.5985 | -2.7162 | 8.4398 |
| 0.8332 | 2.5006 | 5.9104 | 0.3026 | -0.6161 | -0.6386 | -2.5884 | 8.4398 |
| 0.8484 | 2.6179 | 5.9104 | 0.2686 | -0.6552 | -0.6768 | -2.4606 | 8.4398 |
| 0.8635 | 2.7352 | 5.9104 | 0.2393 | -0.6932 | -0.7131 | -2.3329 | 8.4398 |
| 0.8787 | 2.8524 | 5.9104 | 0.2140 | -0.7302 | -0.7475 | -2.2051 | 8.4398 |
| 0.8938 | 2.9697 | 5.9104 | 0.1921 | -0.7663 | -0.7799 | -2.0773 | 8.4398 |
| 0.9089 | 3.0870 | 5.9104 | 0.1730 | -0.8016 | -0.8104 | -1.9496 | 8.4398 |
| 0.9241 | 3.2042 | 5.9104 | 0.1563 | -0.8361 | -0.8389 | -1.8218 | 8.4398 |
| 0.9392 | 3.3215 | 5.9104 | 0.1416 | -0.8699 | -0.8655 | -1.6940 | 8.4398 |
| 0.9543 | 3.4388 | 5.9104 | 0.1287 | -0.9030 | -0.8902 | -1.5663 | 8.4398 |
| 0.9695 | 3.5560 | 5.9104 | 0.1173 | -0.9355 | -0.9129 | -1.4385 | 8.4398 |
| 0.9846 | 3.6733 | 5.9104 | 0.1071 | -0.9674 | -0.9337 | -1.3107 | 8.4398 |

Table 3 - Curvature of vapor pressure curve $\boldsymbol{k}$ and acentric factor $\omega$ of $\boldsymbol{p}-\boldsymbol{H}_{2}$.

| $T^{*}$ | $d P^{*} / d T^{*}$ <br> Eq.(5) | $d^{2} P^{*} / d T^{2 *}$ <br> Eq.(6) | $k$ <br> Eq.(4) | $\omega$ <br> Eq.(8) | $\omega$ <br> Eq.(9) | $d \omega / d T^{*}$ <br> Eq.(10) | $d^{2} \omega / d T^{* 2}$ <br> Eq.(11) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.4190 | -0.7439 | 6.0372 | 3.1182 | 1.2615 | 1.1579 | -6.0848 | 8.4398 |
| 0.4250 | -0.6968 | 6.0372 | 3.3345 | 1.2124 | 1.1211 | -6.0335 | 8.4398 |
| 0.4554 | -0.4611 | 6.0372 | 4.5212 | 0.9804 | 0.9418 | -5.7773 | 8.4398 |
| 0.4706 | -0.3433 | 6.0372 | 5.1082 | 0.8748 | 0.8551 | -5.6492 | 8.4398 |
| 0.4858 | -0.2254 | 6.0372 | 5.6046 | 0.7752 | 0.7703 | -5.5211 | 8.4398 |
| 0.5009 | -0.1076 | 6.0372 | 5.9339 | 0.6809 | 0.6875 | -5.3930 | 8.4398 |
| 0.5161 | 0.0102 | 6.0372 | 6.0362 | 0.5916 | 0.6066 | -5.2648 | 8.4398 |
| 0.5313 | 0.1281 | 6.0372 | 5.8916 | 0.5067 | 0.5277 | -5.1367 | 8.4398 |
| 0.5465 | 0.2459 | 6.0372 | 5.5282 | 0.4260 | 0.4507 | -5.0086 | 8.4398 |
| 0.5617 | 0.3638 | 6.0372 | 5.0105 | 0.3491 | 0.3756 | -4.8805 | 8.4398 |
| 0.5768 | 0.4816 | 6.0372 | 4.4152 | 0.2756 | 0.3025 | -4.7524 | 8.4398 |
| 0.5920 | 0.5994 | 6.0372 | 3.8094 | 0.2054 | 0.2313 | -4.6243 | 8.4398 |
| 0.6072 | 0.7173 | 6.0372 | 3.2392 | 0.1382 | 0.1621 | -4.4961 | 8.4398 |
| 0.6224 | 0.8351 | 6.0372 | 2.7300 | 0.0738 | 0.0948 | -4.3680 | 8.4398 |
| 0.6376 | 0.9529 | 6.0372 | 2.2905 | 0.0120 | 0.0295 | -4.2399 | 8.4398 |
| 0.6527 | 1.0708 | 6.0372 | 1.9196 | -0.0475 | -0.0339 | -4.1118 | 8.4398 |
| 0.6679 | 1.1886 | 6.0372 | 1.6108 | -0.1046 | -0.0954 | -3.9837 | 8.4398 |
| 0.6831 | 1.3065 | 6.0372 | 1.3556 | -0.1598 | -0.1549 | -3.8556 | 8.4398 |
| 0.6983 | 1.4243 | 6.0372 | 1.1454 | -0.2129 | -0.2124 | -3.7274 | 8.4398 |
| 0.7135 | 1.5421 | 6.0372 | 0.9723 | -0.2642 | -0.268 | -3.5993 | 8.4398 |
| 0.7286 | 1.6600 | 6.0372 | 0.8295 | -0.3138 | -0.3217 | -3.4712 | 8.4398 |
| 0.7438 | 1.7778 | 6.0372 | 0.7114 | -0.3618 | -0.3734 | -3.3431 | 8.4398 |
| 0.7590 | 1.8956 | 6.0372 | 0.6132 | -0.4082 | -0.4232 | -3.2150 | 8.4398 |
| 0.7742 | 2.0135 | 6.0372 | 0.5313 | -0.4532 | -0.471 | -3.0869 | 8.4398 |
| 0.7894 | 2.1313 | 6.0372 | 0.4627 | -0.4969 | -0.5169 | -2.9587 | 8.4398 |
| 0.8045 | 2.2491 | 6.0372 | 0.4048 | -0.5393 | -0.5608 | -2.8306 | 8.4398 |
| 0.8197 | 2.3670 | 6.0372 | 0.3558 | -0.5805 | -0.6028 | -2.7025 | 8.4398 |


| 0.8349 | 2.4848 | 6.0372 | 0.3142 | -0.6206 | -0.6429 | -2.5744 | 8.4398 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0.8501 | 2.6027 | 6.0372 | 0.2785 | -0.6596 | -0.681 | -2.4463 | 8.4398 |
| 0.8653 | 2.7205 | 6.0372 | 0.2479 | -0.6976 | -0.7172 | -2.3182 | 8.4398 |
| 0.8804 | 2.8383 | 6.0372 | 0.2215 | -0.7346 | -0.7514 | -2.1900 | 8.4398 |
| 0.8956 | 2.9562 | 6.0372 | 0.1986 | -0.7708 | -0.7837 | -2.0619 | 8.4398 |
| 0.9108 | 3.0740 | 6.0372 | 0.1787 | -0.8061 | -0.814 | -1.9338 | 8.4398 |
| 0.9260 | 3.1918 | 6.0372 | 0.1613 | -0.8407 | -0.8424 | -1.8057 | 8.4398 |
| 0.9412 | 3.3097 | 6.0372 | 0.1461 | -0.8745 | -0.8688 | -1.6776 | 8.4398 |
| 0.9563 | 3.4275 | 6.0372 | 0.1326 | -0.9077 | -0.8933 | -1.5495 | 8.4398 |
| 0.9715 | 3.5454 | 6.0372 | 0.1208 | -0.9403 | -0.9158 | -1.4213 | 8.4398 |
| 0.9867 | 3.6632 | 6.0372 | 0.1103 | -0.9724 | -0.9364 | -1.2932 | 8.4398 |
| 1.0000 | 3.7664 | 6.0372 | 0.1020 | -1.0001 | -0.9529 | -1.1810 | 8.4398 |

## Acentric factor

Pitzer's acentric factor $\omega_{P}$ is defined[47] such that its value is zero for spherical molecules. That is,
$\omega_{P}=-1-\log _{10}\left(P_{v p}^{*}\right.$ at $\left.T^{*}=0.7\right)$
where $P_{v p}^{*}=P_{v p} / P_{c}$.
On the other hand, the acentric factor $\omega$ being a measure of deviation in the properties of nonspherical molecules from that of spherical molecules. This deviation depends on temperature. Hence, Eq.(7) may be generalized to get the temperature-dependent acentric factor as

$$
\begin{equation*}
\omega(T)=-1-\log _{10}\left(P_{v p}^{*}\right) \tag{8}
\end{equation*}
$$

Using the data on $P_{v p}^{*}$ obtained by the correlation given by Eq.(1) for $n-H_{2}, o-H_{2}$ and $p-H_{2}$, their acentric factors are calculated by Eq.(8). The results are presented in Tables 1-3 and depicted in Fig. 4


Fig. 4 - Acentric factor of $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$.

For $n-H_{2}, o-H_{2}$ and $p-H_{2}$, the values of $\omega$ may be fitted to a second degree polynomial in temperature. That is,

$$
\begin{equation*}
\omega=B_{1}+B_{2} T^{*}+B_{3} T^{* 2} \tag{9}
\end{equation*}
$$

The coefficients in Eq.(1) are $B_{1}=4.4480, B_{2}=-9.6208$ and $B_{3}=4.2199$. The correlation given by Eq.(9) is characterized by the correlation coefficient $(R)$ of 0.9992 and the coefficient of determination ( $R^{2}$ ) of 0.9985 , respectively.

Eq.(9) gives the temperature derivatives of acentric factor as
$\frac{d \omega}{d T^{*}}=B_{2}+2 B_{3} T^{*}$
$\frac{d^{2} \omega}{d T^{* 2}}=2 B_{3}$
For $n-H_{2}, o-H_{2}$ and $p-H_{2}$, the derivatives $d \omega / d T^{*}$ and $d^{2} \omega / d T^{* 2}$ are determined by Eq.(10) and (11), respectively. The results are presented in Tables 1-3 and depicted in Fig. 5


Fig. $5-d \omega / d T^{*}$ versus $T^{*}$ for $n-H_{2}, o-H_{2}$ and $p-H_{2}$.

## Riedel's parameter

Riedel's parameter $\alpha_{R}$ is a measure of the temperature-dependence of vapour pressure in the vapour-liquid critical region. It determines the slope of the fluid-phase equilibrium curve in the critical region. The Riedel's parameter $\alpha_{R}$ is correlated[48] to the Pitzer's acentric factor $\omega_{P}$ by
$\alpha_{R}=4.919 \omega_{P}+5.811$

Pitzer's acentric factor $\omega_{P}$ and Riedel's parameter $\alpha_{R}$ for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are determined by Eqs.(7) and (12). The obtained values are given in Table 4. For comparison, the values of $\omega_{p}$ and ${ }^{\alpha_{R}}$ calculated through other correlations are also presented in Table 4.

Table 4 - The Pitzer's acentric factor and the Riedel's parameter for $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$.

|  | $\omega_{p}$ |  |  | $\alpha_{R}$ |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| Correlation | $n-\mathrm{H}_{2}$ | $o-\mathrm{H}_{2}$ | $p-\mathrm{H}_{2}$ | $n-\mathrm{H}_{2}$ | Eq.(12) <br> $o-H_{2}$ | $p-\mathrm{H}_{2}$ |
| This Work | -0.2170 | -0.2175 | -0.2220 | 4.7438 | 4.7411 | 4.7191 |
| Brandani (B.II)[52] | -0.2565 | -0.2497 | -0.2511 | 4.5493 | 4.5830 | 4.5757 |
| Empirical correlation[49] | -0.2565 | -0.2491 | -0.2506 | 4.5494 | 4.5856 | 4.5783 |
| Edmister[49] | -0.2502 | -0.2427 | -0.2442 | 4.5803 | 4.6169 | 4.6096 |
| Brandani (B.I)[52] | -0.2447 | -0.2382 | -0.2397 | 4.6074 | 4.6392 | 4.6320 |
| Ambrose-Walton(II)[52] | -0.2287 | -0.2222 | -0.2240 | 4.6859 | 4.7179 | 4.7092 |
| CSGC-Reidel Equation[50] | -0.2270 | -0.2204 | -0.2221 | 4.6946 | 4.7267 | 4.7184 |
| Lee-Kesler[49] | -0.2242 | -0.2178 | -0.2195 | 4.7082 | 4.7398 | 4.7314 |
| Twu et al.[52] | -0.2232 | -0.2169 | -0.2187 | 4.7131 | 4.7442 | 4.7353 |
| Ambrose-Walton(I)[51] | -0.2098 | -0.2037 | -0.2053 | 4.7789 | 4.8089 | 4.8011 |

To evaluate the performance characteristics of Eqs.(11) and (12), percentage errors in the Pitzer's acentric factor and the Riedel's parameter determined in this work compared to other correlations are given in Table 5.

Table 5 -Percentage error in the Pitzer's acentric factor and the Riedel's parameter.

| Correlation for comparison | Percentage error \% in $\omega_{p}$ |  |  | Percentage error\% in $\alpha_{R}$ |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $n-\mathrm{H}_{2}$ | $o-\mathrm{H}_{2}$ | $p-\mathrm{H}_{2}$ | $n-\mathrm{H}_{2}$ | $o-\mathrm{H}_{2}$ | $p-\mathrm{H}_{2}$ |
| Brandani (B.II)[52] | 12.8793 | 15.4144 | 11.6135 | -3.4511 | -4.2748 | -3.1354 |
| Empirical correlation[49] | 12.6908 | 15.4128 | 11.4268 | -3.3913 | -4.2743 | -3.0768 |
| Edmister[49] | 10.4002 | 13.2834 | 9.1203 | -2.6898 | -3.5690 | -2.3771 |
| Brandani (B.I)[52] | 8.6964 | 11.3331 | 7.3907 | -2.1965 | -2.9605 | -1.8812 |
| Ambrose-Walton(II)[52] | 2.1259 | 5.1483 | 0.8980 | -0.4926 | -1.2361 | -0.2101 |
| CSGC-Reidel Equation[50] | 1.3344 | 4.4119 | 0.0694 | -0.3061 | -1.0492 | -0.0161 |
| Lee-Kesler[49] | 0.1191 | 3.2273 | -1.1389 | -0.0269 | -0.7559 | 0.2599 |
| Twu et al.[52] | -0.2874 | 2.7956 | -1.5008 | 0.0646 | -0.6512 | 0.3409 |
| Ambrose-Walton(I)[51] | -6.7681 | -3.3988 | -8.1186 | 1.4103 | 0.7340 | 1.7077 |

## Filippov's parameter

Filippov's parameter $\alpha_{F}$ is another measure of the temperature-dependence of vapour pressure in the vapour-liquid critical region. It is defined[48] by

$$
\begin{equation*}
\alpha_{F}=100 \times P_{v p}{ }^{*}\left(\text { at } T^{*}=0.625\right) \tag{13}
\end{equation*}
$$

Filippov's parameter for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are determined by Eq. (13). And its value for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ is $7.2909,7.3207$ and 7.0853 , respectively.

## Correlations between $\omega_{P}, \alpha_{R}, \alpha_{F}$ and the critical compressibility factor $Z_{c}$

The correlations between various characteristic parameters of substances are of significance due to their ability to predict their properties and to reveal the underlying theory. In this respect, Pitzer's acentric factor, Riedel's parameter, Filippov's parameter for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ obtained in this work and the critical compressibility factor[53-58] are shown in Figs. $6 \mathrm{a}, 6 \mathrm{~b}, 6 \mathrm{c}, 6 \mathrm{~d}$ and 6 e .


Fig.6a - $\alpha_{F}$ versus $\alpha_{R}$



Fig.6b $-\omega_{P}$ versus $\alpha_{F}$
Fig.6c - $\alpha_{R}$ versus $Z_{c}[53-58]$

Fig.6d - $\omega_{P}$ versus $Z_{c}[53-58]$


Fig.6e - $\alpha_{F}$ versus $Z_{c}[53-58]$

As seen, these parameters for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ may be correlated by a second degree polynomial of the form:
$Y=a_{0}+a_{1} X+a_{2} X^{2}$

Where $X$ and $Y$ are characteristic parameters such as $\omega_{p}, \alpha_{R}, \alpha_{F}$ and $Z_{c}$. The coefficients in Eq.(14) are presented in Table 6.

Table 6 - The coefficients in Eq.(14).

| $X$ | $Y$ | $a_{0}$ | $a_{1}$ | $a_{2}$ | $R$ | $R^{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\alpha_{R}$ | $\alpha_{F}$ | -19733.2 | 8336.07 | -880.042 | 0.9999 | 0.9998 |
| $\alpha_{F}$ | $\omega_{p}$ | -13.0951 | 3.5589 | -0.2459 | 0.9999 | 0.9998 |
| $Z_{c}$ | $\alpha_{R}$ | -3681.39 | 24181.7 | -39630 | 0.9999 | 0.9998 |
| $Z_{c}$ | $\omega_{p}$ | -97.8133 | 640.25 | -1050 | 0.9999 | 0.9998 |
| $Z_{c}$ | $\alpha_{F}$ | -467.133 | 3095.08 | -5075 | 0.9999 | 0.9998 |

And, the ratios $\alpha_{F} / \alpha_{R}$ and $\alpha_{F} / \omega_{P}$ for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are presented in Table 7.
Table 7- $\alpha_{F} / \alpha_{R}$ and $\alpha_{F} / \omega_{P}$ for $n-H_{2}, o-H_{2}$ and $p-H_{2}$.

| Parameter | $n-\mathrm{H}_{2}$ | $o-\mathrm{H}_{2}$ | $p-\mathrm{H}_{2}$ |
| :---: | :---: | :---: | :---: |
| $\frac{\alpha_{F}}{\alpha_{R}}$ | 1.5369 | 1.5441 | 1.5014 |
| $\frac{\alpha_{F}}{\omega_{p}}$ | -33.5986 | -33.5812 | -31.9158 |

## Results and analysis

Eq.(1) and Fig. 1 reveal that the vapor pressure of $n-H_{2}, o-H_{2}$ and $p-H_{2}$ has a parabolic dependence on temperature. This vapor pressure correlation is characterized by the correlation coefficient $R$ of 0.9976 and the coefficient determination $R^{2}$ of 0.9952 . That's the, the vapor pressure-temperature correlation formulated in this work is excellent from the statistical consideration. Hence, this correlation may be used as a basis for determining other thermodynamic properties of $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$. Based on Eq.(1), the acentric factors of $n-H_{2}, o-H_{2}$ and $p-H_{2}$ have been determined over a wide range of temperature. It is shown that the temperature dependence of the acentric factor for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ may be represented by a second degree polynomial with the correlation coefficient $R$ of 0.9992 and the coefficient determination $R^{2}$ of 0.9985 . Thus, the acentric factor correlation given by Eq.(7) is an excellent one from the statistical view point.

The Pitzer's acentric factors for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are found to be $-0.2170,-0.2175$ and -0.2220 , respectively. The percentage errors in the obtained values of the Pitzer's acentric factor of $n-H_{2}, o-H_{2}$ and $p-H_{2}$ compared to that of other known correlations[49-52] are in the range of about $0.07 \%$ to $15 \%$. That is, the correlation given by Eq. (9) reliably accommodates the Pitzer's acentric factors of $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$ from the literature[49-52]. The values of the Riedel's parameters for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are 4.7438, 4.7411 and 4.7191, respectively. The percentage errors in the obtained values of the Riedel's parameter of $n-\mathrm{H}_{2}, o-\mathrm{H}_{2}$ and $p-\mathrm{H}_{2}$ compared to that of other known correlations[49-52] are in the range of about $-0.02 \%$ to $4.28 \%$. The values of the Filippov's parameter for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ are found to be $7.2909,7.3207$ and 7.0853 , respectively. The ratio of the Filippov's parameter to the Riedel's parameter for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ is about 1.53, 1.54 and 1.50 , respectively. And, the ratio of the Filippov's parameter to the Pizter's acentric factor for $n-\mathrm{H}_{2}, o-H_{2}$ and $p-\mathrm{H}_{2}$ is about $-33.03,-33.58$ and -31.92 , respectively. As seen from Eq.(14) and Table 6, the Pitzer's acentric factor, Riedel's parameter, Filippov's parameter and the critical compressibility factor of $n-H_{2}, o-H_{2}$ and $p-H_{2}$ may be correlated by a second degree polynomial. For $n-H_{2}$, $o-H_{2}$ and $p-H_{2}$, the curvature of the vapor pressure curve is found to have a maximum at $17.11 \mathrm{~K}, 17.12 \mathrm{~K}$ and 17.00 K , respectively.

## Conclusion

Temperature-dependence correlations of vapor pressure and acentric factor for $n-\mathrm{H}_{2}, \mathrm{o}-\mathrm{H}_{2}$ and $p-H_{2}$, have been formulated. The obtained correlations are statistically excellent. The characteristic parameters such as the Pitzer's acentric factor, Riedel's parameter, Filippov's parameter have been determined for $n-H_{2}, o-H_{2}$ and $p-H_{2}$. And, the curvatures of vapor pressure curve for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ have been determined in a wide range of temperature.

It is found that the curvatures of vapor pressure curve for $n-H_{2}, o-H_{2}$ and $p-H_{2}$ have a maximum at about $17.11 \mathrm{~K}, 17.12 \mathrm{~K}$ and 17.00 K , respectively.

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